



$$\begin{aligned}
|u(x_0) - u(x_0)| &= \left| \int_{\partial B} K(x, y) (\varphi(y) - \varphi(x_0)) ds_y \right| \\
&= \int_{|y-x_0| \leq \delta} K(x, y) |\varphi(y) - \varphi(x_0)| ds \\
&\quad + \int_{|y-x_0| > \delta} K(x, y) |\varphi(y) - \varphi(x_0)| ds_y \\
&\leq \varepsilon + 2M \frac{(R^2 - |x|^2) R^{n-2}}{(\delta/2)^n}
\end{aligned}$$

#

Other Green's Functions.

Ex. Half-space 

Ex. Half-ball 

Ex. Quarter ball 

Finally, we are ready to discuss Perron's Method. The main goal of Perron's Method is the following theorem

Theorem I.13. Under some smoothness condition on  $\Omega$  (e.g.  $\partial\Omega \in C^2$ ), the

$\forall g \in C^0(\partial\Omega), \exists!$  a sol'n to

$$\begin{cases} \Delta u = 0 & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases}$$

We first define:  $u \in C^2(\Omega)$  is called subharmonic (superharmonic) in a domain  $\Omega$  if  $\Delta u \geq 0$  (or  $\Delta u \leq 0$ ) in  $\Omega$

Def: A fcn  $u \in C^0(\Omega)$  is called subharmonic in  $\Omega$  if  $\forall B \subset \subset \Omega$  and  $\forall$  harmonic  $h$  in  $B$ , i.e.  $h \in C^2(B) \cap C(\bar{B})$  and  $\Delta h = 0$  in  $B$ , satisfying  $u \leq h$  on  $\partial B$  we have  $u \leq h$  in  $B$

Corollary: A harmonic function in  $\Omega$  is both superharmonic and a sub-harmonic fcn.

Lemma 1. (Strong M.P.). Assume  $\Omega$  is connected. If a subharmonic fcn  $u$  attains its supremum in  $\Omega$ , then  $u \equiv \text{Constant}$  in  $\Omega$ .

Lemma 2.  $u$  subharmonic,  $v$  superharmonic,  $u - v \leq 0$  on  $\partial\Omega$

Then either  $v > u$  in  $\Omega$  or  $v \equiv u$ .

Proof:  $\{x \mid \sup(u-v) = M\} = \Omega_1$

open:  $x_1 \in \Omega_1$ ,  $B_p(x_1) \subset \subset \Omega$ . Then  $B_p(x_1) \in \Omega_1$ . If not,  $\exists$  a ball  $B_p(x^2)$ ,  $0 < p < p_0$ ,  $x^2 \in \partial B$ ,  $(u-v)(x^2) < M$ . Let  $h_1, h_2$  be harmonic in  $B$ ,  $h_1 = u$  on  $\partial B$ ,  $h_2 = v$  on  $\partial B$ . Then, if  $x \in B$

$$M \geq \max_{\partial B} (u-v) = \max_{\partial B} (h_1 - h_2)$$

$$\geq h_1(x) - h_2(x) \geq u(x) - v(x)$$

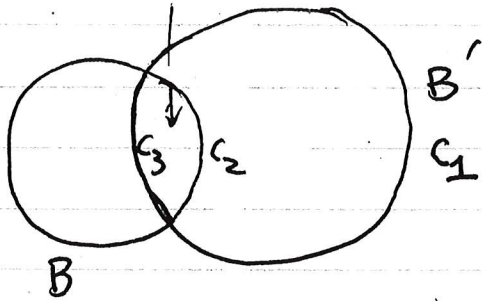
$$\text{Set } x = x^1 \Rightarrow u(x^1) - v(x^1) = M \Rightarrow h_1 - h_2 \equiv \text{Constant} \quad \#$$

Harmonic Lifting. Let  $u$  be subharmonic in  $\Omega$ ,  $B \subset \subset \Omega$  a ball,  $\bar{u}$  harmonic in  $B$  s.t.  $\bar{u} = u$  on  $\partial B$ . Then

$$U(x) = \begin{cases} \bar{u} & : x \in B \\ u(x) & : x \in \Omega \setminus \bar{B} \end{cases} \quad \text{is called a harmonic lifting of } u \text{ in } \bar{B}$$

Lemma 3.  $U$  is subharmonic in  $\Omega$

Proof: Let  $B' \subset \subset \Omega$  and  $h$  harmonic in  $B'$ ,  $h \geq U$  on  $\partial B'$ . We have to show that  $h \geq U$  in  $B'$ .



$$\begin{cases} \Delta h = \Delta \bar{u} = 0 \\ h \geq u = \bar{u} \text{ on } C_2 \\ h \geq U = \bar{u} \text{ on } C_3 \end{cases}$$

$$\Rightarrow h \geq \bar{u} = U \text{ on } B' \cap B$$

on  $C_1 = \partial B' \setminus B$  we have  $h \geq U \equiv u$

$$\underbrace{B' \setminus B}_{h \geq u} \begin{cases} \Delta h = 0 \text{ in } B' \\ h \geq U = u \text{ on } C_1 \\ h \geq U = \bar{u} \geq u \text{ on } C_3 \end{cases}$$

on  $C_3 = \partial B' \cap B$ , according to the definition of  $\bar{u}$ ,  $h \geq U = \bar{u} \geq u$

on  $C_2 = \partial B \cap B'$ ,  $\bar{u} = u = U$

Combining these inequalities,  $\Rightarrow h \geq u$  on  $\partial B'$ , hence  $h \geq u$  in  $B'$

Then  $U \leq h$  in  $B' \setminus B$

Since  $U \equiv u$  in  $B' \setminus B$ . It remains to show that  $U \leq h$  in  $B' \cap B$ .

on  $\partial(B' \cap B)$  we have  $h \geq U$ , and assumption  $h \geq U$  in  $B'$ . since

$U \equiv \bar{u}$  in  $B \cap B'$ ,  $h$  is harmonic  $\Rightarrow h \geq U$  in  $B \cap B'$

Lemma 4. Let  $u_1, \dots, u_n$  be subharmonic in  $\Omega$ . Then

$u = \max \{u_1, \dots, u_n\}$  is also subharmonic

Def:  $u$  is subharmonic,  $u \leq \phi$  on  $\partial \Omega \Leftrightarrow u$  is a sub-sol'n

Lemma 5.  $u$  is a subfcn,  $\bar{u}$  is a superfcn with respect to  $\varphi$ .  
Then  $u \leq \bar{u}$  in  $\Omega$

Set  $S_\varphi = \{v \in C(\Omega) \text{ subharmonic in } \Omega, v \leq \varphi \text{ on } \partial\Omega\}$ .

First  $S_\varphi$  is not empty,  $v = \inf_{\partial\Omega} \varphi$

Theorem I.14. (Perron). The function

$$u(x) := \sup_{v \in S_\varphi} v(x)$$

is harmonic in  $\Omega$

Proof: (i) We have in  $\Omega$   
 $\inf_{\partial\Omega} \varphi \leq u(x) \leq \sup_{\partial\Omega} \varphi$

(ii) Let  $y \in \Omega$  be fixed. Then  $\exists$  a sequence  $v_n \in S_\varphi$  with  $\lim_{n \rightarrow +\infty} v_n(y) = u(y)$ .

Let  $B = B_R(y) \subset \subset \Omega$ ,  $R$  sufficiently small, and let  $V_n$  be the harmonic lifting of  $v_n$  in  $B$ . Then  $v_n \in S_\varphi$

$$\lim_{n \rightarrow +\infty} V_n(y) = u(y) \quad - (2.5)$$

Proof of (2.5):  $v_n(y) \leq V_n(y)$  since  $v_n = V_n$  on  $\partial B$ ,  $V_n$  is harmonic in  $B$  and  $v_n$  is subharmonic.

$$u(y) \leq \lim_{n \rightarrow +\infty} V_n(y)$$

On the other hand, since  $v_n \in S_\varphi$ , we have

$$v_n(y) \leq \sup_{v \in S_\varphi} v(y) = u(y)$$

and  $u(y) \leq \lim_{n \rightarrow +\infty} v_n \leq u(y)$ .



(iii)  $\forall h$  harmonic in  $B$ , we have

$$\sup_{B_p(y)} |D^k h| \leq C \sup_{B_R} |h|$$

$\exists$  a subsequence  $V_{n_k} \rightarrow$  uniformly in  $B_p(y)$  to a harmonic fcn  $v$ . and

$$v(x) \leq u(x), \quad x \in B_R(y)$$

since  $V_n(x) \leq u(x)$  on  $B_R(y)$ .

At the center  $y$  it is,

$$v(y) = u(y). \quad (9)$$

(iv). Claim  $v(x) = u(x)$ ,  $x \in B$ .

Proof: If not,  $\exists z \in B$ ,  $v(z) < u(z)$ . Then  $\exists u_0 \in S_p$ ,  $v(z) < u_0(z)$ .

Set

$$w_k(x) = \max(u_0(x), V_{n_k}(x)).$$

Let  $W_k$  be the harmonic lifting of  $w_k$  in  $B$ .

A subsequence of  $W_k$  converges uniformly on cpt of  $B$  to  $w$  harmonic in  $B$  s.t.

$$v(x) \leq w(x) \leq u(x), \quad x \in B = B_R(y). \quad (10)$$

By (9)-(10),  $v(y) = w(y) = u(y)$ .

Since  $\Delta v = \Delta w = 0$ .

Strong M.P.  $v(x) = w(x)$ ,  $x \in B$ .

$$w(z) = v(z) < u_0(z).$$

By the definition of  $w_n(x)$  and  $W_n$

$$u_0(x) \leq w_{n_k}(x) \leq W_{n_k}(x)$$

$$u_0(x) \leq w(x), \quad x \in B \quad \Rightarrow \quad v(z) < u_0(z) \leq w(z), \quad \text{Contradiction!}$$

## Boundary Behavior

Def: A  $C(\bar{\Omega})$ -fcn  $w = w_{\xi}$  is called a barrier at  $\xi \in \partial\Omega$  relative to  $\Omega$  if

(i)  $w$  is superharmonic in  $\Omega$

(ii)  $w > 0$  in  $\bar{\Omega} \setminus \{\xi\}$  and  $w(\xi) = 0$

An important feature of the barrier concept is that it is a local property of the boundary  $\partial\Omega$ . Namely, let us define  $w$  to be a local barrier at  $\xi \in \partial\Omega$  if  $\exists$  a nbhd  $N$  of  $\xi$  such that  $w$  satisfies the definition in  $\Omega \cap N$ .

Let  $B$  be a ball satisfying  $\xi \in B \subset \subset N$  and  $m = \inf_{N-B} w > 0$ .

$$\bar{w}(x) = \begin{cases} \min(m, w(x)), & x \in \bar{\Omega} \cap B \\ m, & x \in \bar{\Omega} \setminus B \end{cases}$$

is then a barrier at  $\xi$  relative to  $\Omega$ .

$\xi$  is called a regular point if  $\exists$  a local barrier

Theorem I.14. Let  $u$  be a harmonic fcn defined in  $\Omega$  by the Perron method with boundary data  $\phi$ . If  $\xi$  is a regular point of  $\partial\Omega$  and if  $\phi$  is continuous at  $\xi$ , then

$$\lim_{x \rightarrow \xi, x \in \Omega} u(x) = \phi(\xi)$$

Proof: Fix  $\epsilon > 0$ . Then  $\exists \delta > 0$  s.t.  $|\phi(x) - \phi(\xi)| < \epsilon$ ,  $\forall x \in \partial\Omega, |x - \xi| < \delta$ .

Let  $M = \sup_{\partial\Omega} |\phi|$ . Let  $w$  be a barrier at  $\xi$ . Then  $\exists$  a  $k = k(\epsilon)$  s.t.

$$k w(x) > 2M \quad \text{if } |x - \xi| > \delta$$

Then  $\phi(z) + \varepsilon + k w(x)$  is a super-sol'n relative to  $\varphi$ .

$\phi(z) - \varepsilon - k w(x)$  is a sub-sol'n

$$u(x) = \sup_{v \in S_\varphi} v(x)$$

$$\phi(z) - \varepsilon - k w(x) \leq u(x) \leq \phi(z) + \varepsilon + k w(x)$$

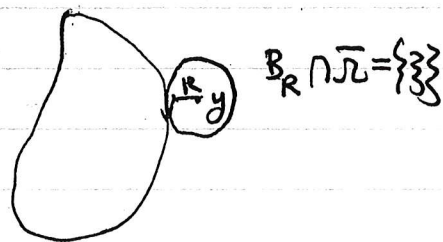
$$|u(x) - \phi(z)| \leq \varepsilon + k w(x).$$

#

Examples of Local barrier.

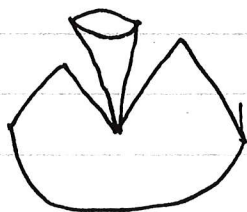
1). Exterior sphere.

$$w(x) = \begin{cases} R^{2-n} - |x-y|^{2-n}, & n \geq 3 \\ \log\left(\frac{|x-y|}{R}\right) \end{cases}$$



is a local barrier

2). Exterior cone.



$$r^\lambda f(\theta).$$

$$r^\lambda \omega(\mu\theta).$$

$$\Delta w = r^{\lambda-2} (r^2 - \mu^2) \omega(\mu\theta), \quad |\mu\theta| \leq \frac{\pi}{2}.$$

Perron's Method not easy to execute numerically. Another

method is Dirichlet Energy Method

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases} \quad (1)$$

$$E[u] = \frac{1}{2} \int_{\Omega} |\nabla u|^2 - \int_{\Omega} f u$$

$$\Lambda_0 = \left\{ \begin{array}{l} u = g \text{ on } \partial\Omega \\ u \in C^2(\Omega) \cap C(\bar{\Omega}) \end{array} \right\}$$

Theorem: 1) If  $u \in \Lambda_0$  is a sol'n to (1), then

$$E[u] \leq E[v], \quad \forall v \in \Lambda_0$$

2) If  $u \in \Lambda_0$  such that

$$E[u] \leq E[v], \quad \forall v \in \Lambda_0$$

then  $u$  satisfies (1).

Proof:  $v = u + w$ ,  $w = 0$  on  $\partial\Omega$

$$\int_{\Omega} |\nabla u + w|^2 = \int_{\Omega} 2\nabla u \cdot \nabla w + |\nabla w|^2$$

$$= 2 \int_{\Omega} \nabla(w \nabla u) - w \Delta u + |\nabla w|^2$$

$$= + \int_{\Omega} w f + |\nabla w|^2 + \int_{\Omega} w(-\Delta u)$$

$$E[u+w] = E[u] + \frac{1}{2} \int_{\Omega} |\nabla w|^2$$

$\Rightarrow$  true  $\forall w \in C^2(\Omega) \cap C(\bar{\Omega})$   
 $w=0$



So if we consider

$$\inf_{u \in \Lambda_0} E[u] = c_0$$

if  $c_0$  is attained by some  $u$  and  $u \in C^2(\Omega) \cap C(\bar{\Omega})$ , then we are done.

Problem 1: ~~Q~~ Is  $c_0$  attained?

Ex. 1. (Hadamard 1906): Let  $D = \{x \in \mathbb{R}^2 \mid |x|^2 < 1\}$

Let  $u : D \rightarrow \mathbb{R}$  be given by

$$u(r, \theta) = \sum_{n=1}^{+\infty} n^{-2} r^{n!} \sin(n\theta)$$

It is easy to check  $\Delta u = 0$

series converges absolutely uniformly in  $\bar{D}$ . Hence  $u$  is harmonic in  $D$  and continuous in  $\bar{D}$ .

On the other hand

$$\begin{aligned} E(u) &= \int_D |\nabla u|^2 \geq \int_0^{2\pi} \int_0^{\rho} |\nabla u|^2 r dr d\theta \\ &= \sum_{n=1}^{+\infty} \frac{\pi n!}{2n^4} \rho^{2n!} \geq \sum_{n=1}^m \frac{\pi n!}{2n^4} \rho^{2n!} \end{aligned}$$

$\forall \rho < 1, \forall$  integer  $m$ . Hence  $E(u) = \infty$ .

To conclude, there exists a Dirichlet datum  $g \in C(\partial D)$  for which the Dirichlet problem is perfectly solvable, but the sol'n cannot be obtained by minimizing the Dirichlet energy.

Ex. 2. Consider  $\begin{cases} \Delta u = 0 & \text{in } D = B_1 \\ u = 0 & \text{on } \partial B_1 \end{cases}$

$$k \in \mathbb{N}, \quad u_k(r, \theta) = \begin{cases} k a_k, & r < e^{-2k} \\ -a_k(k + \log r), & e^{-2k} < r < e^{-k} \\ 0, & e^{-k} < r < 1 \end{cases}$$

$$\begin{aligned} & \frac{n-2}{2} \eta\left(\frac{x}{\epsilon}\right) \quad (n \geq 3) \\ & \frac{1}{\sqrt{\log \frac{1}{\epsilon}}} \left[ \log(\epsilon^2 + r^2) - k \right] \end{aligned}$$

$u_k$  is continuous, piecewise smooth

$$E(u_k) = 2\pi \int_0^1 |2r u_k|^2 r dr = 2\pi a_k^2 \log r \Big|_{e^{-2k}}^{e^{-k}} = 2\pi k a_k^2$$

choosing  $a_k = k^{-2/3}$ . Then  $E(u_k) \rightarrow 0$

However  $u_k(\omega) = k a_k = k^{1/3}$  diverges as  $k \rightarrow +\infty$ .

---

Remedy:

- First, we show that  $E$  has a minimizer in a class that contains  $\Lambda_0$  as a subset (Sobolev space)
  - Show that minimizer is in fact in  $\Lambda_0$  (Regularity Theory)
- 

### Part I.3 Heat Equation

$$\begin{cases} u_t = \Delta u + f, & x \in \mathbb{R}^n, t > 0 \\ u(x, 0) = g(x), & x \in \mathbb{R}^n, t = 0 \end{cases}$$

$$u \in C^2, x \in \mathbb{R}^n, t > 0.$$

A formal solution is obtained by Fourier transformation. We first consider  $f=0$

$$\hat{u}(\xi) = (2\pi)^{-n/2} \int_{\mathbb{R}^n} e^{-i\xi \cdot x} u(x) dx$$

$$u(x) = (2\pi)^{-n/2} \int e^{i\xi \cdot x} \hat{u}(\xi) d\xi$$

$$\begin{cases} \hat{u}_t + |\mathbb{z}|^2 \hat{u} = 0 \\ \hat{u}(0) = \hat{g} \end{cases}$$

$$u(x, t) = (2\pi)^{-n/2} \int e^{i(x-\mathbb{z}) - |\mathbb{z}|^2 t} \hat{g}(\mathbb{z}) d\mathbb{z} = \int K(x, y, t) f(y) dy$$

$$K(x, y, t) = (2\pi)^{-n} \int e^{i(x-\mathbb{z}) - |\mathbb{z}|^2 t} d\mathbb{z}$$

$$\mathbb{z} = \frac{i(x-y)}{2t} + \frac{1}{\sqrt{4t}} \eta$$

$$\int e^{-|\eta|^2} d\eta = \left( \int_{-\infty}^{\infty} e^{-s^2} ds \right)^n = \pi^{n/2}$$

$$K(x, y, t) = (4\pi t)^{-n/2} e^{-|x-y|^2 / 4t}$$

We obtain Poisson's formula

Theorem 1: Let  $f$  be continuous and bounded for  $x \in \mathbb{R}^n$ . Then

$$u(x, t) = \int K(x, y, t) f(y) dy$$

$$u \in C^\infty, \quad u_t = \Delta u$$

$$\lim_{t \rightarrow 0} u(x, t) = f(x).$$

The proof follows from basic properties of  $K$

(a)  $K \in C^\infty$

(b)  $(\frac{\partial}{\partial t} - \Delta_x) K(x, y, t) = 0, \quad t > 0$

(c)  $K(x, y, t) > 0, \quad t > 0$

(d)  $\int K(x, y, t) dy = 1, \quad \forall x \in \mathbb{R}^n, t > 0$

(e)  $\forall \delta > 0, \quad \lim_{t \rightarrow 0} \int_{|x-y| > \delta} K(x, y, t) dy = 0, \quad \text{uniformly } \forall x \in \mathbb{R}^n$

then

$$|u(x,t) - f(z)| = \left| \int K(x,y,t) (f(y) - f(z)) dy \right|$$

$$\leq \int_{|x-y| < \delta} + \int_{|x-y| > \delta}$$

By the same type of argument one proves that if  $f$  is measurable and

$$|f(x)| \leq M e^{a|x|^2}, \quad a, M > 0$$

then  $u(x,t)$  satisfies  $u_t = \Delta u$ ,  $0 < t < \frac{1}{4a}$

infinite-speed

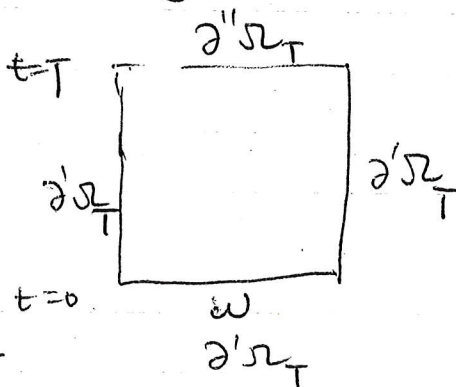
Maximum Principle, uniqueness, and regularity

$$\Omega_T = \{(x,t) \mid x \in \Omega, 0 < t < T\}$$

$$\partial' \Omega_T = \{(x,t) \mid x \in \partial \Omega, 0 \leq t \leq T \text{ or } x \in \Omega, t=0\}$$

$$\partial'' \Omega_T = \{(x,t) \mid x \in \Omega, t=T\}$$

Theorem:  $u_t - \Delta u \leq 0$ . Then  
 $\max_{\bar{\Omega}_T} u = \max_{\partial' \Omega_T} u$



Maximum Principle  $\Rightarrow$  uniqueness in  $\Omega_T$

Proof:  $v = u - kt$   
 $v_t - \Delta v < 0$  #

We can extend the maximum principle and uniqueness theorem to the case where  $\Omega$  is the "slab"

$$\Omega = \{(x,t) \mid x \in \mathbb{R}^n, 0 < t < T\}$$



Theorem.  $u_t - \Delta u \leq 0, \quad 0 < t < T, \quad x \in \mathbb{R}^n$

$$u(x, t) \leq M e^{ax^2} \quad 0 < t < T, \quad x \in \mathbb{R}^n$$

$$u(x, 0) = f(x) \quad x \in \mathbb{R}^n$$

Then  $u(x, t) \leq \sup_z f(z) \quad 0 \leq t \leq T, \quad x \in \mathbb{R}^n,$

Corollary:  $\begin{cases} u_t - \Delta u = 0 & 0 < t < T \\ u(x, 0) = f(x) \end{cases}$

is unique if  $|u(x, t)| \leq M e^{a|x|^2}, \quad 0 < t < T$

Proof: It is sufficient to show that under  $4aT < 1$  since we can divide  $0 < t < T$  into small intervals.

$$u(x, t) \leq \sup_y u(y, k\tau) \leq \sup_y u(y, 0) \quad k\tau \leq t \leq (k+1)\tau.$$

Let  $4a(T+\epsilon) < 1$ .

Given a fixed  $y, \mu > 0$

$$v_\mu(x, t) = u(x, t) - \mu (4\pi(T+\epsilon-t))^{-n/2} \exp[-|x-y|^2/4(T+\epsilon-t)]$$

~~$v_\mu(x, t)$~~   $0 \leq t \leq T$

$$\partial_t v_\mu - \Delta v_\mu = u_t - \Delta u \leq 0$$

Consider

$$\Omega_T = \{(x, t) \mid |x-y| < \rho, \quad 0 < t < T\}$$

$$v_\mu(y, t) \leq \max_{\partial' \Omega} v_\mu$$

on the plane part  $\partial' \Omega, \quad v_\mu(x, 0) \leq u(x, 0) \leq \sup_z f(z)$

on the curved part  $|x-y| = \rho, \quad 0 \leq t < T$

$$v_\mu(x, t) \leq M e^{a|x|^2} - \mu (4\pi(T+\epsilon-t))^{-n/2} \exp[-\rho^2/4(T+\epsilon-t)]$$

$$\leq M e^{a(|y|+\rho)^2} - \mu (4\pi(T+\epsilon))^{-n/2} e^{-\rho^2/4(T+\epsilon)}$$

$$\leq \sup f(z)$$

if  $\rho$  is large enough,  $\rho = \rho(\mu, \epsilon, T)$ .

$$\max_{\partial \Omega} V_{\mu}(y, t) \leq \sup f(z)$$

$$V_{\mu}(y, t) = u(y, t) - \mu (4(T+\epsilon-t))^{-n/2} \leq \sup f(z)$$

Letting  $\mu \rightarrow 0$ .

#

The growth condition is necessary.

$$\begin{cases} u_t = u_{xx}, & t > 0 \\ u(x, 0) = g(x), & 0 \end{cases}$$

Solve 
$$\begin{cases} u_t = u_{xx} \\ u(0, t) = g(t), \\ u_x(0, t) = 0 \end{cases}$$

$$u = \sum_{j=0}^{\infty} g_j(t) x^j$$

$$g_0 = g, g_1 = 0, g_j'(t) = (j+2)(j+1)g_{j+2} \Rightarrow g_{j+2} = \frac{1}{(j+1)(j+2)} g_j'(t)$$

$$u(x, t) = \sum_{k=0}^{\infty} \frac{g^{(k)}(t)}{(2k)!} x^{2k}$$

Now we choose

$$g(t) = \begin{cases} \exp[-t^{-\alpha}], & t > 0 \\ 0, & t \leq 0, \end{cases} \quad \alpha > 1$$

Then  $\exists \theta = \theta(\alpha) > 0$  s.t.

$$|g^{(k)}(t)| \leq \frac{k!}{(\theta t)^k} \exp[-\frac{1}{2}t^{-\alpha}]$$

Since  $k! / (2k)! < 1/k!$

$$\sum_{k=0}^{\pm\infty} \left| \frac{g^{(k)}(t)}{(2k)!} x^{2k} \right| \leq \sum_{k=0}^{\pm\infty} \frac{|x|^{2k}}{k! (\theta t)^k} \exp \left[ -\frac{1}{2} t^{-\alpha} \right]$$

$$= \exp \left[ \frac{1}{t} \left( \frac{|x|^2}{\theta} - \frac{1}{2} t^{-\alpha} \right) \right]$$

$$\begin{cases} u_t = \Delta u \\ u(x, 0) = 0 \end{cases} \quad u \in C^\infty(\mathbb{R}^{n+1})$$

$u$  is not bounded by  $e^{a|x|^2}$  (for  $t$  small)

Inhomogeneous heat equation

$$\begin{cases} u_t - \Delta u = f \\ u = 0 \end{cases}$$

Duhamel's Principle: We solve

$$\begin{cases} u_t(\cdot, s) - \Delta u(\cdot, s) = 0, & t > s \\ u(\cdot, s) = f(\cdot, s) & t = s \end{cases}$$

Then

$$u(x, t) = \int_0^t u(x, t; s) ds$$

In fact, formally,  $u_t = u(x, t; t) + \int_0^t \cancel{u(x, t; s)}$   
 $= f(x, t) - \Delta u$

$$u(x, t) = \int_0^t \frac{1}{(4\pi(t-s))^{n/2}} \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4(t-s)}} f(y, s) dy ds$$

Theorem:  $f \in C_1^2(\mathbb{R}^n \times [0, \infty))$  and  $f$  has compact support (32)

Then

$$\begin{cases} u_t - \Delta u = f \\ u(x, 0) = 0 \end{cases}$$

Proof:  $u(x, t) = \int_0^t \int_{\mathbb{R}^n} \Phi(y, s) f(x-y, t-s) dy ds$

$$u_t = \int_0^t \int_{\mathbb{R}^n} \Phi(y, s) f_t(x-y, t-s) dy ds + \int_{\mathbb{R}^n} \Phi(y, t) f(x-y, 0) dy$$

$$\Delta u = \int_0^t \int_{\mathbb{R}^n} \Phi(y, s) \Delta_x f(x-y, t-s) dy ds$$

$$u_t - \Delta u = \int_0^t \int_{\mathbb{R}^n} \Phi(y, s) [\partial_t - \Delta_x] f(x-y, t-s) dy ds$$

$$+ \int_{\mathbb{R}^n} \Phi(y, t) f(x-y, 0) dy$$

$$= \int_{\varepsilon}^t \int_{\mathbb{R}^n} \Phi(y, s) [(\partial_s - \Delta_y)] f(x-y, t-s) dy ds$$

$$+ \int_0^{\varepsilon} \int_{\mathbb{R}^n} \Phi(y, s) [ -\partial_s - \Delta_y ] f(x-y, t-s) dy ds$$

$$+ \int_{\mathbb{R}^n} \Phi(y, t) f(x-y, 0) dy$$

$$= I_{\varepsilon} + J_{\varepsilon} + K$$

$$|J_{\varepsilon}| \leq (\|f_t\|_{\infty} + \|\Delta^2 f\|_{\infty}) \int_0^{\varepsilon} \int_{\mathbb{R}^n} \Phi(y, s) dy ds \leq \varepsilon C$$

$$I_{\varepsilon} = \int_{\varepsilon}^t \int_{\mathbb{R}^n} [(\partial_s - \Delta_y) \Phi(y, s)] f(x-y, t-s) dy ds$$

$$+ \int_{\mathbb{R}^n} \Phi(y, \varepsilon) f(x-y, t-\varepsilon) dy - \int_{\mathbb{R}^n} \Phi(y, t) f(x-y, 0) dy$$



$$u_t - \Delta u = \lim_{\varepsilon \rightarrow 0} \int \Phi(y, \varepsilon) f(x-y, t-\varepsilon) dy$$

$$= f(x, t). \quad (\text{as } \varepsilon \rightarrow 0)$$

#

Summary

1)  $f \in C^2_{\perp}$ , compact support  $\Rightarrow$   
 $g \in C^0$

$$u(x, t) = \int \Phi(x-y, t) g(y) dy + \int_0^t \int \Phi(x-y, t-s) f(y, s) dy ds$$

Solves

$$\begin{cases} u_t - \Delta u = f \\ u = g \end{cases}$$

2) This sol'n is unique if

$$|u(x, t)| \leq e^{a|x|^2}$$

2) Maximum Principle works only for parabolic boundary

Finally, we discuss the local smoothness of heat equation.

Fix  $(x_0, t_0)$ .

parabolic cylinder



$$Q_r(x_0, t_0) = B_r(x_0) \times (t_0 - r^2, t_0)$$

$$Q_r \neq Q_r(0, 0)$$

Thm. Suppose  $u \in C^2_1(Q_r)$ ,  $u_t - \Delta u = 0$  in  $Q_r$

Then  $u \in C^\infty(Q_r)$  and

$$\max_{Q_{\frac{r}{2}}} |D_x^k D_t^l u| \leq \frac{C_{k,l}}{r^{k+2l+n+2}} \|u\|_{L^1(Q_r)}$$

Observe dimension balance

Proof: 1) First assume  $u \in C^\infty(Q_r)$ .

Let  $\zeta(x,t)$  be a cut-off for

$$\zeta \in C^\infty_{x,t}, \zeta|_{Q_{3/4}} = 1, \zeta|_{Q_{7/8}^c} = 0$$

Let  $v = \zeta u$

$$v_t - \Delta v = (\zeta_t - \Delta \zeta)u - 2\nabla \zeta \cdot \nabla u = f$$

$$v = \tilde{v}(x,t) = \int_{-1}^t \int_{\mathbb{R}^n} \Phi(x-y, t-s) f(y,s) dy ds$$

integrate by parts

$u = v$  in  $Q_{3/4}$

In  $Q_{1/2}$   $D_x^k D_t^l u = D_x^k D_t^l \tilde{v} = \int_{-1}^t \int_{\mathbb{R}^n} D_x^k D_t^l [\underbrace{\zeta_t - \Delta \zeta}_{\text{smooth}} + 2\nabla \zeta \cdot \nabla u] f$

$$|D_x^k D_t^l u(x,t)| \leq \int_{-1}^t \int_{B_1} |D_x^k D_t^l (\Phi(x-y, t-s)) (\zeta_t - \Delta \zeta)| |u(y,s)| dy ds$$
$$\leq \int_{-1}^t \int_{B_1} C_{k,l} |u(y,s)| dy ds = C_{k,l} \|u\|_{L^1(Q_r)}$$

2) The general case

$u^\epsilon = u * \eta$ ,  $\eta_\epsilon$  mollifier in  $x,t$

$$u^\epsilon(x,t) = \iint_{Q_1} K(x,t; y,s) u^\epsilon(y,s) dy ds \quad \epsilon \rightarrow 0$$