

Recall: Let $S(x, t) = \frac{1}{\sqrt{4k\pi t}} e^{-\frac{x^2}{4kt}}$, $t > 0$

be the source function. Then the solution to the diffusion equation

$$\begin{cases} u_t = k u_{xx}, & t > 0 \\ u(x, 0) = \phi(x), & -\infty < x < +\infty \end{cases}$$

is given by

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{4k\pi t}} \int_{-\infty}^{+\infty} e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy \\ &= \int_{-\infty}^{+\infty} S(x-y, t) \phi(y) dy \end{aligned}$$

$$\boxed{u := S(t) \phi}$$

Next we want to solve diffusion with a source

$$(*) \begin{cases} u_t - k u_{xx} = f(x, t), & -\infty < x < +\infty, t > 0 \\ u(x, 0) = \phi(x) \end{cases}$$

Theorem: The solution to (*) is given by

$$u(x, t) = \int_{-\infty}^{+\infty} S(x-y, t) \phi(y) dy + \int_0^t \int_{-\infty}^{+\infty} S(x-y, t-s) f(y, s) dy ds$$

(**)

Proof: We just need to show that the solution to (*) (2)

with $\phi \equiv 0$.

$$\text{Let } u(x, t) = \int_0^t \int_{-\infty}^{+\infty} s(x-y, t-s) f(y, s) dy ds$$

Let us compute

$$\begin{aligned} u_t &= \int_{-\infty}^{+\infty} S(x-y, 0) f(y, t) dy \\ &\quad + \int_0^t \int_{-\infty}^{+\infty} \frac{\partial}{\partial t} S(x-y, t-s) f(y, s) dy ds \\ &= f(x, t) + \int_0^t \int_{-\infty}^{+\infty} k \frac{\partial^2 S}{\partial x^2}(x-y, t-s) f(y, s) dy ds \\ &= f(x, t) + k \frac{\partial^2}{\partial x^2} \left[\int_0^t \int_{-\infty}^{+\infty} S(x-y, t-s) f(y, s) dy ds \right] \\ &= f(x, t) + k \frac{\partial^2}{\partial x^2} u \end{aligned}$$

$$\text{so } u_t - k u_{xx} = f(x, t)$$

On the other hand, clearly we have

$$u(x, 0) = 0$$

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General solution:

$$u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{+\infty} e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy + \int_0^t \int_{-\infty}^{+\infty} \frac{1}{\sqrt{4\pi k(t-s)}} e^{-\frac{(x-y)^2}{4k(t-s)}} f(y, s) dy ds$$

Example 1: Solve

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$$\begin{cases} u_t - k u_{xx} = 1 \\ u(x, 0) = 0 \end{cases}$$

Solution:

$$u(x, t) = \int_0^t \int_{-\infty}^{+\infty} \frac{1}{\sqrt{4\pi k(t-s)}} e^{-\frac{(x-y)^2}{4k(t-s)}} dy ds$$

$$= \int_0^t 1 ds = t$$

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Now we discuss well-posedness of diffusion equation

$$\begin{cases} u_t - k u_{xx} = f(x, t), & -\infty < x < +\infty, t > 0 \\ u(x, 0) = \phi(x) \end{cases}$$

Well-posedness:

(1) Existence: Given $(f(x, t), \phi(x))$, there exists

a solution

$$u(x, t) = \int_{-\infty}^{+\infty} S(x-y, t) \phi(y) dy + \int_0^t \int_{-\infty}^{+\infty} S(x-y, t-s) f(y, s) dy ds$$

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② Uniqueness: Given (f, ϕ) , the solution is unique.
 discuss later

③ stability: From the formula (**), we have

$$\begin{aligned} \max_{0 \leq t \leq T} |u(x, t)| &\leq \int_{-\infty}^{+\infty} S(x, y, t) |\max |\phi(y)|| dy \\ &+ \int_0^t \int_{-\infty}^{+\infty} S(x, y, t-s) \max_{0 \leq s \leq T} |f(y, s)| dy ds \\ &\leq \max |\phi| + T \max_{\substack{-\infty < x < \infty \\ 0 \leq t \leq T}} |f(x, t)| \\ &\leq (1+T) (\max |\phi| + \max |f(x, t)|) \end{aligned}$$

So if ϕ, f are small, then u is small.

Finally, we prove uniqueness: by energy method:

Consider diffusion equation

$$u_t - k u_{xx} = 0, \quad -\infty < x < +\infty, \quad t > 0$$

Consider its energy

$$E(t) = \frac{1}{2} \int_{-\infty}^{+\infty} u_x^2(x, t) dx$$

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$$\text{let } E(t) = \frac{1}{2} \int_{-\infty}^{+\infty} v^2(x,t) dx$$

$$\text{Then } E(0) = \frac{1}{2} \int_{-\infty}^{+\infty} v^2(x,0) dx = 0$$

$$\frac{dE}{dt} = -k \int_{-\infty}^{+\infty} v_x^2 \leq 0$$

$$\text{So } 0 \leq E(t) \leq E(0) = 0 \text{ for } t \geq 0$$

$$\Rightarrow E(t) = 0, \forall t > 0 \Rightarrow v(x,t) \equiv 0, \forall x, \forall t > 0 \quad \#$$

Now we go back to wave equation.

$$\begin{cases} u_{tt} - c^2 u_{xx} = f(x) \\ u(x,0) = \phi(x), u_t(x,0) = \psi(x) \end{cases}$$

We can use similar energy method to prove uniqueness:

$$\text{Let } E(t) = \frac{1}{2} \int_{-\infty}^{+\infty} u_t^2(x,t) dx + \frac{1}{2} c^2 \int_{-\infty}^{+\infty} u_x^2(x,t) dx$$

Then

$$\begin{aligned} \frac{dE}{dt} &= \int (u_t u_{tt} + c^2 u_x u_{xt}) \\ &= \int (u_t c^2 u_{xx} + c^2 u_x u_{xt}) \\ &= c^2 \int (u_t u_x)_x dx \\ &= 0 \end{aligned}$$

Then

$$\begin{aligned}
\frac{dE}{dt} &= \int_{-\infty}^{\infty} u_{xx} u_{xxx} dx \\
&= k \int_{-\infty}^{\infty} u_x u_{xx} dx \\
&= -k \int_{-\infty}^{\infty} u_x^2 dx + k \int_{-\infty}^{\infty} (u u_x)_x dx \\
&= -k \int_{-\infty}^{\infty} u_x^2 dx + k u u_x \Big|_{-\infty}^{\infty} \\
&= -k \int_{-\infty}^{\infty} u_x^2 dx
\end{aligned}$$

Consequence: $\frac{dE}{dt} \leq 0 \Rightarrow$ energy is decreasing
 \Rightarrow dissipative

Proof of Uniqueness: Suppose there are two solutions

u_1, u_2 to the problem

$$\begin{cases} u_t - k u_{xx} = f(x, t), & t > 0 \\ u(x, 0) = \phi(x) \end{cases}$$

Then $v(x, t) = u_1 - u_2$ satisfies

$$\begin{cases} v_t - k v_{xx} = 0, & t > 0 \\ v(x, 0) = 0 \end{cases}$$

As a consequence, for wave equation, the energy is "conserved"

Conservation of Energy: $E(t) = E(0), \forall t$

Proof of uniqueness: As before, $u_t = u_1 - u_2$ satisfies

$$\begin{cases} u_{tt} - c^2 u_{xx} = 0 \\ u(x,0) = 0, u_t(x,0) = 0 \end{cases}$$

So $u_x(x,0) = 0, u_t(x,0) = 0$

$\Rightarrow E(0) = 0 = E(t) \Rightarrow u_t = u_x = 0 \Rightarrow u \equiv \text{Constant}$

$\Rightarrow u \equiv 0$. #

Summary:	Formula	Energy	Well-posedness	Discontinuity
Diffusion Equation $\begin{cases} u_t - k u_{xx} = f \\ u(x,0) = \phi(x) \end{cases}$	$u = \int_0^x S(x-y,t) \phi(y) dy + \int_0^t \int_0^x S(x-y,t-s) f(y,s) dy ds$	decreasing	Yes	becomes smooth when $t > 0$
Wave Eqn $\begin{cases} u_{tt} - c^2 u_{xx} = f \\ u(x,0) = \phi(x) \\ u_t(x,0) = \psi(x) \end{cases}$	$u = \frac{1}{2} [\phi(x-ct) + \phi(x+ct)] + \frac{c}{2c} \int_{x-ct}^{x+ct} \psi(s) ds + \frac{1}{2c} \int_0^t \int_{x-c(t-s)}^{x+c(t-s)} f(y,s) dy ds$	conserved	Yes	Propagate along the characteristics

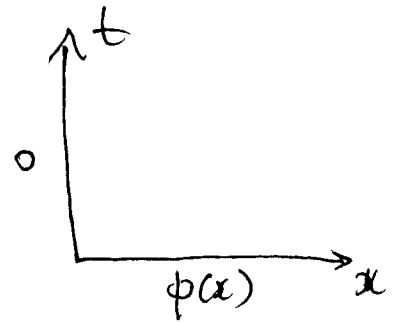
Diffusion Equation with boundaries: Method of Reflection

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Ex. 1

Dirichlet BC

$$\begin{cases} u_t - k u_{xx} = f(x, t), & x > 0, t > 0 \\ u(x, 0) = \phi(x) \\ u(0, t) = 0 \end{cases}$$



Solution: Reflection: extend ϕ, f , oddly to $x < 0$

$$f_{\text{ext}}(x, t) = \begin{cases} f(x, t), & x > 0 \\ -f(-x, t), & x < 0 \end{cases}$$

$$\phi_{\text{ext}}(x, t) = \begin{cases} \phi(x), & x > 0 \\ -\phi(-x), & x < 0 \end{cases}$$

Solve

$$\begin{cases} u_t - k u_{xx} = f_{\text{ext}}(x, t) \\ u(x, 0) = \phi_{\text{ext}} \end{cases}$$

$$u_{\text{ext}}(x, t) = \int S(x-y, t) \phi_{\text{ext}} + \int_0^t \int S(x-y, t-s) f_{\text{ext}}$$

check: $u_{\text{ext}}(0, t) = 0$.

By uniqueness (same proof as before) $u(x, t) = u_{\text{ext}}(x, t)$

Example 1: Solve

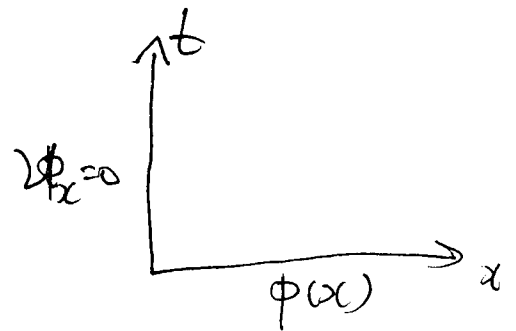
$$\begin{cases} u_t - k u_{xx} = 0, & x > 0 \\ u(x, 0) = 1 \\ u(0, t) = 0 \end{cases}$$

Solution

$$\begin{aligned}
u(x,t) &= \int_{-\infty}^{+\infty} S(x-y, t) \phi_{\text{ext}}(y) dy \\
&= \int_{-\infty}^{+\infty} S(x-y, t) dy - \int_{-\infty}^0 S(x-y, t) dy \\
&= \int_0^{+\infty} S(x-y, t) dy - \int_0^{+\infty} S(x+y, t) dy \\
&= \int_{-\infty}^{\frac{x}{\sqrt{4kt}}} e^{-p^2} dp / \sqrt{\pi} - \int_{\frac{x}{\sqrt{4kt}}}^{+\infty} e^{-q^2} dq / \sqrt{\pi} \\
&= \left[\frac{1}{2} + \frac{1}{2} \operatorname{Erf} \left(\frac{x}{\sqrt{4kt}} \right) \right] - \left[\frac{1}{2} - \frac{1}{2} \operatorname{Erf} \left(\frac{x}{\sqrt{4kt}} \right) \right] \\
&= \operatorname{Erf} \left(\frac{x}{\sqrt{4kt}} \right) \quad \#
\end{aligned}$$

Ex.2 Neumann BC

$$\begin{cases}
u_t - k u_{xx} = f(x, t), & x > 0, t > 0 \\
u(x, 0) = \phi(x) \\
u_x(0, t) = 0
\end{cases}$$



Solution: Reflection: extend ϕ, f evenly to $x > 0$

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$$f_{\text{ext}}(x,t) = \begin{cases} f(x,t), & x > 0 \\ f(-x,t), & x < 0 \end{cases}$$

$$\phi_{\text{ext}}(x,t) = \begin{cases} \phi(x), & x > 0 \\ \phi(-x), & x < 0 \end{cases}$$

Solve $\begin{cases} u_t - k u_{xx} = f_{\text{ext}}(x,t) \\ u(x,0) = \phi_{\text{ext}} \end{cases}$

$$u_{\text{ext}}(x,t) = \int S(x-y,t) \phi_{\text{ext}} + \int_0^t \int S(x-y,t-s) f_{\text{ext}}$$

check: $u_{\text{ext},x}(0,t) = 0$

By uniqueness, $u(x,t) = u_{\text{ext}}(x,t)$

Example 2 Solve $\begin{cases} u_t - k u_{xx} = 0, & x > 0 \\ u(x,0) = 1 \\ u(0,t) = 0 \end{cases}$

Sol'n: $u(x,t) = \int S(x-y,t) \phi_{\text{ext}}(y) dy = \int S(x-y,t) dy = 1$
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Source on a half-line

$$\text{Ex. 3} \quad \begin{cases} u_t - k u_{xx} = f(x, t) \\ u(0, t) = h(t) \\ u(x, 0) = \phi(x) \end{cases}$$

Sol'n: Let $V(x, t) = u(x, t) - h(t) \Rightarrow$

$$\begin{cases} V_t - k V_{xx} = f(x, t) - h'(t) \\ V(0, t) = 0 \\ V(x, 0) = \phi(x) - h(0) \end{cases}$$

$$\text{Ex. 4} \quad \begin{cases} u_t - k u_{xx} = f(x, t) \\ u'_x(0, t) = h(t) \\ u(x, 0) = \phi(x) \end{cases}$$

Sol'n: Let $V(x, t) = u(x, t) - x h(t)$

$$\begin{cases} V_t - k V_{xx} = f(x, t) \\ V'_x(0, t) = 0 \\ V(x, 0) = \phi(x) - x h(0) \end{cases}$$

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What about two boundary conditions

\Rightarrow Lecture 9

$$\begin{cases} u_t - k u_{xx} = f(x, t), 0 < x < l \\ u(x, 0) = \phi(x) \\ u(0, t) = 0, u(l, t) = 0 \end{cases}$$