

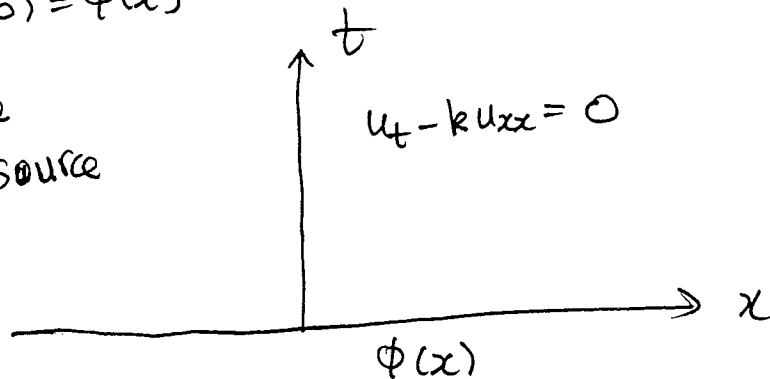
Chapter 2.4 Diffusion Equation on the whole line Lecture 7 ①

AIM: Obtain explicit solution formula for diffusion equation on the whole line \mathbb{R}

$$(IVP) \begin{cases} u_t = k u_{xx} + f(x,t), & -\infty < x < \infty, \quad 0 < t < +\infty \\ u(x,0) = \phi(x) \end{cases}$$

← source

First we assume that \exists no source



Recall: for wave equation, we can factorize the equation.
For diffusion eqn, we can't factorize.

Idea: $u_t = k u_{xx} \rightarrow$

- ① Find a particular solution
- ② use properties of $u_t - k u_{xx}$ to get general solutions

Properties of diffusion eqn: $u_t - k u_{xx} = 0$

- (a) $u(x-y, t)$ is another sol'n (translation invariance)
- (b) u_x, u_t, u_{xx} , etc is again a solution
- (c) linear combination of sol'n's is again a sol'n

(d) An integral is a sol'n

(2)

$$v(x, t) = \int_{-\infty}^{\infty} u(x-y, t) g(y) dy \quad \text{is a sol'n}$$

$$(\because \approx \sum u(x-y_i, t) g(y_i) \Delta y_i)$$

(c) (Dilation Invariance) If $u(x, t)$ is a solution, then the

dilation function

$u(\sqrt{a}x, at)$ is also a sol'n, $\forall a > 0$

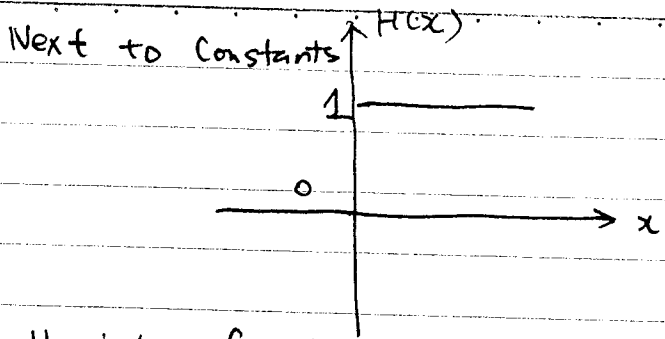
$$\bar{x} = \sqrt{a}x, \bar{t} = at, \quad \frac{\bar{x}^2}{\bar{t}} = \frac{x^2}{t} \quad (\text{invariance}).$$

We will use property (c) to obtain a special solution.

First we also need to choose an initial data

$\phi(x)$ (different from constant) which has the

dilation invariance.



$$H(x) = \begin{cases} 1, & x > 0 \\ 0, & x < 0 \end{cases}$$

Heaviside function

solve:

$$\begin{cases} Q_t - k Q_{xx} = 0 \\ Q(x, t) = H(x) \end{cases}$$

Since $H(x)$ is dilation invariant, we look for solns of the following form:

$$Q(x, t) = g(p), \quad p = \frac{x}{\sqrt{4kt}}$$

This reduces the problem to an ODE

Let us compute

$$Q_t = \frac{dg}{dp} \frac{\partial p}{\partial t} = -\frac{1}{2t} \frac{x}{\sqrt{4kt}} g'(p)$$

$$Q_{xx} = \frac{1}{4kt} g''(p), \quad Q_{xxx} = \frac{1}{4kt} g'''(p)$$

$$g_{pp} + 2pg_p = 0$$

$$g(p) = c_1 \int_0^p e^{-p^2} dp + c_2$$

$$\begin{aligned} Q(x, t) &= c_1 \int_0^p e^{-p^2} dp + c_2 \\ &= c_1 \int_0^{\frac{x}{\sqrt{4kt}}} e^{-p^2} dp + c_2 \end{aligned}$$

We now decide the two constants

If $x > 0$, $1 = \lim_{t \rightarrow 0} Q = c_1 \frac{\sqrt{\pi}}{2} + c_2$

If $x < 0$, $0 = \lim_{t \rightarrow 0} Q = -c_1 \frac{\sqrt{\pi}}{2} + c_2$

$$Q(x, t) = \frac{1}{2} + \frac{1}{\sqrt{4kt}} \int_0^{\frac{x}{\sqrt{4kt}}} e^{-p^2} dp$$

Now we differentiate

$$S(x, t) = \frac{\partial Q}{\partial x} \text{ is again a sol'n}$$

$$Q = \frac{1}{\sqrt{4kt}} e^{-\frac{x^2}{4kt}} \text{ for } t > 0$$

theorem: $u(x, t) = \int_{-\infty}^{+\infty} S(x-y, t) \phi(y) dy, \text{ for } t > 0$

proof: $u(x, t) = \int_{-\infty}^{+\infty} \frac{\partial Q}{\partial x}(x-y, t) \phi(y) dy$

$$= - \int_{-\infty}^{+\infty} \frac{\partial}{\partial y} [Q(x-y, t)] \phi(y) dy$$

$$= \int_{-\infty}^{+\infty} Q(x-y, t) \phi'(y) dy - Q(x-y, t) \phi(y) \Big|_{y=-\infty}^{y=+\infty}$$

$$u(x, 0) = \int_{-\infty}^{+\infty} S(x-y, 0) \phi'(y) dy = \phi \Big|_{-\infty}^x = \phi(x)$$

†

$$u(x, t) = \frac{1}{\sqrt{4kt}} \int_{-\infty}^{+\infty} e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy$$

$$S(x, t) = \frac{1}{\sqrt{4kt}} e^{-\frac{x^2}{4kt}} \text{ - fundamental sol'n's}$$

- Green's function

- Source Function

- Gaussian Potential

Properties of $S(x, t)$

$$1. \int_{-\infty}^{+\infty} S(x, t) dx = 1 \quad \Rightarrow \quad \int_{-\infty}^{+\infty} S(x-y, t) dy = 1, \quad \forall y$$

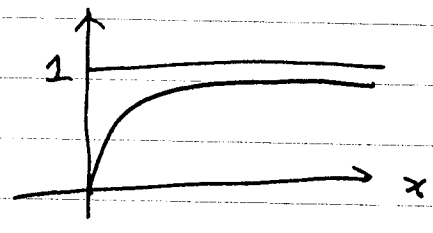
$$2. S(x, t) = S(-x, t)$$

$$3. \text{Fix } s > 0, \max_{|x| \geq s} S(x, t) \rightarrow 0 \text{ as } t \rightarrow 0$$

To use the formula, it is convenient to introduce the following notation

$$\text{Erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-p^2} dp$$

$$\text{Erf}(0) = 0, \quad \text{Erf}(+\infty) = 1$$



$$\text{x. 1. } Q(x) = \begin{cases} 1, & x > 0 \\ 0, & x < 0 \end{cases}$$

$$Q(x, t) = \frac{1}{2} + \frac{1}{2} \text{Erf}\left(\frac{x}{\sqrt{4kt}}\right)$$

$$2. \phi(x) = e^{-x}$$

$$u(x, t) = \frac{1}{\sqrt{4k\pi t}} \int_{-\infty}^{+\infty} e^{-\frac{(x-y)^2}{4kt}} e^{-y} dy$$

$$= \frac{1}{\sqrt{4k\pi t}} \int_{-\infty}^{+\infty} e^{-\frac{y^2 - 2xy + 4kt y + x^2}{4kt}} dy$$

$$= \frac{1}{\sqrt{4k\pi t}} \int_{-\infty}^{+\infty} e^{-\frac{(y+2kt-x)^2 + x^2 - (2kt-x)^2}{4kt}} dy$$

$$= \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{+\infty} e^{-\frac{(y+2kt-x)^2}{4kt}} dy e^{-\frac{4kt(x-4k^2t^2)}{4kt}}$$

$$= \int_{-\infty}^{+\infty} S(y-2kt+x, t) dy e^{-x+kt}$$

$$= e^{kt-x}$$

x.3. $\phi(x) = 1$ for $|x| < a$, $\phi(x) = 0$ for $|x| > a$

$$u(x,t) = \int_{-\infty}^{+\infty} S(x-y,t) \phi(y) dy$$

$$= \frac{1}{\sqrt{4\pi kt}} \int_{-a}^a e^{-\frac{(x-y)^2}{4kt}} dy \quad (kt y - x = \sqrt{4kt} p)$$

$$= \frac{1}{\sqrt{4\pi kt}} \int_{\frac{x-a}{\sqrt{4kt}}}^{\frac{x+a}{\sqrt{4kt}}} e^{-p^2} \sqrt{4kt} dp$$

$$= \frac{1}{\sqrt{\pi}} \int_{\frac{x-a}{\sqrt{4kt}}}^{\frac{x+a}{\sqrt{4kt}}} e^{-p^2} dp$$

$$= \frac{1}{2} \left(-\text{Erf} \left(\frac{a-x}{\sqrt{4kt}} \right) + \text{Erf} \left(\frac{a+x}{\sqrt{4kt}} \right) \right) \neq 0 \quad \forall x, t$$

This shows that diffusion eqn has ∞ speed

Diffusion with a source

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$$\begin{cases} u_t - k u_{xx} = f(x, t) & -\infty < x < +\infty, t > 0 \\ u(x, 0) = \phi(x) & \end{cases} \quad (*)$$

Theorem: The sol'n to (*) is given by

$$u(x, t) = \int_{-\infty}^{+\infty} S(x-y, t) \phi(y) dy + \int_0^t \int_{-\infty}^{+\infty} S(x-y, t-s) f(y, s) dy ds$$

Proof: Let $u_1(x) = \int_{-\infty}^{+\infty} S(x-y, t) \phi(y) dy$

$$u_2(x) = \int_0^t \int_{-\infty}^{+\infty} S(x-y, t-s) f(y, s) dy ds$$

Let us compute u_2 :

$$(u_2)_t = \int_{-\infty}^{+\infty} S(x-y, 0) f(y, t) dy$$

$$= f(x, t) + \int_0^t \int_{-\infty}^{+\infty} \frac{\partial}{\partial t} S(x-y, t-s) f(y, s) dy ds$$

$$= f(x, t) + \int_0^t \int_{-\infty}^{+\infty} k \frac{\partial^2 S}{\partial x^2}(x-y, t-s) f(y, s) dy ds$$

$$= f(x, t) + k \frac{\partial^2}{\partial x^2} \left[\int_0^t \int_{-\infty}^{+\infty} S(x-y, t-s) f(y, s) dy ds \right]$$

$$= f(x, t) + k \frac{\partial^2}{\partial x^2} u_2$$

If $\phi = 0$, $u(x, t) = \int_0^t \int_{-\infty}^{+\infty} \frac{1}{\sqrt{4\pi k(t-s)}} e^{-\frac{(x-y)^2}{4k(t-s)}} f(y, s) dy ds$

general solution:

$$u(x, t) = \int_{-\infty}^{+\infty} S(x-y, t) \phi(y) dy + \int_0^t \int_{-\infty}^{+\infty} S(x-y, t-s) f(y, s) dy ds$$

Idea $u_t - k u_{xx} = 0 \Rightarrow u = e^{-tA} \phi(0)$

$$u_t + Au = f(t) \Rightarrow u = e^{-tA} u(0) + \int_0^t e^{-(s-t)A} f(s) ds$$

Source on a half-line

$$\begin{cases} v_t - kv_{xx} = f(x,t), & x > 0 \\ v(x,0) = \phi(x) \\ v(0,t) = h(t) \end{cases}$$

Let $v(x,t) = u(x,t) - h(t)$

$$\begin{cases} u_t - kv_{xx} = f(x,t) - h'(t) \\ u(x,0) = \phi(x) - h(0) \\ u(0,t) = 0 \end{cases}$$

method of reflection

$$f_{\text{odd}} \rightarrow \phi_{\text{odd}} \rightarrow v_{\text{ext}} \rightarrow v(x,t)$$

Neumann problem

$$\begin{cases} v_t - kv_{xx} = f(x,t), & x > 0 \\ v(x,0) = \phi(x) \\ v_x(0,t) = h(t) \end{cases}$$

$$v(x,t) = \phi(x) - x h(t)$$