

Math 412, lecture 14

Last time: Rigid body motion

Today: Hamiltonian mechanics

- (1) convexity & Legendre transformations
- (2) Hamiltonians & Hamilton's equations

Physics $H = pr - L$ what does this mean?

(1) Convexity

Fix a f.d. real vector space V

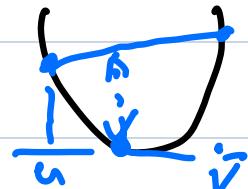
Def: A subset $C \subset V$ is **convex** if $\forall u, v \in C, t \in [0, 1]$
 $(1-t)u + tv \in C$.

A function $f: C \rightarrow \mathbb{R}$ defined on convex set is
convex if $\forall u, v \in C, t \in [0, 1]$

$$f((1-t)u + tv) \leq (1-t)f(u) + t f(v)$$

strictly convex if $f((1-t)u + tv) < (1-t)f(u) + t f(v)$ when $u \neq v, t \neq 0, 1$.

Geometrically, f convex



secant line is
above graph

Ex: f is convex iff the epigraph $\{(u, a) \mid f(u) \leq a\}$ is convex in $V \times \mathbb{R}$.

Remark: Allow f to take value ∞ , then might as well have them defined on V , check: $\{u \mid f(u) < \infty\}$ is convex

Ex: Epigraph is closed iff f is lower semicontinuous
 $\forall v \forall \varepsilon > 0 \exists u \text{ nbd of } v \text{ s.t. if } u \in U,$
 $f(u) \geq f(v) - \varepsilon$

Ex: let $f: I \rightarrow \mathbb{R}$ be convex, $I \subset \mathbb{R}$ interval
(1) slope of secant line from $(u, f(u))$ to $(v, f(v))$,
is monotone in u, v .
(2) f has left & right derivatives at every point
& if $u \in V$

$$f'_L(u) \leq f'_R(u) \leq f'_L(v) \leq f'_R(v)$$

(3) At most countably many points where f is non-diff, which are jump discontinuities, f' is monotone.

Cor: we can parametrize points u by slopes $f'(u)$
(if f is strictly convex, f' strictly increasing)

The mapping $u \mapsto f'(u)$ is a map $V \rightarrow V^*$

$$df_u$$

Def: The Fenchel conjugate or convex conjugate or Legendre transform of $f: V \rightarrow \mathbb{R} \cup \{\infty\}$ is the function f^* defined on V^* by

$$f^*(p) = \sup_{v \in V} [\langle p, v \rangle - f(v)]$$

(f is "proper" if $f(v) \neq \infty$ for some v)

Example: (1) $f(x) = \langle m, x \rangle + b$

$$\sup_{x \in V} (\langle p-m, x \rangle - b) = \begin{cases} -b & p=m \\ \infty & p \neq m \end{cases}$$

(2) For $1 < r < \infty$, if $f(x) = \frac{1}{r} |x|^r$, $f^*(p) = \frac{1}{r^*} |p|^{r^*}$

$$\frac{1}{r} + \frac{1}{r^*} = 1.$$

(3) Let $g: V \rightarrow V^*$ be symmetric, pos. def.

and let $f(x) = \frac{1}{2} \langle gx, x \rangle + \langle m, x \rangle + b$

then

$$f^*(p) = \frac{1}{2} \langle p-m, g^{-1}(p-m) \rangle - b$$

Pf: Given p let $h = g^{-1}(p-m) \in V$.

$$\text{Then } \langle p, x \rangle - f(x) = \langle p-m, x \rangle - \frac{1}{2} \langle gx, x \rangle - b$$

completing square

$$= -\frac{1}{2} \langle g(x-h), (x-h) \rangle + \frac{1}{2} \langle gh, h \rangle - b$$

maximized if $x=h$ whence claim

$$\langle gh, h \rangle = \langle p-m, g^{-1}(p-m) \rangle$$

Observation: Fenchel's inequality, $f(v) + f^*(p) \leq \langle p, v \rangle$

Lemma: let f be any function

(1) f^* is lsc & convex

(2) $\{(v, a) : f^{**}(v) \leq a\} = \overline{\text{Conv}}(\{(v, a) : f(v) \leq a\})$
closed convex hull

Theorem: (Fenchel - Moreau) i.: f is proper, convex,
'closed' (lsc.) then $f^{***} = f$.

Key point: K convex, closed, $x \notin K$ $\exists p^* \in p(K) \subset K$
 $p^*(x) < p(x)$

If f is strictly convex, diff, choose $u \in \text{dom}(f)$
consider $f^*(df_u)$: function $v \mapsto \langle df_u, v \rangle - f(v)$
has its unique critical point at u

$$d[\langle df_u, v \rangle - f(v)] = df_u - df_v$$

$$\Rightarrow f^*(df_u) = \langle df_u, u \rangle - f(u)$$

Phase space & the Hamiltonian

Def: Phase space or the cotangent bundle is
the space

$$T^*X = \{(x, p) \mid \begin{array}{l} x \in X \\ p \in (T_x X)^* \end{array}\}$$

Now fix a Lagrangian $L: TX \times E^1 \rightarrow \mathbb{R}$

Assume $v \mapsto L(x, v; t)$ is strictly convex on $T_x E$ for x, t fixed

Def The **Hamiltonian** of the system is the function

$$H: T^* E \times E^1 \rightarrow \mathbb{R}$$

s.t. for fixed x, t , $p \mapsto H(x, p; t)$ is
the convex conjugate of $v \mapsto L(x, v; t)$

Example: $L(x, v) = \frac{1}{2} \langle M(x)v, v \rangle - U(x)$

then $H(x, p) = \frac{1}{2} \langle M^{-1}(x)p, p \rangle + U(x)$

(if only dependence on v is in quadratic kinetic term, then $H = T + U$)

To think on $\frac{1}{2} \langle Mv, v \rangle$, $\frac{1}{2} \langle M^{-1}p, p \rangle$ as "same function", identify state of motion (x, v, t) with phase point (x, p, t)

where $p = \frac{\partial L}{\partial v} \Big|_{(x, v, t)}$

Observes This depends on choice of L

Aside: in optimization, how functions f on V (say f is convex) want to minimize f .

Paradigm: start at point $x_0 \in V$ step down against gradient ∇f , in this time have ODE $\dot{x} = -\nabla f(x)$ then $\frac{d}{dt} f(x) = -\langle \nabla f, \dot{x} \rangle \leq 0$

But the derivative of f is $df_x \in V^*$
need a rule mapping $V^* \rightarrow V$ (i.e. a metric / inner product)

Still for quadratic Lagrangian,

$$L = \frac{1}{2} \langle M v, v \rangle - U(r)$$

$$p = \frac{\partial L}{\partial v} = Mv, \quad v = M^{-1}p$$

Then $H = p v - L$

§2 The Hamiltonian flow

The physical path $\gamma(t) = x(t)$ is a path of states of motion $(x(t), v(t); t)$ $v(t) = \dot{x}(t)$

\Rightarrow path in phase space (x, p, t)

To start with $\dot{x} = v = \frac{\partial H}{\partial p}$

what about \dot{p} ? $\dot{p} = \frac{\partial L}{\partial x}$ (Euler-Lagrange)

$L = pV - H$ so maybe $\frac{\partial L}{\partial x} = -\frac{\partial H}{\partial x}$?

This is nonsense: $\frac{\partial L}{\partial x} = \left(\frac{\partial L}{\partial x}\right)_V$.
natural $\frac{\partial H}{\partial x}$ is $\left(\frac{\partial H}{\partial x}\right)_p$.

Prove of this nonsense

think of H as function on TX

$$\begin{aligned} dH &= d(pV - L) = \langle dp, V \rangle + \langle p, dV \rangle \\ &\quad - \underbrace{\frac{\partial L}{\partial v} dv}_{-pdv} = \frac{\partial L}{\partial x} dx - \frac{\partial L}{\partial t} dt \\ &= \langle dp, V \rangle - \frac{\partial L}{\partial x} dx - \frac{\partial L}{\partial t} dt \end{aligned}$$

$$\Rightarrow V = \frac{\partial H}{\partial p}, \quad -\left(\frac{\partial L}{\partial x}\right)_V = \left(\frac{\partial H}{\partial x}\right)_p, \quad \frac{\partial H}{\partial t} = -\frac{\partial L}{\partial t}$$

Now push everything to T^*X

$$\Rightarrow \begin{cases} \dot{x} = \frac{\partial H}{\partial p} \\ \dot{p} = -\frac{\partial H}{\partial x} \end{cases} \quad \Leftarrow \text{Hamilton's equations}$$

$$\Rightarrow z = (x, p) \quad \dot{z} = X_H \quad X_H = \text{vector field on } T^*X$$

If have $dF = a dx + b dy$

if $V = \begin{pmatrix} V_x \\ 0 \end{pmatrix}$ then $\langle dF, V \rangle = a \cdot V_x$

$$\text{ie. } a = \left(\frac{\partial F}{\partial x}\right)_y$$

UBC Library has online copy of Landau & Lifshitz.