

I. Finite systems (Ruelle, pp. 3-4)

Recall defn. of entropy of probability vector: $\bar{p} = (p_1, \dots, p_n)$:

$$H(\bar{p}) = - \sum_i p_i \log p_i$$

(p_i is probability of “micro-state” i .)

Let $u_i = U(i)$ be energy contribution from state i .

Let μ_U be the probability vector (p_1, \dots, p_n) defined by

$$p_i = \frac{e^{-u_i}}{Z(U)}$$

where Z is the normalization factor:

$$Z(U) = \sum_j e^{-u_j}.$$

Prop (from class): Let $E = E_{\mu_U}(U)$. Then

$$\max_{\bar{p}: E_{\bar{p}}U=E} H(\bar{p}) = \log Z(U) + E$$

and achieved uniquely by μ_U .

Interpretation: Given energy function U and expected value of energy, you get a well-defined “most likely state” μ_U

Prop (Ruelle, bottom of p. 3):

$$\max_{\bar{p}} H(\bar{p}) - E_{\bar{p}}(U) = \log Z(U)$$

and achieved uniquely by μ_U .

Proof of Ruelle Prop:

$$H(\bar{p}) - E_{\bar{p}}(U) = \sum_i p_i (-\log(p_i) - u_i)$$

$$= \sum_i p_i \log \frac{e^{-u_i}}{p_i} \leq \log \sum_i p_i \frac{e^{-u_i}}{p_i} = \log Z(U)$$

with equality iff $\frac{e^{-u_i}}{p_i}$ is constant ($Z(U)$). by Jensen.

Note: class prop follows since

$$\max_{\bar{p}: E_{\bar{p}}U=E} H(\bar{p}) - E \leq \max_{\bar{p}} H(\bar{p}) - E_{\bar{p}}(U)$$

and LHS includes μ_U .

Observe:

max is a non-probabilistic quantity.

achieved uniquely by an explicit probability “measure”, a “equilibrium state”

maximizing measure has a certain form, a “Gibbs state”

So, Gibbs states = Equilibrium states

II. Variational Principle

Theorem 1 (Ruelle, p. 6):

Let M be a compact metric space.

Let T be a continuous Z^d -action on M and $f : M \rightarrow R$ a continuous function.

Let $\mathcal{M}(T)$ be the set of all Borel probability measures invariant under T .

For $\mu \in \mathcal{M}(T)$, let $h_\mu(T)$ denote the measure-theoretic entropy of T w.r.t. μ .

Let $P(T, f)$ denote the pressure (of f, T). Then

$$P(T, f) = \sup_{\mu \in \mathcal{M}(T)} h_\mu(T) + \int f d\mu$$

Dictionary of notation:

These notes	Ruelle
M	Ω
T	τ
f	A
μ	σ
$P(T, f)$	$P(A)$
α	\mathfrak{A}
$x, y \in M$	$\xi, \eta \in \Omega$
$\mathbf{m} \in Z^d$	$x \in Z^d$

Note: Under certain conditions, sup is *achieved* and under stronger conditions achieved *uniquely*.

Will now define the terms, with examples, in Theorem 1:

Compact metric space: for Ruelle, “*metrizable*” because particular quantities do not depend upon specific choice of metric.

continuous Z^d -action

Defn: group homomorphism: $Z^d \rightarrow \text{Homeo}(M, M)$

Action generated by pairwise commuting homeos T_1, \dots, T_d of M and for $(m_1, \dots, m_d) \in Z^d$,

$$T^{(m_1, \dots, m_d)}(x) = T_1^{m_1} \circ \dots \circ T_d^{m_d}(x)$$

$d = 1$:

$$T^m(x) = T_1^m(x)$$

Main Example: *Full Z -shift*:

$M = F^Z$ with product topology (configurations on Z with finite alphabet F).

Metric: $d(x, y) = 2^{-k}$, where x, y agree on $2k + 1$ interval centered at origin but not larger k .

$T_1 =$ left-shift map;

Special Class: *Z-shift of finite type (SFT)*:

“Forbid *finitely* many configurations on finite intervals”

Examples:

Golden Mean

$F = \{0, 1\}$: forbid 11

RLL(1,2)

$F = \{0, 1\}$: forbid 11 and 000

Very special class: *Topological Markov Chain (n.n. Z-SFT)*:

Defn: Let C be a square 0-1 (transition) matrix (say $m \times m$).

Let $F = \{0, \dots, m - 1\}$.

Let $M_C = \{x \in F^{\mathbb{Z}} : C_{x_i, x_{i+1}} = 1 \text{ for all } i \in \mathbb{Z}\}$

“allowed” viewpoint

T_C : left shift on M_C

For golden mean:

$$C = \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix}$$

$d = 2$:

Main Example: *Full Z^2 -shift*:

$M = F^{Z^2}$ with product topology (configurations on 2-dimensional integer lattice with finite alphabet F).

Metric: $d(x, y) = 2^{-k}$, where x, y agree on a $(2k + 1) \times (2k + 1)$ square centered at origin but no larger k .

$T_1 =$ left-horizontal shift; $T_2 =$ down-vertical shift,

Special Example: Z^2 -shift of finite type (SFT):

“Forbid finitely many configurations on finite shapes”

Given finite sets $\Delta_1, \dots, \Delta_n \subset Z^2$ and $u_1 \in F^{\Delta_1}, \dots, u_n \in F^{\Delta_n}$,

$$M = \{x \in F^{Z^2} : \forall \mathbf{m} \in Z^2, x_{\Delta_i + \mathbf{m}} \neq u_i, i = 1, \dots, n\}$$

Note: *translation-invariant* condition

Ruelle: defines SFT by “allowed” configs (Ω in mid-page 7):

a finite set $\Delta \subset Z^2$ (think rectangle) and set $G \subset F^\Delta$:

$$M = \{x \in F^{Z^2} : \forall \mathbf{m} \in Z^2, x_{\Delta + \mathbf{m}} \in G\}$$

Examples:

Hard square

$F = \{0, 1\}$: forbid 11 horizontally and vertically

$$RLL(1, 2)^{\otimes 2}$$

Dominos (Dimers)

$F = \{L, R, T, B\}$: forbid horizontal configs LL, LT, LB, RR, TR, BR, and vertical configs.

Monomer-Dimers

Z^2 -TMC (n.n. Z^2 -SFT): horizontal and vertical transition matrices

(Topological) Entropy:

Defn: For finite open cover α of M ,

$$N(\alpha) = \text{minimum size of subcover of } \alpha$$

$$H(\alpha) = \log N(\alpha).$$

Defn: for finite open covers α, β of M

$$\alpha \vee \beta = \{A_i \cap B_j : \text{nonempty}\}$$

$$\text{For } \mathbf{m} \in \mathbb{Z}^d, T^{-\mathbf{m}}(\alpha) = \{T^{-\mathbf{m}}(A_i)\}$$

For set $\Lambda \subset \mathbb{Z}^d$:

$$\alpha_\Lambda = \bigvee_{\mathbf{m} \in \Lambda} T^{-\mathbf{m}}(\alpha).$$

Consider d -dimensional prisms $\Lambda = \Lambda(a_1, \dots, a_d)$,

$$h(T, \alpha) = \lim_{a_1, \dots, a_d \rightarrow \infty} (1/|\Lambda|) \log N(\alpha_\Lambda)$$

Defn:

$$h(T) = \sup_{\alpha} h(T, \alpha)$$

Theorem: If α is a top. generator (i.e., distinguishes points), then $h(T) = h(T, \alpha)$.

Here, “distinguishes points” means: letting $\alpha(x) = \bigcup_{i: x \in A_i} A_i$, if $x, y \in M$ and $x \neq y$, then for some $u \in \mathbb{Z}^d$, $\alpha(T^u(x)) \cap \alpha(T^u(y)) = \emptyset$.

For full shift and SFT's, the standard cover:

$\alpha = \{\{x \in M : x_0 = a\} : a \in F\}$ is a topological generator.

So, $h(T)$ is a growth rate of counts.

$d = 1$:

$$h(T) = \lim_{n \rightarrow \infty} (1/n) \log(\# \text{ allowed } n\text{-sequences})$$

Proposition: For a \mathbb{Z} -TMC M_C ,

$$h(T_C) = \log \lambda_C$$

where λ_C is the spectral radius of C , i.e.,

$$\lambda_C = \max\{|\lambda| : \lambda \text{ eigenvalues of } C\}$$

Proof:

$$h(T_C) = \lim_{n \rightarrow \infty} \left(\frac{1}{n}\right) \log \mathbf{1}(C)^{n-1} \mathbf{1}$$

Golden Mean shift: $h(T) = \log(\text{golden mean})$

$$C = \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix}$$

Can deduce formula for top. entropy of any Z -SFT,
 $d = 2$:

$$h(T) = \lim_{n \rightarrow \infty} (1/n^2) \log(\#n \times n \text{ allowed arrays})$$

Hard square: ???

$$\text{Dominos (Dimers): } (1/4) \int_0^1 \int_0^1 (4 - 2 \cos(2\pi s) - 2 \cos(2\pi t)) ds dt$$

— Monomer-Dimer: ???

Pressure:

Defn given in Ruelle, p. 5:

Same as top, entropy except:

For $\Lambda = \Lambda(a_1, \dots, a_d)$, replace $N(\alpha_\Lambda)$ by:

$$Z_\Lambda(f, \alpha) = \min_{\text{subcover } \beta \text{ of } \alpha_\Lambda} \sum_j \exp\left(\sup_{x \in B_j} \sum_{u \in \Lambda} f(T^u(x))\right)$$

$$\beta = \{B_1, B_2, \dots, B_{n(\beta)}\}$$

So,

$$P(T, f, \alpha) = \lim_{a_1, \dots, a_d \rightarrow \infty} (1/|\Lambda|) \log Z_\Lambda(f, \alpha)$$

and

$$P(T, f) = \sup_{\alpha} P(T, f, \alpha)$$

Theorem: If α is a topological generator, then $P(T, f) = P(T, f, \alpha)$.

Note:

$$Z_\Lambda(0, \alpha) = N(\alpha_\Lambda).$$

$$P(T, 0) = h(T).$$

Examples:

$d = 1$:

T_C is TMC:

$$f(x) = f(x_0x_1)$$

$$P(T_C, f) = \lim_{n \rightarrow \infty} (1/n) \log \left(\sum_{x_0 \dots x_n} \exp(f(x_0x_1) + \dots + f(x_{n-1}x_n)) \right)$$

Note: No min and No sup.

Prop:

$$P(T_C, f) = \log \lambda(C_f)$$

where

$$(C_f)_{ij} = C_{ij} e^{f(ij)}.$$

Proof:

$$P(T_C, f) = \lim_{n \rightarrow \infty} \left(\frac{1}{n} \right) \log \mathbf{1}(C_f)^{n-1} \mathbf{1}$$

$d = 2$:

1. Hard square with activity.

$T =$ Hard square SFT

Let $c \in \mathbb{R}$ and define

$$f_c(x) = c \text{ if } x_0 = 1$$

$$f_c(x) = 0 \text{ if } x_0 = 0.$$

$P(T, f_c)$ = growth rate of number of allowed arrays, with 1's weight by e^c and 0's weighted by 1.

$a = e^c$; activity level

No exact known formula for $P(T, f_c)$ known.

2. Ising model

T = full shift on $F = \{\pm 1\}$

f : Ising model

Given constants β, J, H ,

$$f(x) = \beta(Bx_{0,0} + J(x_{0,0}x_{1,0} + x_{0,0}x_{0,1}))$$

β : inverse temperature

J : interaction strength

B : external magnetic field strength

$P(T, f)$: growth rate of number of allowed arrays, weighted by e^f , which incorporates interactions on adjacent sites (horizontal and vertical) and magnetic field (on individual sites) .

Onsager: exact solution for $P(T, f)$, when $B = 0$.

Measure-theoretic entropy

Let T be an MPT Z^d -action on probability space (X, \mathcal{A}, μ) .

Defn: For finite, measurable *partition* α ,

$$H_\mu(\alpha) = - \sum_i \mu(A_i) \log \mu(A_i)$$

where $\alpha = \{A_i\}$.

For finite set $\Lambda \subset Z^d$:

$$\alpha_\Lambda = \bigvee_{\mathbf{m} \in \Lambda} T^{-\mathbf{m}}(\alpha).$$

Consider d -dimensional prisms $\Lambda = \Lambda(a_1, \dots, a_d)$,

$$h_\mu(T, \alpha) = \lim_{a_1, \dots, a_d \rightarrow \infty} (1/|\Lambda|) H_\mu(\alpha_\Lambda)$$

Defn:

$$h_\mu(T) = \sup_{\alpha} h_\mu(T, \alpha)$$

Theorem: If α is a meas.-theo. generator (i.e., $\alpha_{Z^d} = \mathcal{A}$ a.e.), then $h_\mu(T) = h_\mu(T, \alpha)$.

$d = 1$:

X is a *stationary* process with law μ and $T =$ left-shift, then

$$h_\mu(T) = h(X) = \lim_{n \rightarrow \infty} (1/n) H(X_1, \dots, X_n)$$

. where $H(X_1, \dots, X_n)$ is the entropy of (X_1, \dots, X_n) as a random vector.

Examples:

$\mu = \text{iid}(\bar{p})$:

$$h_\mu(T) = H(\bar{p})$$

μ : stationary (first-order) Markov with probability transition matrix P with stationary vector π :

$$h_\mu(T) = - \sum_{ij} \pi_i P_{ij} \log P_{ij}$$

$d = 2$:

X is a stationary Z^2 -process with law μ and $T^{(m,n)}$: shift by translation (m, n) . Then

$$h_\mu(T) = h(X) = \lim_{n \rightarrow \infty} (1/n^2) H(X_{i,j} : 1 \leq i, j \leq n).$$

where H is the entropy of the random vector (array):

$$X_{i,j} : 1 \leq i, j \leq n$$

Examples:

1. $\mu = \text{iid}(p)$:

$$h_\mu(T) = H(p)$$

2. Markov chains replaced by Gibbs measures/Markov random fields.

Few explicit results.

Back to Variational Principle:

$$P(T, f) = \sup_{\mathcal{M}(T)} h_\mu(T) + \int f d\mu$$

Defn: An *equilibrium state* for T, f is a measure $\mu \in \mathcal{M}(T)$ which achieves $P(T, f)$.

Let $I_{T,f}$ denote the set of equilibrium states (which can be empty).

Defn: T is *expansive* if there exists $\delta > 0$ s.t. $\forall x \neq y \in M, \exists \mathbf{m} \in Z^d$ s.t. $\text{dist}(T^{\mathbf{m}}x, T^{\mathbf{m}}y) > \delta$.

Fact: Any Z^d -SFT is expansive.

Theorem: If T is expansive, then for every continuous f , $I_{T,f} \neq \emptyset$.

Proof uses upper semi-continuity of $\mu \mapsto h_\mu(T)$

Non-uniqueness corresponds to *phase transition*.

$d = 1$:

Special case:

Theorem (Variational Principle for irreducible TMC) Let T be TMC and $f(x) = f(x_0x_1)$. Let

$$(C_f)_{ij} = C_{ij}e^{f(ij)}$$

Then

$$P(T_C, f) = \log \lambda_{C_f} = \sup_{\mu \in \mathcal{M}} h_\mu(T_C) + \int f d\mu$$

and the sup is achieved uniquely by an explicitly describable Markov chain:

$$P_{ij} = C_{ij}e^{f(ij)} \frac{v_j}{\lambda_{C_f} v_i}$$

where v is a right eigenvector for matrix C_f and eigenvalue λ_{C_f} .

Example: Golden mean with $f = f_c$, ($a = e^c$).

$$C_f = \begin{bmatrix} a & 1 \\ a & 0 \end{bmatrix}$$

$$\lambda = \frac{a + \sqrt{a^2 + 4a}}{2}$$

$$v = \begin{bmatrix} \lambda \\ a \end{bmatrix}$$

No phase transition!

See lecture notes from Entropy class for proof in case $c = 0$.

$d = 2$:

1. Hard core with activity $a = e^c$: unique equilibrium state up to some critical threshold.

2. Ising model: unique equilibrium state up to some critical threshold in β , when $B = 0$.

Gibbs measures

Let $T : M \rightarrow M$ be a nearest neighbour Z^2 -SFT.

Let $C_1 = F$, the alphabet (a.k.a. configurations on single nodes)

Let C_2 be all *allowed* configurations on domino shapes (i.e., configurations on 1×2 and 2×1 rectangles).

Let $\Phi : C_1 \cup C_2 \rightarrow R$, (*nearest-neighbour interaction*).

A *translation invariant (stationary) nearest-neighbour Gibbs measure* on M is a T -invariant measure μ on M such that for all finite subsets $\Lambda \subset Z^d$ and a.e. $x \in M$,

$$\mu(x|_{\Lambda} \mid x|_{\Lambda^c}) \sim \left(\prod_{v \in \Lambda} \exp(\Phi(x_v)) \right) \left(\prod_{i=1}^d \prod_{\{v \in \Lambda, v+e_i \in \Lambda \cup \partial\Lambda\}} \exp(\Phi((x_v, x_{v+e_i}))) \right). \quad (1)$$

In particular, $\mu(x|_{\Lambda} \mid x|_{\Lambda^c}) = \mu(x \mid x|_{\partial\Lambda})$.

Let

$$f_{\Phi}(x) = \Phi(x_{(0,0)}) + \Phi(x_{(0,0)}, x_{(1,0)}) + \Phi(x_{(0,0)}, x_{(0,1)})$$

Theorem:

$\{ \text{Equilibrium states for } f_{\Phi} \} \subseteq \{ \text{translation invariant Gibbs states for } \Phi \}$

Assuming Condition D (a mixing condition) on the SFT M (Ruelle, p. 57), in particular, for the full shift,

$\{ \text{Equilibrium states for } f_{\Phi} \} = \{ \text{translation invariant Gibbs states for } \Phi \}$

There is a much more general version of this (see Ruelle, Theorem 3, p.8 and Theorem 4.2, p. 58):

1. Begin with an Interaction: function Φ on allowed configurations on finite sets (see Chapter 1)

2. Form the Energy function: f_Φ , a sum of interaction values
3. A Gibbs measure is a measure μ that satisfies: whenever $x, y \in M$ disagree at only finitely many sites, then

$$\mu(x|_\Lambda \mid x|_{\Lambda^c}) = \left[\sum_{y: y_{Z^d \setminus \Lambda} = x_{Z^d \setminus \Lambda}} \prod_{u \in Z^d} \exp(f_\Phi(T^u(y)) - f_\Phi(T^u(x))) \right]^{-1}$$

(in nearest neighbour special case above, this is equivalent to (1))

In Ruelle (pp, 7-8), there is no mention of interaction Φ . Gibbs measure is defined for any function f on M that has exponentially decreasing dependence (equivalently Holder continuous). In Ruelle (chapters 3 and 4), $f = f_\Phi$ where Φ has satisfies a summability condition.

Equilibrium states and derivative of pressure

Let $C^\alpha(M, R)$ denote the set of Holder continuous functions, with exponent α from M to R .

For a topologically mixing Z -SFT $T : M \rightarrow M$, and $f, g \in C^\alpha(M, R)$, if μ_f is the unique equilibrium state for T, f , then

$$\frac{d}{dt} P(f + tg) = \int g d\mu_f$$

Thus, a unique equilibrium state can be viewed as a derivative of the pressure map

$$P : C^\alpha \rightarrow R$$

Phase transitions correspond to discontinuities in derivative of pressure (as well as non-uniqueness of equilibrium states).