

[ **Last lecture:** ... Hopf bifurcation ... ]

**Example 3.C**, continued. *Brusselator* chemical reaction model

$$\begin{aligned} \dot{x}_1 &= \gamma - x_1 - \alpha x_1 + x_1^2 x_2 \\ \dot{x}_2 &= \alpha x_1 - x_1^2 x_2 \end{aligned} \tag{3.C.0}$$

We suppose  $\gamma > 0$  is fixed, and treat  $\alpha$  as the bifurcation parameter.

Family of equilibria can be found explicitly (**Exercise**),

$$x_1 = p_1^0(\alpha) = \gamma, \quad x_2 = p_2^0(\alpha) = \frac{\alpha}{\gamma}.$$

The first coordinate change

$$x_1 = \gamma + u_1, \quad x_2 = \frac{\alpha}{\gamma} + u_2 \tag{I}$$

transforms (3.C.0) into

$$\begin{pmatrix} \dot{u}_1 \\ \dot{u}_2 \end{pmatrix} = \underbrace{\begin{pmatrix} \alpha - 1 & \gamma^2 \\ -\alpha & -\gamma^2 \end{pmatrix}}_{A(\alpha)} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} + \begin{pmatrix} 1 \\ -1 \end{pmatrix} \left( \frac{\alpha}{\gamma} u_1^2 + 2\gamma u_1 u_2 + u_1^2 u_2 \right) \tag{1}$$

or

$$\dot{u} = A(\alpha) u + \hat{f}^{(2)}(u, \alpha) + \hat{f}^{(3)}(u, \alpha). \tag{3.C.1}$$

By linear stability analysis, it is easy to verify that (3.C.0) satisfies (H.0.i) and (H.0.ii) with

$$\alpha_0 = 1 + \gamma^2, \quad p_0^0 = \left( \gamma, \frac{1 + \gamma^2}{\gamma} \right) \quad \omega_0 = \gamma,$$

where  $\gamma > 0$ .

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It is also easy to verify (H.1):

Now, to verify (H.2): put  $\alpha = \alpha_0 = 1 + \gamma^2$  in (3.C.1),

Solve (carefully!)  $A_0 q = i\gamma q$ ,  $A_0^T p = -i\gamma p$ ,  $\langle p, q \rangle = 1$ :

Put  $u = z_1 q + \bar{z}_1 \bar{q}$ , i.e.

in (3.C.1), take the inner product with  $p$  to get

Expand in powers of  $z_1$  and  $\bar{z}_1$

identify the three important coefficients

and find the cubic normal form coefficient

By Theorem 3.6, the dynamics of (3.C.0) at  $(p_0^0, \alpha_0)$  can be deduced from the normal form system in polar coordinates

so there are bifurcating stable limit cycles for (3.C.0) if  $\alpha > 1 + \gamma^2$ , at least for  $\alpha$  sufficiently near  $1 + \gamma^2$ .

A schematic (**AUTO**-style) one-parameter *local* branching diagram (showing max and min values of  $x_2^0(t)$  on the limit cycle that exists for  $\alpha > 1 + \gamma^2$ )

A two-parameter bifurcation diagram in the  $(\alpha, \gamma)$ -plane with *local* phase portraits for (3.C.0), and for  $\alpha$  *near*  $1 + \gamma^2$ :

To study branches of limit cycles for  $\alpha$  farther from the bifurcation value, typically need numerical computation, e.g. **AUTO** or **MatCont**.

## Centre manifolds

Centre manifold theory is applied, to a vector field or a map, to locally “reduce” the dynamical system to a dynamical system in a lower state space dimension. The theory is easily adapted to families of vector fields or maps. With centre manifold theory, we can analyze local bifurcations (fold, Hopf, etc.) in an  $n$ -dimensional system, even if  $n$  is large.

We summarize the theory for a vector field

$$\dot{x} = f(x), \quad x \in \mathbb{R}^n \quad (3.5)$$

(the theory for a map is similar). At a nonhyperbolic equilibrium  $p^0$ , the  $n \times n$  matrix  $A = f_x(p^0)$  has a centre subspace  $T^c$  of dimension  $n_0 > 0$ , and there is a smooth, locally invariant **local centre manifold**  $W_{loc}^c(p^0)$  with the same dimension  $n_0$ . More precisely, there is the following theorem:

**Theorem 3.7. (Local Centre Manifold)** *If  $f : \mathbb{R}^n \rightarrow \mathbb{R}^n$  is  $C^p$  ( $p \geq 1$ ) in an open set containing  $p^0$ , if  $f(p^0) = 0$ , and if  $A = f_x(p^0)$  has  $n_0 > 0$  eigenvalues  $\lambda_j$ , counting multiplicities, with  $\operatorname{Re} \lambda_j = 0$ , then there exists a  $C^p$  submanifold  $W_{loc}^c(p^0)$  in  $\mathbb{R}^n$ , of dimension  $n_0$ , that is locally invariant for (3.5), contains  $p^0$ , and is tangent to the centre subspace at  $p^0$ . Moreover, there is an open neighbourhood  $U$  of  $p^0$  in  $\mathbb{R}^n$ , such that if a solution  $x(t)$  for (3.5) satisfies  $x(t) \in U$  for all  $t \geq 0$  [for all  $t \leq 0$ ], then  $x(t) \rightarrow W_{loc}^c(p^0)$  as  $t \rightarrow +\infty$  [as  $t \rightarrow -\infty$ ].*



We develop a method for using the centre manifold theorem (the Reduction Principle) to determine dynamics restricted to a local centre manifold. Suppose the centre, stable and unstable subspaces of the linearization  $A = f_x(p^0)$  have dimensions

$$\dim T^c = n_0, \quad \dim T^s = n_-, \quad \dim T^u = n_+,$$

respectively, with  $0 < n_0 < n$ ,  $0 < n_{\pm} \triangleq n_- + n_+ < n$ ,  $n_0 + n_{\pm} = n$ .

Let

$$T^{su} = T^s \oplus T^u$$

be the **stable-unstable** subspace. Then we have

$$\mathbb{R}^n = T^c \oplus T^{su}, \quad \dim T^c = n_0, \quad \dim T^{su} = n_{\pm},$$

and there exists a corresponding coordinate shift (I) and linear change of coordinates (II), from  $x \in \mathbb{R}^n$ , to  $(u, v) \in \mathbb{R}^{n_0} \times \mathbb{R}^{n_{\pm}}$ , so that in the new coordinates  $(u, v)$  the equilibrium is the origin  $(0, 0)$  and the linearization of the vector field at the equilibrium has a block-diagonal form (for example, real normal form),

$$\begin{pmatrix} \dot{u} \\ \dot{v} \end{pmatrix} = \begin{pmatrix} B & O \\ O & C \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} + \begin{pmatrix} g(u, v) \\ h(u, v) \end{pmatrix}, \quad (u, v) \in \mathbb{R}^{n_0} \times \mathbb{R}^{n_{\pm}}, \quad (3.6)$$

where

and both nonlinear functions

$$g : \mathbb{R}^{n_0} \times \mathbb{R}^{n_{\pm}} \rightarrow \mathbb{R}^{n_0}$$

$$h : \mathbb{R}^{n_0} \times \mathbb{R}^{n_{\pm}} \rightarrow \mathbb{R}^{n_{\pm}}$$

are locally defined at the origin  $(u, v) = (0, 0)$ , are  $C^p$  and  $O(\|(u, v)\|^2)$  (if  $p \geq 2$ ).

In these coordinates, the local centre manifold can be represented as the graph of a  $C^p$  function

where

is defined and smooth in an open neighbourhood of  $u = 0$  in  $\mathbb{R}^{n_0}$ , and  $V(u) = O(\|u\|^2)$  (i.e.  $V(0) = 0$  and  $V_u(0) = 0$ ). The dynamics restricted to  $W_{loc}^c(p^0)$  essentially determine the local dynamics of the full system: