ON THE TOPOLOGICAL BOUNDARY OF THE RANGE OF SUPER-BROWNIAN MOTION

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We show that if $\partial \mathcal{R}$ is the boundary of the range of super-Brownian motion and dim denotes Hausdorff dimension, then with probability one, for any open set $U, U \cap \partial \mathcal{R} \neq \emptyset$ implies

$$\dim(U \cap \partial \mathcal{R}) = \begin{cases} 4 - 2\sqrt{2} \approx 1.17 & \text{if } d = 2\\ \frac{9 - \sqrt{17}}{2} \approx 2.44 & \text{if } d = 3. \end{cases}$$

This improves recent results of the last two authors by working with the actual topological boundary, rather than the boundary of the zero set of the local time, and establishing a local result for the dimension.

1. Introduction. We consider a d-dimensional super-Brownian motion (SBM), $(X_t, t \geq 0)$, starting at X_0 under \mathbb{P}_{X_0} with $d \leq 3$. Here $X_0 \in M_F(\mathbb{R}^d)$, the space of finite measures on \mathbb{R}^d with the weak topology, X is a continuous $M_F(\mathbb{R}^d)$ -valued strong Markov process, and \mathbb{P}_{X_0} denotes any probability under which X is as above. We write $X_t(\phi)$ for the integral of ϕ with respect to X, and take our branching rate to be one, so that for any non-negative bounded Borel functions ϕ , f on \mathbb{R}^d ,

(1.1)
$$\mathbb{E}_{X_0}\left(\exp\left(-X_t(\phi) - \int_0^t X_s(f)ds\right)\right) = \exp(-X_0(V_t(\phi, f))).$$

Here $V_t(x) = V_t(\phi, f)(x)$ is the unique solution of the mild form of

(1.2)
$$\frac{\partial V}{\partial t} = \frac{\Delta V_t}{2} - \frac{V_t^2}{2} + f, \quad V_0 = \phi,$$

that is,

$$V_t = P_t(\phi) + \int_0^t P_s \left(f - \frac{V_{t-s}^2}{2} \right) ds.$$

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In the above (P_t) is the semigroup of standard d-dimensional Brownian motion. See Chapter II of [21] for the above and further properties. Note that X has an a.s. finite extinction time, and therefore we can define the so-called total occupation time measure of the super-Brownian motion as a finite measure,

$$I(A) = \int_0^\infty X_s(A)ds.$$

 $\operatorname{Supp}(\mu)$ will denote the closed support of a measure μ . We define the range, \mathcal{R} , of X to be

$$\mathcal{R} = \operatorname{Supp}(I)$$
.

A slightly smaller set is often used in the literature (see [3] or Corollary 9 in Ch. IV of [17]) but the definitions agree under \mathbb{P}_{δ_x} or the canonical measures \mathbb{N}_x defined below, and also give the same outcomes for $\mathcal{R} \cap \operatorname{Supp}(X_0)^c$ and $\partial \mathcal{R} \cap \operatorname{Supp}(X_0)^c$. Therefore the two definitions will be equivalent for our purposes. In dimensions $d \leq 3$, the occupation measure I has a density, L^x , which is called (total) local time of X, that is,

$$I(f) = \int_0^\infty X_s(f) ds = \int_{\mathbb{R}^d} f(x) L^x dx$$
 for all non-negative measurable f .

Moreover, $x \mapsto L^x$ is lower semicontinuous, is continuous on $\operatorname{Supp}(X_0)^c$, and for d=1 is globally continuous (see Theorems 2 and 3 of [24]). From (1.1) and (1.2) it is easy to derive (see Lemma 2.2 in [20])

(1.3)
$$\mathbb{E}_{X_0}(e^{-\lambda L^x}) = \exp\left(-\int_{\mathbb{R}^d} V^{\lambda}(x - x_0) X_0(dx_0)\right),$$

where V^{λ} is the unique solution (see Section 2 of [20] and the references given there) to

(1.4)
$$\frac{\Delta V^{\lambda}}{2} = \frac{(V^{\lambda})^2}{2} - \lambda \delta_0, \quad V^{\lambda} > 0 \text{ on } \mathbb{R}^d.$$

Thus in dimensions $d \leq 3$ we have

$$\mathcal{R} = \overline{\{x : L^x > 0\}},$$

and \mathcal{R} is a closed set of positive Lebesgue measure. In dimensions $d \geq 4$, \mathcal{R} is a Lebesgue null set of Hausdorff dimension 4 (see Theorem 1.4 of [3]), which explains our restriction to $d \leq 3$ in this work.

Our main goal in this paper is to study properties of $\partial \mathcal{R}$ — the topological boundary of \mathcal{R} , and in particular to determine the local Hausdorff dimension

of $\partial \mathcal{R}$ outside the support of X_0 . The related question of the dimension of the boundary of the set where the local time is positive, that is the dimension of

$$(1.5) F = \partial \{x : L^x > 0\},$$

was studied in [20]. To describe this latter result we introduce:

(1.6)
$$p = p(d) = \begin{cases} 3 & \text{if } d = 1\\ 2\sqrt{2} & \text{if } d = 2\\ \frac{1+\sqrt{17}}{2} & \text{if } d = 3, \end{cases}$$

 $d_f = d + 2 - p$, and

(1.7)
$$\alpha = \alpha(d) = \frac{p(d) - 2}{4 - d} = \begin{cases} 1/3 & \text{if } d = 1\\ \sqrt{2} - 1 & \text{if } d = 2\\ \frac{\sqrt{17} - 3}{2} & \text{if } d = 3. \end{cases}$$

Theorem 1.1 ([20]). With \mathbb{P}_{δ_0} -probability one,

$$\dim(F) = d_f = \begin{cases} 0 & \text{if } d = 1\\ 4 - 2\sqrt{2} \approx 1.17 & \text{if } d = 2\\ \frac{9 - \sqrt{17}}{2} \approx 2.44 & \text{if } d = 3. \end{cases}$$

There were also versions of the above in [20] for more general initial conditions X_0 .

I. Benjamini's observation that the boundary of the range exhibited interesting fractal properties in simulations was one motivation for the above. Although F may be a natural object from a stochastic analyst's perspective, the topological boundary of \mathcal{R} , $\partial \mathcal{R}$, is the more natural geometric object and of course was the set Benjamini had in mind. Clearly, $\partial \mathcal{R}$ and F are closely related; it is easy to check that

$$(1.8) \partial \mathcal{R} \subset F.$$

Thus, Theorem 1.1 gives an upper bound on dimension of $\partial \mathcal{R}$. Whether or not $F = \partial \mathcal{R}$ remains open for d = 2 or 3, but Theorem 1.7 in [20] shows that, if d = 1, there exist random variables L, R such that

(1.9)
$$F = \partial \mathcal{R} = \{\mathsf{L}, \mathsf{R}\} \text{ where } \mathsf{L} < 0 < \mathsf{R} \quad \mathbb{P}_{\delta_0} - \text{a.s.},$$

and so we will usually assume d=2 or 3. A point x in F will be in $\partial \mathcal{R}$ iff there are open sets U approaching x s.t. L=0 on U. Note that, for example, any isolated zeros of L will be in F but not in $\partial \mathcal{R}$ but we do not even know if such points exist in d=2,3. It was conjectured in (1.10) of [20] that in d=2,3,

(1.10)
$$\dim(\partial \mathcal{R}) = \dim(F) \quad \mathbb{P}_{\delta_0} - a.s..$$

In this paper we verify this conjecture, and prove the following stronger local version.

THEOREM 1.2. \mathbb{P}_{X_0} -a.s. for any open $U \subset \operatorname{Supp}(X_0)^c$

$$U \cap \partial \mathcal{R} \neq \emptyset \Rightarrow \dim(U \cap \partial \mathcal{R}) = d_f.$$

The following corollary is immediate.

Corollary 1.3. $\mathbb{P}_{X_0} - a.s.$

$$\operatorname{Supp}(X_0)^c \cap \partial \mathcal{R} \neq \emptyset \Rightarrow \dim(\operatorname{Supp}(X_0)^c \cap \partial \mathcal{R}) = d_f.$$

The hypothesis in the above Corollary is needed—see Proposition 1.5 of [20] for an example where it fails with positive probability.

COROLLARY 1.4. \mathbb{P}_{δ_0} -a.s. for any open set U,

$$U \cap \partial \mathcal{R} \neq \emptyset \Rightarrow \dim(U \cap \partial \mathcal{R}) = d_f.$$

In particular, $\dim(\partial \mathcal{R}) = d_f \quad \mathbb{P}_{\delta_0} - a.s.$

PROOF. By Theorem 1.4 and Theorem 1.6 of [11], in d=2 and d=3 we have \mathbb{P}_{δ_0} -a.s. that $L^x \to \infty$ as $x \to 0$. Therefore we can conclude that \mathbb{P}_{δ_0} -a.s. there is some $\delta > 0$ such that $L^x > 0$ for all $|x| < \delta$ and so $0 \notin \partial \mathcal{R}$, which gives $U \cap \partial \mathcal{R} \neq \emptyset \Rightarrow U \setminus \{0\} \cap \partial \mathcal{R} \neq \emptyset$. Now we may apply Theorem 1.2 with $U \setminus \{0\}$ in place of U to finish the proof.

Note that besides confirming (1.10), the above shows that the dimension result holds locally on any open ball intersecting $\partial \mathcal{R}$.

We also consider X and its local time under the canonical measures \mathbb{N}_x . Recall from Section II.7 of [21] that \mathbb{N}_x is a σ -finite measure on the space of continuous finite length $M_F(\mathbb{R}^d)$ -valued excursion paths such that

(1.11)
$$X_t = \int \nu_t \Xi(d\nu) \quad \text{for all } t > 0 \text{ under } \mathbb{P}_{X_0},$$

where Ξ is a Poisson point process with intensity $\mathbb{N}_{X_0}(\cdot) = \int \mathbb{N}_{x_0}(\cdot) X_0(dx_0)$. In this way \mathbb{N}_{x_0} governs the "excursions" of X from a single ancestor at x_0 . The existence of local time L under \mathbb{N}_x follows easily from the above, in fact it is even globally continuous (see [11]). It should not be surprising that Corollary 1.4 continues to hold under the canonical measure, in fact, as we shall see, the proof is a bit easier.

THEOREM 1.5. \mathbb{N}_0 -a.e. for any open set U,

$$U \cap \partial \mathcal{R} \neq \emptyset \Rightarrow \dim(U \cap \partial \mathcal{R}) = d_f.$$

We first say a few words about the argument leading to the proof of Theorem 1.1 in [20]. If a small ball B intersects F, then B contains a point x such that L^x is positive but small. Thus, to get the bounds on the Hausdorff dimension of F, it is useful to understand the asymptotics of $\mathbb{P}_{\delta_0}(0 < L^x < \varepsilon)$, as $\varepsilon \downarrow 0$. Write $f(t) \sim g(t)$ as $t \downarrow 0$ iff f(t)/g(t) is bounded and bounded away from zero for small positive t, and similarly for $f(t) \sim g(t)$ as $t \uparrow \infty$. It was shown in Theorem 1.3 of [20] that for p as in (1.6) and α given by (1.7),

(1.12)
$$\mathbb{P}_{\delta_0}(0 < L^x < \varepsilon) \sim |x|^{-p} \varepsilon^{\alpha}, \text{ as } \varepsilon \downarrow 0.$$

Not very difficult heuristics involving regularity properties of local time and a covering argument explains the upper bound on dimension of F: $\dim(F) \leq d_f$ (see the Introduction of [20]). (1.12) was derived in [20] through a Tauberian theorem we now sketch. Let $\lambda \uparrow \infty$ in (1.3) and (1.4) to see that $V^{\lambda}(x) \uparrow V^{\infty}(x)$ where

(1.13)
$$\mathbb{P}_{\delta_0}(L^x = 0) = \exp(-V^{\infty}(x)).$$

One important simplification available for the analysis of F in [20] is that V^{∞} is explicitly known (see e.g. (2.17) in [20]):

(1.14)
$$V^{\infty}(x) = \frac{2(4-d)}{|x|^2}.$$

In particular V^{∞} solves

(1.15)
$$\frac{\Delta V^{\infty}}{2} = \frac{(V^{\infty})^2}{2} \text{ for } x \neq 0.$$

 V^{∞} sometimes is called the *very singular solution* to (1.15), see, e.g., [2]. Applying a Tauberian theorem one can see that (1.12) can be reduced to verifying

(1.16)
$$\mathbb{E}_{\delta_0}(e^{-\lambda L^x}1(L^x > 0)) \sim |x|^{-p}\lambda^{-\alpha}, \text{ as } \lambda \uparrow \infty.$$

The left-hand side of the above behaves like $d^{\lambda}(x) := V^{\infty}(x) - V^{\lambda}(x)$, and so a substantial part of the argument in [20] was devoted to finding a rate of convergence of V^{λ} to V^{∞} , and showing that it behaves like the right hand side of (1.16).

The upper bound on $\dim(F)$ in [20] also utilized Dynkin's exit measures. For non-empty subsets K_1, K_2 of \mathbb{R}^d , we set

$$(1.17) d(K_1, K_2) := \inf\{|x - y| : x \in K_1, y \in K_2\}.$$

Define

(1.18)

 $\mathcal{O}_{X_0} \equiv \{\text{open sets } D \text{ satisfying } d(D^c, \text{supp}(X_0)) > 0 \text{ and a Brownian}$ path starting from any $x \in \partial D$ will exit D immediately $\}$.

In what follows we always assume that $G \in \mathcal{O}_{X_0}$.

The exit measure of X from such a G under \mathbb{P}_{X_0} or \mathbb{N}_{X_0} is denoted by X_G (see Chapter V of [17] for a good introduction to exit measures). X_G is a random finite measure supported on ∂G , which intuitively corresponds to the mass started at X_0 which is stopped at the instant it leaves G. The Laplace functional of X_G is given by

(1.19)
$$\mathbb{E}_{X_0}(\exp(-X_G(g))) = \exp\left(-\int 1 - \exp(-X_G(g))d\mathbb{N}_{X_0}\right)$$
$$= \exp\left(-\int U^g(x)X_0(dx)\right),$$

where $g:\partial G\to [0,\infty)$ is continuous and $U^g\geq 0$ is the unique continuous function on \overline{G} which is C^2 on G and solves

(1.20)
$$\Delta U^g = (U^g)^2 \text{ on } G, \quad U^g = g \text{ on } \partial G.$$

For this, see Theorem 6 in Chapter V of [17], and the last exercise on p. 86 for uniqueness. Let

$$(1.21) G_{\varepsilon}^{x_0} = G_{\varepsilon}(x_0) = \{x : |x - x_0| > \varepsilon\} \text{ and set } G_{\varepsilon} = G_{\varepsilon}(0).$$

Similarly $B(x_0, \varepsilon) = B_{\varepsilon}(x_0)$ is the open ball centered at x_0 and $B_{\varepsilon} = B_{\varepsilon}(0)$. $\bar{B}_{\varepsilon}(x_0)$ will denote the closed ball centered at x_0 . Proposition 3.4 of [20] gives an upper bound on $\mathbb{P}_{\delta_0}(0 < X_{G_{\varepsilon}^x}(1) < \varepsilon)$ as $\varepsilon \downarrow 0$ for $x \neq 0$. This bound is refined to precise asymptotics in Propositions 4.9 and 4.11 in Section 4 below. Intuitively these asymptotics are related to (1.12) since a small exit measure from G_{ε}^x suggests small values of the local time inside $B_{\varepsilon}(x)$. Estimates for exit measures of smalls balls are also considered in [1]. There the

interest is in the exit measure from a fixed open set and how much mass it puts on a small ball around an exit point on the boundary.

Consider next the ideas underlying Theorem 1.2, where exit measures play a more central role. To show that a point x is near $\partial \mathcal{R}$, it is not enough to show that the local time at x is small and positive, or that the exit measure from some G_{ε}^x is small. In addition, there should be balls B near x on which the local time is zero, or equivalently $X_{\bar{B}^c} = 0$. To this end we will study $\mathbb{P}_{\delta_0}(0 < X_{G_{\varepsilon}^x}(1) \leq K\varepsilon^2, X_{G_{\varepsilon/2}^x}(1) = 0)$ and show (see Theorem 4.1 and Proposition 4.9)

(1.22)
$$\mathbb{P}_{\delta_0}(0 < X_{G_{\varepsilon}^x}(1) \le K\varepsilon^2, X_{G_{\varepsilon/2}^x}(1) = 0) \sim \varepsilon^{p-2}, \text{ as } \varepsilon \downarrow 0.$$

The proof of (1.22) requires asymptotics for solutions to (1.20) with varying boundary conditions, rather than solutions to (1.4). For $\varepsilon > 0$ and $\lambda \geq 0$ we let $U^{\lambda,\varepsilon}$ denote the unique continuous function on $\{|x| \geq \varepsilon\}$ such that (cf. (1.20))

(1.23)
$$\Delta U^{\lambda,\varepsilon} = (U^{\lambda,\varepsilon})^2 \text{ for } |x| > \varepsilon, \text{ and } U^{\lambda,\varepsilon}(x) = \lambda \text{ for } |x| = \varepsilon.$$

Uniqueness of solutions implies the scaling property

(1.24)
$$U^{\lambda,\varepsilon}(x) = \varepsilon^{-2} U^{\lambda\varepsilon^2,1}(x/\varepsilon) \quad \text{for all } |x| \ge \varepsilon,$$

and also shows $U^{\lambda,\varepsilon}$ is radially symmetric, thus allowing us to write $U^{\lambda,\varepsilon}(|x|)$ for the value at $x \in \mathbb{R}^d$. By (1.19) we have for any finite initial measure satisfying $\operatorname{Supp}(X_0) \subset G_{\varepsilon}$,

(1.25)
$$\mathbb{E}_{X_0}(\exp(-\lambda X_{G_{\varepsilon}}(1))) = \exp(-X_0(U^{\lambda,\varepsilon})).$$

Let $\lambda \uparrow \infty$ in the above to see that $U^{\lambda,\varepsilon} \uparrow U^{\infty,\varepsilon}$ on G_{ε} and

(1.26)
$$\mathbb{P}_{X_0}(X_{G_{\varepsilon}}(1)=0) = \exp(-X_0(U^{\infty,\varepsilon})).$$

Proposition 9(iii) of [17] readily implies (see (3.5) and (3.6) of [20])

$$(1.27) U^{\infty,\varepsilon} \text{ is } C^2 \text{ and } \Delta U^{\infty,\varepsilon} = (U^{\infty,\varepsilon})^2 \text{ on } G_{\varepsilon},$$

$$\lim_{|x|\to\varepsilon, |x|>\varepsilon} U^{\infty,\varepsilon}(x) = +\infty, \ \lim_{|x|\to\infty} U^{\infty,\varepsilon}(x) = 0.$$

Clearly a key step in deriving (1.22) are asymptotics for

$$\mathbb{P}_{\delta_0}(0 < X_{G_{\varepsilon}^x}(1) \le K\varepsilon^2) \sim U^{\infty,\varepsilon}(x) - U^{K^{-1}\varepsilon^{-2},\varepsilon}(x) \text{ as } \varepsilon \downarrow 0,$$

where the above equivalence is by a Tauberian theorem. In Section 4.1 we show (see Corollary 4.7)

$$U^{\infty,\varepsilon}(x) - U^{K^{-1}\varepsilon^{-2},\varepsilon}(x) \sim \varepsilon^{p-2}.$$

This and a special Markov property (Propositions 2.2 and 2.3) then give (1.22). To get a lower bound on $\partial \mathcal{R}$, essentially by an inclusion-exclusion argument, in addition to the lower bound in (1.22), we will also need an upper bound on (see Proposition 5.1)

$$(1.28) \mathbb{P}_{\delta_0}(0 < X_{G_{\varepsilon}^{x_1}}(1) \le K\varepsilon^2, \ 0 < X_{G_{\varepsilon}^{x_2}}(1) \le K\varepsilon^2).$$

Although involved, this argument is quite similar to the proof of Proposition 6.1 in [20] and so is omitted (it can be found in Supplementary Material [13]). The above estimates allow us to show that the lower bound on the dimension of $\partial \mathcal{R}$ holds with positive probability—see Proposition 5.3. To conclude the proof of Theorem 1.2, we will show that the lower bound on local dimension, in fact holds with probability one. This will be a consequence of the following proposition:

PROPOSITION 1.6. Let $x_1 \in \mathbb{R}^d$, $r_0 > 0$, satisfy $B_{2r_0}(x_1) \subset \operatorname{Supp}(X_0)^c$. If $r_1 \in (0, r_0)$, then \mathbb{P}_{X_0} -a.s.

$$(1.29) X_{G_{r_1}^{x_1}}(1) = 0 and X_{G_{r_0}^{x_1}}(1) > 0 imply \dim(B_{r_0}(x_1) \cap \partial \mathcal{R}) \ge d_f.$$

The main ingredient in the proof of Proposition 1.6 is a version under the canonical measure.

PROPOSITION 1.7. Let $x_1 \in \mathbb{R}^d$, $r_0 > 0$, satisfy $B_{2r_0}(x_1) \subset \operatorname{Supp}(X_0)^c$. If $r_1 \in (0, r_0)$, then \mathbb{N}_{X_0} -a.e.

$$(1.30) \quad \begin{cases} X_{G_{r_1}^{x_1}}(1) = 0 \text{ and } X_{G_{r_0}^{x_1}}(1) > 0 \text{ imply} \\ \dim(B_r(x_1) \cap \partial \mathcal{R}) \ge d_f \text{ for every } r > r_1 \text{ s.t. } X_{G_r^{x_1}}(1) > 0. \end{cases}$$

The paper is organized as follows. In Section 2, preliminary results on super-Brownian motion, Brownian snakes, exit measures and their special Markov property are presented. In Section 3 we prove Theorems 1.2 and 1.5, assuming Propositions 1.6, 1.7.

In Section 4 left tail asymptotics of exit measures are given. First in Section 4.1 we derive necessary bounds on solutions to the boundary value problems (1.23) and (1.27), and then in Section 4.2 we prove (1.22) (see

Theorem 4.1 and Proposition 4.9). In Section 5, we show that the lower bound on the local dimension of $\partial \mathcal{R}$ holds with positive probability; see Proposition 5.3 and Lemma 5.4.

In Section 6, in preparation for the proofs of Propositions 1.6 and 1.7, we analyze the process of exit measures $(X_{G_r}(1), 0 < r \le r_0)$ through the rescaled and time-changed process, $Z(t) = X_{G_{r_0e^{-t}}}(1)/(r_0e^{-t})^2$, $t \ge 0$. We refer the reader to Section 6 for the precise definition of a continuous state branching process (CSBP).

PROPOSITION 1.8. Assume X_0 is a finite initial measure and $B_{2r_0} \subset Supp(X_0)^c$. Under \mathbb{N}_{X_0} , $(Z(t), t \geq 0)$ has a cadlag version with only non-negative jumps which is a CSBP such that for $0 \leq t_1 < t_2$ and $\lambda \geq 0$,

$$\mathbb{N}_{X_0}(\exp(-\lambda Z(t_2))|Z(s), s \le t_1) = \exp(-Z(t_1)u(t_2 - t_1, \lambda)),$$

where
$$u(t,\lambda) = e^{2t}U^{\lambda,1}(e^t)$$
 for $t,\lambda \geq 0$.

This will follow from Proposition 6.2(c) in Section 6. The absence of negative jumps in Z is important in the proof of Proposition 1.7, but we also believe that Z and its associated measure-valued process are of independent interest. Some information about the associated branching mechanism Ψ of Z and its Lévy measure are given in Remark 6.3. This complements results on the exit measure process from balls (as opposed to their complements) in [10].

The proofs of Proposition 1.6 and 1.7 are concluded in Section 7. For the proof of Proposition 1.6 one shows that for $r < r_0$ sufficiently small there is a single excursion of X (see (1.11)) governed by \mathbb{N}_{X_0} that enters B_r and thus by the monotonicity of dimension, Proposition 1.6 follows from Proposition 1.7. Proposition 1.7 (with $x_1 = 0$ without loss of generality) is proved by studying the martingale

$$(1.31) M_r = \mathbb{N}_{x_0}(\dim(B_{r_0} \cap \partial \mathcal{R}) \ge d_f|\mathcal{E}_r) 0 \le r < r_0, |x_0| > 2r_0,$$

where \mathcal{E}_r is the σ -field generated by the Brownian snake observed inside G_{r_0-r} (see Section 2 for a careful definition). In particular we analyze

$$M_r$$
 as $r \uparrow T_0 = \inf\{r : X_{G_{r_0}-r}(1) = 0\}$ on $\{0 < T_0 \le r_0 - r_1\}$,

where r_0, r_1 are as in Proposition 1.7. The special Markov property and results from Sections 5 and 6 will show $M_r \ge q > 0$, for r close enough to T_0 , and on the above set. The last step is then to show that (see Lemma 7.3)

 $\{\dim(B_{r_0} \cap \partial \mathcal{R}) \geq d_f\} \in \mathcal{E}_{T_0-}$. Now let $r \uparrow T_0$ in (1.31) to see that on $\{0 < T_0 \leq r_0 - r_1\}$, $1(\dim(B_{r_0} \cap \partial \mathcal{R}) \geq d_f) \geq q > 0$, as required.

Note that the methods used in [20] (see Theorem 1.4 and the ensuing discussion of that work) would have required the stronger hypothesis $\operatorname{Conv}(X_0)^c \cap \partial \mathcal{R} \neq \emptyset$ in Corollary 1.3, where $\operatorname{Conv}(X_0)$ is the closed convex hull of $\operatorname{Supp}(X_0)$. This is because exit measures from hyperplanes were used in [20], instead of the process of exit measure from the complements of shrinking balls. This refinement also leads to the purely local result on dimension in Theorem 1.2.

Convention on Functions and Constants. Constants whose value is unimportant and may change from line to line are denoted $C, c, c_d, c_1, c_2, \ldots$, while constants whose values will be referred to later and appear initially in say, Lemma i.j are denoted $c_{i,j}$, or $\underline{c}_{i,j}$ or $C_{i,j}$.

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2. Exit Measures and the Special Markov Property. Notation. Let \mathcal{K} be the space of compact subsets of \mathbb{R}^d equipped with the Hausdorff metric, where we add \emptyset as a discrete point. That is let $K^{\varepsilon} = \{x : d(x, K) \leq \varepsilon\}$ where $d(x, K) = \inf\{|x - y| : y \in K\}$ and for K_1, K_2 non-empty compacts, set

$$\rho(K_1, K_2) = \inf\{\varepsilon > 0 : K_1 \subset K_2^{\varepsilon} \text{ and } K_2 \subset K_1^{\varepsilon}\} \wedge 1,$$

and $\rho(\emptyset, K) = 1$ for K non-empty compact. (K, ρ) is then a complete separable metric space. If U is an open set in \mathbb{R}^d we let C(U) be the space of continuous functions on U with the compact-open topology.

We start with a measurability result requiring a bit of care; a proof is given in the Appendix.

LEMMA 2.1. (a) For any R > 0, $\psi_a : \mathcal{K} \to \mathcal{K}$ is a Borel map, where $\psi_a(K) = K \cap \overline{B_R}$.

(b) For any $\alpha, R > 0$, $\psi_b : \mathcal{K} \to \mathbb{R}$ is a universally measurable map, where $\psi_b(K) = 1(\dim(\partial K) \cap B_R) < \alpha$.

We will use Le Gall's Brownian snake construction of a SBM X, with initial state $X_0 \in M_F(\mathbb{R}^d)$. Set $\mathcal{W} = \bigcup_{t \geq 0} C([0,t],\mathbb{R}^d)$ with the natural metric (see page 54 of [17]), and let $\zeta(w) = t$ be the lifetime of $w \in C([0,t],\mathbb{R}^d) \subset \mathcal{W}$. The Brownian snake $W = (W_t, t \geq 0)$ is a \mathcal{W} -valued continuous strong Markov process and, abusing notation slightly, let \mathbb{N}_x denote its excursion measure starting from the path at $x \in \mathbb{R}^d$ with lifetime zero. As usual we

let $\hat{W}(t) = W_t(\zeta(W_t))$ denote the tip of the snake at time t, and $\sigma(W) > 0$ denote the length of the excursion path. We refer the reader to Ch. IV of [17] for the precise definitions. The construction of super-Brownian motion, X = X(W) under \mathbb{N}_x or \mathbb{P}_{X_0} , may be found in Ch. IV of [17]. The "law" of X(W) under \mathbb{N}_x is the canonical measure of SBM starting at x described in the last Section (and also denoted by \mathbb{N}_x). If $\Xi = \sum_{j \in J} \delta_{W_j}$ is a Poisson point process on \mathcal{W} with intensity $\mathbb{N}_{X_0}(dW) = \int \mathbb{N}_x(dW) X_0(dx)$, then by Theorem 4 of Ch. IV of [17] (cf. (1.11))

(2.1)
$$X_t(W) = \sum_{j \in J} X_t(W_j) = \int X_t(W) \Xi(dW) \text{ for } t > 0$$

defines a SBM with initial measure X_0 . We will refer to this as the standard set-up for X under \mathbb{P}_{X_0} .

Recall $\mathcal{R} = \overline{\{x : L^x > 0\}}$ is the range of the SBM X under \mathbb{P}_{X_0} or \mathbb{N}_{X_0} . Under \mathbb{N}_{X_0} we have (see (8) on p. 69 of [17])

$$(2.2) \mathcal{R} = {\hat{W}(s) : s \in [0, \sigma]}.$$

Let $G \in \mathcal{O}_{X_0}$. Then

(2.3) X_G is a finite random measure supported on $\mathcal{R} \cap \partial G$ a.s..

Under \mathbb{N}_{X_0} this follows from the definition of X_G on p. 77 of [17] and the ensuing discussion, and (2.2). Although [17] works under \mathbb{N}_x for $x \in G$ the above extends immediately to \mathbb{P}_{X_0} because as in (2.23) of [20],

(2.4)
$$X_G = \sum_{j \in J} X_G(W_j) = \int X_G(W) d\Xi(W),$$

where Ξ is a Poisson point process on \mathcal{W} with intensity \mathbb{N}_{X_0} . Working under \mathbb{N}_{X_0} and following [16], we define

$$S_G(W_u) = \inf\{t \le \zeta_u : W_u(t) \notin G\} \quad (\inf \emptyset = \infty),$$
$$\eta_s^G(W) = \inf\{t : \int_0^t 1(\zeta_u \le S_G(W_u)) \, du > s\},$$
$$\mathcal{E}_G = \sigma(W_{\eta_s^G}, s \ge 0) \vee \{\mathbb{N}_{X_0} - \text{null sets}\},$$

where $s \to W_{\eta_s^G}$ is continuous (see p. 401 of [16]). Write the open set $\{u: S_G(W_u) < \zeta_u\}$ as countable union of disjoint open intervals, $\cup_{i \in I}(a_i, b_i)$. Clearly $S_G(W_u) = S_G^i < \infty$ for all $u \in [a_i, b_i]$ and we may define

$$W_s^i(t) = W_{(a_i+s) \wedge b_i}(S_G^i + t) \text{ for } 0 \le t \le \zeta_{(a_i+s) \wedge b_i} - S_G^i.$$

Therefore for $i \in I$, $W^i \in C(\mathbb{R}_+, \mathcal{W})$ are the excursions of W outside G. Proposition 2.3 of [16] implies X_G is \mathcal{E}_G -measurable and Corollary 2.8 of the same reference implies

(2.5) $\begin{cases} \text{Conditional on } \mathcal{E}_G, \text{ the point measure } \sum_{i \in I} \delta_{W^i} \text{ is a Poisson} \\ \text{point measure with intensity } \mathbb{N}_{X_G}. \end{cases}$

If D is an open set in \mathcal{O}_{X_0} such that $\bar{G} \subset D$ and $d(D^c, \bar{G}) > 0$ (recall (1.17)), then the definition (and existence) of $X_D(W)$ applies equally well to each $X_D(W^i)$ and it is easy to check that

(2.6)
$$X_D(W) = \sum_{i \in I} X_D(W^i).$$

If U is an open subset of $\operatorname{Supp}(X_0)^c$, then L_U , the restriction of the local time L^x to U, is in C(U). Here are some simple consequences of (2.5).

PROPOSITION 2.2. (a) Let G_1, G_2 be open sets in \mathcal{O}_{X_0} such that $\overline{G_1} \subset G_2$ and $d(G_2^c, \overline{G_1}) > 0$.

(i) If $\psi_1: C(\overline{G_1}^c) \to [0, \infty)$ is Borel measurable, then

$$\mathbb{N}_{X_0}(\psi_1(L_{\overline{G_1}^c})|\mathcal{E}_{G_1}) = \mathbb{E}_{X_{G_1}}(\psi_1(L_{\overline{G_1}^c})).$$

(ii) If $\psi_2: M_F(\mathbb{R}^d) \to [0, \infty)$ is Borel measurable then

$$\mathbb{N}_{X_0}(\psi_2(X_{G_2})|\mathcal{E}_{G_1}) = \mathbb{E}_{X_{G_1}}(\psi_2(X_{G_2})).$$

(b) If $0 < R_2 < R_1$, $d(Supp(X_0), \overline{B_{R_1}}) > 0$, and $\psi_3 : \mathcal{K} \to [0, \infty)$ is Borel measurable, then

$$\mathbb{N}_{X_0}(\psi_3(\mathcal{R}\cap \overline{B_{R_2}})|\mathcal{E}_{G_{R_1}}) = \mathbb{E}_{X_{G_{R_1}}}(\psi_3(\mathcal{R}\cap \overline{B_{R_2}})),$$

where G_R is as in (1.21).

PROOF. (a) (i) is Proposition 2.6(b) of [20]. (a)(ii) follows in a similar manner from (2.5), (2.6) and (2.4).

(b) Define $S: C(B_{R_1}) \to \mathcal{K}$ by $S(f) = \operatorname{Supp}(f) := \overline{\{x: f(x) > 0\}}$, where the closure is taken in all of \mathbb{R}^d . Then it is easy to see that S is Borel measurable, for example by considering the inverse images of closed balls in \mathcal{K} . In addition the map $K \to \overline{B_{R_2}} \cap K$ is measurable on \mathcal{K} by Lemma 2.1(a). Now observe that $\mathcal{R} \cap \overline{B_{R_2}} = S(L_{B_{R_1}}) \cap \overline{B_{R_2}}$, and so by the above observations is a measurable function of $L_{B_{R_1}}$. Therefore (b) now follows from (a)(i) with $G_i = G_{R_i}$.

We will need a version of the above under \mathbb{P}_{X_0} as well.

PROPOSITION 2.3. For $X_0 \in M_F(\mathbb{R}^d)$ and an open set G in \mathcal{O}_{X_0} , let Ψ be a bounded measurable function on $C(\overline{G}^c)$ and Φ_i , i = 0, 1 be bounded measurable functions on $M_F(\mathbb{R}^d)$ and $M_F(\mathbb{R}^d)^n$, respectively. Then $(a) \mathbb{E}_{X_0}(\Phi_0(X_G)\Psi(L)) = \mathbb{E}_{X_0}(\Phi_0(X_G)\mathbb{E}_{X_G}(\Psi(L)))$.

(b) (i) Let D_i be open sets in \mathcal{O}_{X_0} , such that $d(D_i^c, \bar{G}) > 0$, $\forall 1 \leq i \leq n$. Then

$$\mathbb{E}_{X_0}\Big(\Phi_0(X_G)\Phi_1(X_{D_1},\ldots,X_{D_n})\Big) = \mathbb{E}_{X_0}\Big(\Phi_0(X_G)\mathbb{E}_{X_G}\Big(\Phi_1(X_{D_1},\ldots,X_{D_n})\Big)\Big).$$

(ii) If
$$0 < R_2 < R_1$$
 and $d(Supp(X_0)^c, \overline{B_{R_1}}) > 0$, then

$$\mathbb{E}_{X_0}(\Phi_0(X_{G_{R_1}})1(\mathcal{R}\cap\overline{B_{R_2}}\neq\emptyset))=\mathbb{E}_{X_0}(\Phi_0(X_{G_{R_1}})\mathbb{P}_{X_{G_{R_1}}}(\mathcal{R}\cap\overline{B_{R_2}}\neq\emptyset)).$$

PROOF. (a) is Proposition 2.6(c) of [20]. (b)(i) follows by the same reasoning there, using (2.5), (2.4) (the latter for each D_i , as well as G), and Proposition 2.2(a)(ii), trivially extended to accommodate $(X_{D_1}, \ldots, X_{D_n})$ in place of X_{G_2} . (b)(ii) follows from (a), as in the proof of Proposition 2.2(b).

3. Proof of Theorems 1.2 and 1.5. We will see in this section that (using the upper bound on $\dim(F)$ in Theorem 1.1) Theorem 1.2 is a simple consequence of Proposition 1.6, and similarly Theorem 1.5 can be derived from Proposition 1.7.

Proof of Theorem 1.2. Let $x_1 \in \mathbb{R}^d$ and $0 < r_1 < r_0 \le 1$ satisfy $B_{2r_0}(x_1) \subset \operatorname{Supp}(X_0)^c$. From (2.3) we have $\mathbb{P}_{X_0} - \operatorname{a.s.}$,

(3.1)
$$\partial G_{r_1}^{x_1} \cap \mathcal{R} = \emptyset \Rightarrow X_{G_{r_1}^{x_1}} = 0.$$

Proposition 2.3(b)(ii) and translation invariance imply

$$\mathbb{P}_{X_0}(X_{G_{r_0}^{x_1}} = 0, \mathcal{R} \cap B_{r_0/2}(x_1) \neq \emptyset)$$

$$= \mathbb{E}_{X_0}(1(X_{G_{r_0}^{x_1}} = 0)\mathbb{P}_{X_{G_{r_0}^{x_1}}}(\mathcal{R} \cap B_{r_0/2}(x_1) \neq \emptyset)) = 0.$$

It follows that \mathbb{P}_{X_0} – a.s.,

(3.2)
$$\mathcal{R} \cap B_{r_0/2}(x_1) \neq \emptyset \Rightarrow X_{G_{r_0}^{x_1}}(1) > 0.$$

Fix ω outside a \mathbb{P}_{X_0} -null set so that (1.29) of Proposition 1.6, (3.1), and (3.2) all hold for all $x_1 \in \mathbb{Q}^d$ and all rational numbers $0 < r_1 < r_0 \le 1$ satisfying $B_{2r_0}(x_1) \subset \operatorname{Supp}(X_0)^c$. Assume U is an open set in $\operatorname{Supp}(X_0)^c$

which intersects $\partial \mathcal{R}$ and choose $x_0 \in U \cap \partial \mathcal{R}$. Pick a rational r_0 in (0,1] so that

$$B_{3r_0}(x_0) \subset U \subset \operatorname{Supp}(X_0)^c$$
,

then choose $x_1 \in \mathbb{Q}^d \cap B_{r_0/2}(x_0) \cap \mathbb{R}^c$, and finally select a rational $r_1 \in (0, r_0)$ such that

(3.3)
$$B_{2r_1}(x_1) \subset \mathcal{R}^c \text{ and so } \partial G_{r_1}^{x_1} \cap \mathcal{R} = \emptyset.$$

Clearly we have

$$(3.4) B_{2r_0}(x_1) \subset B_{3r_0}(x_0) \subset U \subset \text{Supp}(X_0)^c,$$

and

(3.5)
$$x_0 \in B_{r_0/2}(x_1) \cap \partial \mathcal{R} \text{ and so } \mathcal{R} \cap B_{r_0/2}(x_1) \neq \emptyset.$$

Our choice of ω and (3.4) allow us to conclude from (3.1), (3.3) and (3.2), (3.5) respectively, that

$$(3.6) \hspace{1cm} X_{G^{x_1}_{r_1}}(1) = 0 \text{ and } X_{G^{x_1}_{r_0}}(1) > 0, \text{respectively}.$$

By (3.4) and our choice of ω we may also apply Proposition 1.6 with (3.6) and conclude that

$$\dim(U \cap \partial \mathcal{R}) \ge \dim(B_{r_0}(x_1) \cap \partial \mathcal{R}) \ge d_f$$

where we have used (3.4) in the first inequality. On the other hand we know from Theorem 1.4(a) of [20] and $\partial \mathcal{R} \subset \partial \{x : L^x > 0\}$ that

$$\dim(U \cap \partial \mathcal{R}) \leq \dim(\operatorname{Supp}(X_0)^c \cap \partial \mathcal{R}) \leq d_f$$

and the proof is complete.

Proof of Theorem 1.5. As in the above proof of Theorem 1.2 from Proposition 1.6, we can derive from Proposition 1.7 that \mathbb{N}_{X_0} -a.e. for any open set $U \subset \operatorname{Supp}(X_0)^c$,

$$(3.7) U \cap \partial \mathcal{R} \neq \emptyset \Rightarrow \dim(U \cap \partial \mathcal{R}) = d_f.$$

Now we turn to the \mathbb{N}_0 case in Theorem 1.5. By (4.4) in [12], we have \mathbb{N}_0 -a.e. that $L^0 > 0$. Although the arguments in [12] were given for d = 1, they work in any dimension $d \leq 3$. Since L^x is globally continuous under \mathbb{N}_0 by Theorem 1.2 of [11], we can conclude that \mathbb{N}_0 -a.e. there is some $\delta > 0$ such that $L^x > 0$ for all $|x| < \delta$, and so $0 \notin \partial \mathcal{R}$, which gives $U \cap \partial \mathcal{R} \neq \emptyset \Rightarrow U \setminus \{0\} \cap \partial \mathcal{R} \neq \emptyset$. Now we may apply (3.7) with $U \setminus \{0\}$ in place of U to finish the proof.

4. Lower Bound on the Exit Measure Probability. Throughout this Section we fix $\varepsilon_0 \in (0,1)$. As noted in the Introduction, the goal of this section, stated below, is a key estimate for the lower bound on the dimension of $\partial \mathcal{R}$. Although we are interested in d=2,3, we assume $d\leq 3$ throughout this section as the arguments remain valid.

THEOREM 4.1. There are positive constants $R_{4.1}$, $K_1(\varepsilon_0) < K_2(\varepsilon_0) < \infty$ and $c_{4.1}(\varepsilon_0)$ such that, for all $\varepsilon_0 \leq |x| \leq \varepsilon_0^{-1}$,

$$\mathbb{P}_{\delta_x}(K_1 \le \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le K_2, X_{G_{\varepsilon/2}(1)} = 0) \ge c_{4.1}(\varepsilon_0)\varepsilon^{p-2}, \ \forall \ 0 < \varepsilon < \varepsilon_0/R_{4.1}.$$

The next subsection is devoted to proving necessary bounds on solutions to the boundary value problems (1.23), (1.27). These bounds will be used for proving Theorem 4.1 in Section 4.2.

4.1. Bounds on Solutions to some Boundary Value Problems. Recall $U^{\lambda,R}$ and $U^{\infty,R}$ from (1.23) and (1.27), respectively. A simple application of (1.27), (1.15) and the maximum principle implies

$$(4.1) V^{\infty}(x) \le U^{\infty,1}(x) \quad \forall |x| > 1.$$

We will need an upper bound on $U^{\infty,1}$ which shows this bound is asymptotically sharp for large |x|. We briefly include d=1 in our analysis. Recall that p is as in (1.6).

PROPOSITION 4.2. There exist constants $C_{4,2} > 1$ and $c_{4,2} \ge 0$ such that

$$U^{\infty,1}(x) \le V^{\infty}(x)(1 + c_{4.2}|x|^{2-p}) \quad \forall |x| \ge C_{4.2}.$$

PROOF. We will write u(r) for $U^{\infty,1}(r)$ and v(r) for $V^{\infty}(r)$. Let

$$q(t) = \frac{u(e^{t/4})}{v(e^{t/4})} = \frac{1}{4}u(e^{t/4})e^{t/2}, t \ge 0, \text{ in } d = 2,$$

$$q(t) = \frac{u(e^{t/3})}{v(e^{t/3})} = \frac{1}{2}u(e^{t/3})e^{2t/3}, t \ge 0, \text{ in } d = 3,$$

$$q(t) = \frac{u(e^{t/5})}{v(e^{t/5})} = \frac{1}{6}u(e^{t/5})e^{2t/5}, t \ge 0, \text{ in } d = 1.$$

Consider first d=2. A simple calculation gives

$$(4.2) \qquad \frac{1}{2}q'' - \frac{1}{2}q' + \beta(q - q^2) = 0 \text{ on } (0, \infty), \ q(0) = \infty, \ \lim_{t \to \infty} q(t) = 1,$$

where $\beta=\frac{1}{8}$ in d=2 (similarly $\beta=\frac{3}{25}$ in d=1 and $\beta=\frac{1}{9}$ in d=3). Note that $q(0)=\infty$ by the definitions of $U^{\infty,1}$ and V^{∞} . Also note that the last limit is derived the same way as (3.10) in [20] by taking $U^{\infty,1}$ instead of $U^{\delta_0,1}$, $\tilde{y}(t)=y(t+2)$ instead of y(t) and $\tilde{z}(t)=z(t+2)$ instead of z(t) there.

Note that $q(x) \ge 1$ for all $x \ge 0$ by (4.1). Define

$$z = q - 1$$
,

and thus by the above $z(x) \ge 0$ for all $x \ge 0$. Then z satisfies the following equation:

$$(4.3) \quad \frac{1}{2}z'' - \frac{1}{2}z' - \beta z(z+1) = 0 \text{ on } (0,\infty), \ z(0) = \infty, \ \lim_{t \to \infty} z(t) = 0.$$

Let w be the unique solution to

$$(4.4) \quad \frac{1}{2}w''(t) - \frac{1}{2}w'(t) - \beta w(t) = 0, \ t > 1, w(1) = z(1), \ \lim_{t \to \infty} w(t) = 0.$$

By the comparison principle and using $z + 1 \ge 1$, we get

(4.5)
$$w(t) \ge z(t), t \ge 1.$$

We leave it for the reader to check that

$$w(t) = z(1)e^{-\lambda}e^{\lambda t}, \ t \ge 1,$$

with

$$\lambda = \frac{1}{2} - \sqrt{\frac{1}{4} + 2\beta} < 0.$$

By the definition of β we have $\lambda = 1/2 - \frac{1}{\sqrt{2}}$ in d=2 (similarly $\lambda = -0.2$ for d=1 and $\lambda = 1/2 - \frac{\sqrt{17}}{6}$ for d=3). This and (4.5) imply that for $C=z(1)e^{-\lambda} \geq 0$ we have

$$z(x) \le Ce^{\lambda x}, \ x \ge 1,$$

and since $\lambda < 0$ we get that z decreases to zero exponentially fast. Recall the definition of q to get

$$q(x) \le 1 + Ce^{\lambda x}, \ x \ge 1.$$

Then (as d=2) we have,

$$u(e^{t/4}) \le v(e^{t/4})(1 + Ce^{\lambda t}), \ t \ge 1.$$

and so

$$u(s) \le v(s)(1 + Cs^{4\lambda}) = v(s)(1 + Cs^{2-p}), \ s \ge e^{1/4}.$$

Similar algebra shows the result in d = 1, 3.

Recalling (1.14) and that p > 2, we may immediately conclude:

COROLLARY 4.3. There are constants $C_{4.3}, c_{4.3} > 0$ such that for all $x \in \mathbb{R}^d$, with $|x| \geq C_{4.3}$, we have

(4.6)
$$U^{\infty,1}(x) \le V^{\infty}(x) + \frac{c_{4.3}}{|x|^p}.$$

In particular, there is some constant $K_{4,3} > 2$ such that

$$(4.7) U^{\infty,1}(x) \le 3(4-d)|x|^{-2}, \ \forall |x| \ge K_{4.3}.$$

If $D^{\lambda}=U^{\infty,1}-U^{\lambda,1}\geq 0$ for $\lambda>0$, then the Feynmann-Kac formula (as in (3.8) in [20]) easily gives

$$D^{\lambda}(x) = D^{\lambda}(R)E_x\left(1_{(\tau_R < \infty)} \exp\left(-\int_0^{\tau_R} \left(\frac{U^{\infty, 1} + U^{\lambda, 1}}{2}\right)(B_s)ds\right)\right), |x| \ge R > 1,$$

where B denote a d-dimensional Brownian motion starting at x under P_x and $\tau_R = \inf\{t \ge 0 : |B_t| \le R\}$ for $|x| \ge R > 1$.

We will frequently use the following lemmas. For $\gamma \in \mathbb{R}$, let (ρ_t) denote a γ -dimensional Bessel process starting from r > 0 under $P_r^{(\gamma)}$. For R > 0, let $\tau_R = \inf\{t \geq 0 : \rho_t \leq R\}$. The following result is from Lemma 5.3 of [20].

LEMMA 4.4. Assume $0 < 2\gamma \le \nu^2$ and q > 2. Then

$$\sup_{r\geq 1} E_r^{(2+2\nu)} \left(\exp\left(\int_0^{\tau_1} \frac{\gamma}{\rho_s^q} \, ds \right) \middle| \tau_1 < \infty \right) \leq c_{4.4}(q,\nu) < \infty.$$

Lemma 4.5. Let $q>2,\ a\in\mathbb{R},\ \zeta\in[0,2(4-d)),\ \nu_{\zeta}=\sqrt{\nu^2-\zeta}$ and $p_{\zeta}=\nu_{\zeta}+\mu,\ where$

(4.9)
$$\mu = \begin{cases} -1/2 & \text{if } d = 1, \\ 0 & \text{if } d = 2, \text{ and } \nu = \sqrt{\mu^2 + 4(4 - d)}, \\ 1/2 & \text{if } d = 3, \end{cases}$$

Then for all R < |x|,

$$(4.10) \quad E_x \left(1(\tau_R < \infty) \exp\left(\int_0^{\tau_R} \frac{a}{|B_s|^q} ds \right) \exp\left(-\int_0^{\tau_R} \frac{2(4-d) - \zeta/2}{|B_s|^2} ds \right) \right)$$
$$= E_{|x|}^{(2+2\nu_{\zeta})} \left(\exp\left(\int_0^{\tau_R} \frac{a}{\rho_s^q} ds \right) \Big| \tau_R < \infty \right) (R/|x|)^{p_{\zeta}}.$$

PROOF. The proof is based on arguments from [20] (see the proof of Lemma 5.4 there), which in turn go back to Marc Yor [26] and is deferred to Appendix A.2.

PROPOSITION 4.6. There are positive universal constants $C_{4.6}$, $c_{4.6} > 0$, $K_{4.6} > K_{4.3}$, and $R_{4.6} > 2$ such that

(a)
$$D^{\lambda}(x) \le \frac{R^p}{|x|^p} D^{\lambda}(R), \ \forall |x| \ge R > 1, \ \lambda \ge 2(4 - d).$$

(b)
$$D^{\lambda}(x) \le C_{4.6} \frac{R^p}{|x|^p} D^{\lambda}(R), \ \forall |x| \ge R \ge \frac{K_{4.6}}{\lambda}, 0 < \lambda < 1.$$

(c)
$$D^{\lambda}(x) \ge c_{4.6} \frac{R^p}{|x|^p} D^{\lambda}(R) > 0, \ \forall |x| \ge R \ge R_{4.6}, \ \lambda > 0.$$

PROOF. Recall μ , ν introduced in (4.9) so that for p = p(d) defined as in (1.6), we have

$$(4.11) p = \mu + \nu.$$

(a) For
$$\lambda \geq 2(4-d)$$
, clearly we have $U^{\lambda,1}(1) = \lambda \geq 2(4-d) = V^{\infty}(1)$. As
$$\lim_{|x| \to \infty} U^{\lambda,1}(x) \leq \lim_{|x| \to \infty} U^{\infty,1}(x) = 0$$

by (1.27), we may apply the maximum principle to get

(4.12)
$$U^{\lambda,1}(x) \ge V^{\infty}(x) = \frac{2(4-d)}{|x|^2}, \quad \forall |x| > 1.$$

Use (4.1) and the above to see that (4.8) becomes

$$D^{\lambda}(x) \leq D^{\lambda}(R) E_x \Big(1_{(\tau_R < \infty)} \exp \Big(- \int_0^{\tau_R} \frac{2(4-d)}{|B_s|^2} ds \Big) \Big) = D^{\lambda}(R) (R/|x|)^p,$$

the last by Lemma 4.5.

(b) Assume $\lambda \in (0,1)$. Recall Proposition 3.3(b) in [20]:

$$\forall \delta \in (0,1), \ \exists C_{\delta} > 2, \text{ so that } U^{\lambda,1}(x) \ge (1-\delta)V^{\infty}(x) \text{ for all } |x| \ge C_{\delta}/\lambda.$$

For any
$$\delta \in (0,1)$$
, let $\zeta = 2(4-d)\delta \in (0,2(4-d))$. Let μ and ν be as in (4.9). Define $\nu_{\zeta} = \sqrt{\mu^2 + 4(4-d) - \zeta}$ and $p_{\zeta} = \nu_{\zeta} + \mu \to p > 2$ as $\zeta \downarrow 0$.

Choose $\delta \in (0,1)$ small enough so that $p_{\zeta} > 2$. Let $K_{4.6} \equiv C_{\delta} + K_{4.3}$. Now use (4.1), (4.13) and Lemma 4.5 to see that for $|x| \geq R \geq K_{4.6}/\lambda > C_{\delta}/\lambda$, (4.8) implies

$$D^{\lambda}(x) \le D^{\lambda}(R) E_x \left(1(\tau_R < \infty) \exp\left(-\int_0^{\tau_R} \frac{2(4-d) - (\zeta/2)}{|B_s|^2} ds \right) \right)$$
$$= D^{\lambda}(R) (R/|x|)^{p_{\zeta}}.$$

Let $\xi(R) = D^{\lambda}(R)R^{p_{\zeta}}/2$. Then the above gives

$$\left(\frac{U^{\lambda,1} + U^{\infty,1}}{2}\right)(x) \ge U^{\infty,1}(x) - \frac{\xi(R)}{|x|^{p_{\zeta}}} \ge V^{\infty}(x) - \frac{\xi(R)}{|x|^{p_{\zeta}}} \text{ for } |x| \ge R.$$

Use this in (4.8) and then Lemma 4.5 to see that for $|x| \ge R$,

$$D^{\lambda}(x) \leq D^{\lambda}(R) E_{x} \Big(1(\tau_{R} < \infty) \exp\Big(\int_{0}^{\tau_{R}} \frac{\xi(R)}{|B_{s}|^{p_{\zeta}}} ds \Big) \exp\Big(- \int_{0}^{\tau_{R}} \frac{2(4-d)}{|B_{s}|^{2}} ds \Big) \Big)$$
$$= D^{\lambda}(R) E_{|x|}^{(2+2\nu)} \Big(\exp\Big(\int_{0}^{\tau_{R}} \frac{\xi(R)}{\rho_{s}^{p_{\zeta}}} ds \Big) \Big| \tau_{R} < \infty \Big) (R/|x|)^{p}.$$

A scaling argument shows that the above equals

$$D^{\lambda}(R)(R/|x|)^{p} E_{|x|/R}^{(2+2\nu)} \Big(\exp\Big(\int_{0}^{\tau_{1}} \frac{\xi(R)R^{2-p_{\zeta}}}{\rho_{s}^{p_{\zeta}}} ds \Big) |\tau_{1} < \infty \Big).$$

To apply Lemma 4.4 we note that by (4.7), for $R \ge K_{4.6}/\lambda > K_{4.3}$ we have

$$2\gamma \equiv 2\xi(R)R^{2-p_{\zeta}} \le U^{\infty,1}(R)R^2 \le 3(4-d) < \nu^2.$$

So Lemma 4.4 and the above bound show that

(4.14)
$$D^{\lambda}(x) \le D^{\lambda}(R)(R/|x|)^{p} c_{4.4}(p_{\zeta}, \nu).$$

(c) Use (4.6) in Corollary 4.3 to see that for $|x| \geq R > C_{4.3}$, we have

$$\frac{U^{\infty,1} + U^{\lambda,1}}{2}(x) \le U^{\infty,1}(x) \le \frac{2(4-d)}{|x|^2} + \frac{c_{4,3}}{|x|^p}.$$

So (4.8) and Lemma 4.5 imply

$$D^{\lambda}(x) \ge D^{\lambda}(R) E_x \Big(1_{(\tau_R < \infty)} \exp\Big(-\int_0^{\tau_R} \frac{c_{4.3}}{|B_s|^p} ds \Big) \exp\Big(-\int_0^{\tau_R} \frac{2(4-d)}{|B_s|^2} ds \Big) \Big)$$

$$= D^{\lambda}(R) E_{|x|}^{(2+2\nu)} \Big(\exp\Big(-\int_0^{\tau_R} \frac{c_{4.3}}{\rho_s^p} ds \Big) |\tau_R < \infty \Big) (R/|x|)^p,$$

with $p = \mu + \nu$. A scaling argument shows that the above equals

$$(4.15) D^{\lambda}(R)(R/|x|)^{p} E_{|x|/R}^{(2+2\nu)} \Big(\exp\Big(- \int_{0}^{\tau_{1}} \frac{c_{4.3}R^{2-p}}{\rho_{s}^{p}} ds \Big) |\tau_{1} < \infty \Big).$$

To apply Lemma 4.4 note that if $R \ge R_{4.6}$ for some constant $R_{4.6} > 2$,

$$2\gamma \equiv 2c_{4,3}R^{2-p} < 2(4-d) < \nu^2.$$

By Cauchy-Schwartz, we have

$$1 = \left(E_{|x|/R}^{(2+2\nu)} \left(\exp\left(-\int_0^{\tau_1} \frac{c_{4.3}R^{2-p}}{2\rho_s^p} ds \right) \exp\left(\int_0^{\tau_1} \frac{c_{4.3}R^{2-p}}{2\rho_s^p} ds \right) | \tau_1 < \infty \right) \right)^2$$

$$\leq E_{|x|/R}^{(2+2\nu)} \left(\exp\left(-\int_0^{\tau_1} \frac{c_{4.3}R^{2-p}}{\rho_s^p} ds \right) | \tau_1 < \infty \right)$$

$$\times E_{|x|/R}^{(2+2\nu)} \left(\exp\left(\int_0^{\tau_1} \frac{c_{4.3}R^{2-p}}{\rho_s^p} ds \right) | \tau_1 < \infty \right)$$

$$\leq E_{|x|/R}^{(2+2\nu)} \left(\exp\left(-\int_0^{\tau_1} \frac{c_{4.3}R^{2-p}}{\rho_s^p} ds \right) | \tau_1 < \infty \right) c_{4.4}(p, \nu),$$

the last by Lemma 4.4. Hence

$$\inf_{|x| \ge R} E_{|x|/R}^{(2+2\nu)} \left(\exp\left(-\int_0^{\tau_1} \frac{c_{4.3} R^{2-p}}{\rho_s^p} ds \right) \Big| \tau_1 < \infty \right) \ge c_{4.4}(p, \nu)^{-1} > 0,$$

and by (4.15) we are done.

By using the scaling relations of $U^{\infty,\varepsilon}$ and $U^{\lambda\varepsilon^{-2},\varepsilon}$ from (1.24), the following is immediate from the above.

COROLLARY 4.7. For all $\varepsilon > 0$, we have (a)

$$U^{\infty,\varepsilon}(x) - U^{\lambda\varepsilon^{-2},\varepsilon}(x) \le \frac{R^p}{|x|^p} D^{\lambda}(R)\varepsilon^{p-2}, \ \forall |x|/\varepsilon \ge R > 1, \ \lambda \ge 2(4-d).$$

(b)

$$U^{\infty,\varepsilon}(x) - U^{\lambda\varepsilon^{-2},\varepsilon}(x) \le C_{4.6} \frac{R^p}{|x|^p} D^{\lambda}(R) \varepsilon^{p-2}, \ \forall |x|/\varepsilon \ge R \ge \frac{K_{4.6}}{\lambda}, 0 < \lambda < 1.$$

(c)

$$U^{\infty,\varepsilon}(x) - U^{\lambda\varepsilon^{-2},\varepsilon}(x) \ge c_{4.6} \frac{R^p}{|x|^p} D^{\lambda}(R) \varepsilon^{p-2} > 0, \ \forall |x|/\varepsilon \ge R \ge R_{4.6}, \ \lambda > 0.$$

4.2. The left tail of the total exit measure and Proof of Theorem 4.1.

PROPOSITION 4.8. For any $|x| \ge \varepsilon_0$ and $\varepsilon \in (0, \varepsilon_0/2)$, we have

$$\mathbb{P}_{\delta_x} \left(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le \frac{1}{\lambda} \right) \le e^{\frac{2^p}{|x|^p}} D^{\lambda}(2) \varepsilon^{p-2}, \ \forall \lambda \ge 2(4-d).$$

PROOF. Apply Markov's inequality to get

$$\mathbb{P}_{\delta_{x}}\left(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^{2}} \leq \frac{1}{\lambda}\right) \leq e\mathbb{E}_{\delta_{x}}\left(\exp(-\lambda \varepsilon^{-2} X_{G_{\varepsilon}}(1)) 1(X_{G_{\varepsilon}}(1) > 0)\right)$$

$$= e\left(\exp(-U^{\lambda \varepsilon^{-2}, \varepsilon}(x)) - \exp(-U^{\infty, \varepsilon}(x))\right)$$

$$\leq e\left(U^{\infty, \varepsilon}(x) - U^{\lambda \varepsilon^{-2}, \varepsilon}(x)\right),$$

$$(4.16)$$

the equality by (1.25) and (1.26). Note we've assumed $|x|/\varepsilon \geq \varepsilon_0/\varepsilon > 2$ and $\lambda \geq 2(4-d)$ so that we can use Corollary 4.7(a) with R=2 to bound the right-hand side of (4.16) by $e(2/|x|)^p D^{\lambda}(2)\varepsilon^{p-2}$, as required.

PROPOSITION 4.9. There is some $c_{4.9}(\varepsilon_0) > 0$ such that for all $|x| \ge \varepsilon_0$ and $\varepsilon \in (0, \varepsilon_0)$,

$$\mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le \frac{1}{\lambda}) \le c_{4.9}(\varepsilon_0)\lambda^{-(p-2)}\varepsilon^{p-2}, \ \forall \ 0 < \lambda < 1.$$

PROOF. For $\lambda \in (0,1)$ such that

$$|x|/\varepsilon \ge \varepsilon_0/\varepsilon \ge K_{4.6}/\lambda$$
,

we apply Markov's inequality as in (4.16) and use Corollary 4.7(b) with $R = K_{4.6}/\lambda$ to get

$$\mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le \frac{1}{\lambda}) \le e\left(U^{\infty,\varepsilon}(x) - U^{\lambda\varepsilon^{-2},\varepsilon}(x)\right) \le eC_{4.6}U^{\infty,1}(R)(R/|x|)^p \varepsilon^{p-2}$$

$$\le eC_{4.6}(3(4-d)/R^2)(R/\varepsilon_0)^p \varepsilon^{p-2} \le 9eC_{4.6}\varepsilon_0^{-p} K_{4.6}^{p-2} \lambda^{-(p-2)} \varepsilon^{p-2},$$

the second line by (4.7) with $R = K_{4.6}/\lambda > K_{4.6} > K_{4.3}$. For $\lambda \in (0,1)$ such that

$$\varepsilon_0/\varepsilon \leq K_{4.6}/\lambda$$
,

we have

$$\mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le \frac{1}{\lambda}) \le 1 \le (K_{4.6}/\varepsilon_0)^{p-2} \lambda^{-(p-2)} \varepsilon^{p-2}.$$

The result follows by letting $c_{4.9}(\varepsilon_0) = 9eC_{4.6}\varepsilon_0^{-p}K_{4.6}^{p-2} + (K_{4.6}/\varepsilon_0)^{p-2}$.

For $|x| \geq \varepsilon_0$ and $\varepsilon \in (0, \varepsilon_0)$, we define

(4.17)
$$F_{\varepsilon,x}(K) \equiv \mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le K), \ \forall K > 0,$$

and

$$(4.18) \qquad \hat{F}_{\varepsilon,x}(\lambda) \equiv \mathbb{E}_{\delta_x} \Big(\exp\left(-\lambda \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2}\right) \mathbb{1}(X_{G_{\varepsilon}}(1) > 0) \Big), \ \forall \lambda > 0.$$

The dependence on ε and x will at times be suppressed.

LEMMA 4.10. There are constants $R_{4.10} > 2$ and $c_{4.10}(\varepsilon_0) > 0$ so that for any $\varepsilon_0 \leq |x| \leq \varepsilon_0^{-1}$ and $\varepsilon \in (0, \varepsilon_0/R_{4.10})$,

$$\mathbb{E}_{\delta_x} \left(\exp\left(-\lambda \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2}\right) 1(X_{G_{\varepsilon}}(1) > 0) \right) \ge c_{4.10}(\varepsilon_0) D^{\lambda}(R_{4.10}) \varepsilon^{p-2}, \ \forall \lambda > 0.$$

PROOF. By (1.25) and (1.26), for $\hat{F} = \hat{F}_{\varepsilon,x}$ defined as in (4.18) we have

(4.19)
$$\hat{F}(\lambda) = \exp(-U^{\lambda \varepsilon^{-2}, \varepsilon}(x)) - \exp(-U^{\infty, \varepsilon}(x))$$
$$\geq (U^{\infty, \varepsilon}(x) - U^{\lambda \varepsilon^{-2}, \varepsilon}(x)) \exp(-U^{\infty, \varepsilon}(x)).$$

Let $R_{4.10} = K_{4.3} + R_{4.6} > R_{4.6}$. Then for $\varepsilon \in (0, \varepsilon_0/R_{4.10})$ we have $|x|/\varepsilon > R_{4.10}$. Use Corollary 4.7(c) with $R = R_{4.10}$ to get

$$U^{\infty,\varepsilon}(x) - U^{\lambda\varepsilon^{-2},\varepsilon}(x) \ge c_{4.6}D^{\lambda}(R_{4.10})(R_{4.10}/|x|)^{p}\varepsilon^{p-2}$$

$$\ge c_{4.6}D^{\lambda}(R_{4.10})R_{4.10}^{p}\varepsilon_{0}^{-p}\varepsilon^{p-2}.$$

Next since $|x|/\varepsilon \ge \varepsilon_0/\varepsilon > R_{4.10} > K_{4.3}$, we may apply (4.7) to get

$$\exp(-U^{\infty,\varepsilon}(x)) = \exp(-\varepsilon^{-2}U^{\infty,1}(x/\varepsilon)) \ge \exp(-\varepsilon^{-2}3(4-d)(|x|/\varepsilon)^{-2})$$

$$\ge \exp(-9\varepsilon_0^{-2}).$$

So the lemma follows from (4.19) and the above inequalities.

PROPOSITION 4.11. There are positive constants $K_{4.11}(\varepsilon_0)$ and $c_{4.11}(\varepsilon_0)$ such that, for all $\varepsilon_0 \leq |x| \leq \varepsilon_0^{-1}$ we have

$$\mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le K_{4.11}(\varepsilon_0)) \ge c_{4.11}(\varepsilon_0)\varepsilon^{p-2}, \ \forall 0 < \varepsilon < \varepsilon_0/R_{4.10}.$$

PROOF. Recall $F = F_{\varepsilon,x}$ and $\hat{F} = \hat{F}_{\varepsilon,x}$ from (4.17) and (4.18), respectively. We have

$$\hat{F}(\lambda) = \int_0^\infty e^{-\lambda y} dF(y).$$

Let $\lambda=1$ and K>1. Use integration by parts and Proposition 4.9 to see that

$$\begin{split} \hat{F}(1) &= \int_0^\infty e^{-y} F(y) dy \le F(K) + \int_K^\infty e^{-y} F(y) dy \\ &\le F(K) + \int_K^\infty e^{-y} c_{4.9}(\varepsilon_0) y^{p-2} \varepsilon^{p-2} dy \le F(K) + \frac{1}{2} c_{4.10}(\varepsilon_0) D^1(R_{4.10}) \varepsilon^{p-2}, \end{split}$$

where $K = K(\varepsilon_0) > 1$ is large enough. Lemma 4.10, with $\lambda = 1$ and ε , x as in the Proposition, implies

$$F(K) \ge c_{4.10} D^1(R_{4.10}) \varepsilon^{p-2} - \frac{c_{4.10}}{2} D^1(R_{4.10}) \varepsilon^{p-2} = \frac{c_{4.10}}{2} D^1(R_{4.10}) \varepsilon^{p-2}.$$

Proof of Theorem 4.1. Pick $\lambda = \lambda(\varepsilon_0) \geq 2(4-d)$ large enough so that

$$e(2^p/|x|^p)D^{\lambda}(2) \le e2^p \varepsilon_0^{-p}D^{\lambda}(2) \le \frac{1}{2}c_{4.11}(\varepsilon_0).$$

So for $K_1(\varepsilon_0) \in (0, 1/\lambda(\varepsilon_0))$, Proposition 4.8 gives

$$\mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le K_1) \le \frac{1}{2} c_{4.11}(\varepsilon_0) \varepsilon^{p-2}.$$

Let $K_2 = K_{4.11}$ in Proposition 4.11 to see that for x, ε as in the Theorem and $R_{4.1} = R_{4.10}$,

$$\mathbb{P}_{\delta_x}(K_1 \le \frac{X_{G_{\varepsilon}(1)}}{\varepsilon^2} \le K_2) = \mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le K_2) - \mathbb{P}_{\delta_x}(0 < \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^2} \le K_1)$$

$$(4.20) \qquad \qquad \ge \frac{1}{2}c_{4.11}(\varepsilon_0)\varepsilon^{p-2}.$$

Use Proposition 2.3(b)(i) with $G = G_{\varepsilon}$ and $D_1 = G_{\varepsilon/2}$ to see that for x, ε

as above,

$$\mathbb{P}_{\delta_{x}}(K_{1} \leq \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^{2}} \leq K_{2}, X_{G_{\varepsilon/2}}(1) = 0)$$

$$= \mathbb{E}_{\delta_{x}}(1(K_{1} \leq \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^{2}} \leq K_{2}) \mathbb{P}_{X_{G_{\varepsilon}}}(X_{G_{\varepsilon/2}}(1) = 0))$$

$$= \mathbb{E}_{\delta_{x}}(1(K_{1} \leq \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^{2}} \leq K_{2}) \exp(-\int U^{\infty,\varepsilon/2}(y) X_{G_{\varepsilon}}(dy)))$$

$$= \mathbb{E}_{\delta_{x}}(1(K_{1} \leq \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^{2}} \leq K_{2}) \exp(-4\varepsilon^{-2}U^{\infty,1}(2) X_{G_{\varepsilon}}(1))) \text{ (by (1.24))}$$

$$\geq \mathbb{E}_{\delta_{x}}(1(K_{1} \leq \frac{X_{G_{\varepsilon}}(1)}{\varepsilon^{2}} \leq K_{2}) \exp(-4K_{2}U^{\infty,1}(2)))$$

$$\geq \frac{1}{2} c_{4.11}(\varepsilon_{0}) \varepsilon^{p-2} \exp(-4K_{2}U^{\infty,1}(2)),$$

the second equality by (1.26) with $X_0 = X_{G_{\varepsilon}}$ and the last inequality by (4.20). So the theorem follows.

5. Preliminaries for the Lower Bound on the Dimension. In this section, we will show that the lower bound on the local dimension of $\partial \mathcal{R}$ holds with positive probability (see Proposition 5.3). An extended version of this result, which is crucial for the later proof of Proposition 1.7, is given in Lemma 5.4. The next result is important for implementing our program: it plays a role analogous to that of Proposition 6.1 in [20].

PROPOSITION 5.1. There is a $\lambda > 0$ such that for all $\varepsilon_0 > 0$, there is some $c_{5.1}(\varepsilon_0) > 0$ so that for all $|x_i| \ge \varepsilon_0$ and $\varepsilon \in (0, \varepsilon_0)$,

$$\mathbb{E}_{\delta_0} \Big(\prod_{i=1}^2 \lambda \frac{X_{G_{\varepsilon}^{x_i}}(1)}{\varepsilon^2} \exp \Big(- \lambda \frac{X_{G_{\varepsilon}^{x_i}}(1)}{\varepsilon^2} \Big) \Big) \le c_{5.1} (1 + |x_1 - x_2|^{2-p}) \varepsilon^{2(p-2)}.$$

Given the results in Section 4, the proof then follows that of Proposition 6.1 in [20], but as there are some differences, it is included in Supplementary Material [13]. The reader should note that the role of λ in [20] is now played by $\lambda \varepsilon^{-2}$, where λ is chosen to be a fixed large constant.

Recall that we are in the case d=2 or 3. Let $\beta>0$ and $g_{\beta}(r)=r^{-\beta}$. For a finite measure μ on \mathbb{R}^d and Borel subset A of \mathbb{R}^d , let

$$\langle \mu \rangle_{g_{\beta}} = \int \int g_{\beta}(|x-y|) d\mu(x) d\mu(y),$$

and

$$I(g_{\beta})(A) = \inf\{\langle \mu \rangle_{g_{\beta}} : \mu \text{ is a probability supported by } A\}.$$

The g_{β} -capacity of A is $C(g_{\beta})(A) = (I(g_{\beta})(A))^{-1}$. Set

$$\beta = p - 2 = \begin{cases} 2\sqrt{2} - 2, & \text{if } d = 2, \\ (\sqrt{17} - 3)/2, & \text{if } d = 3, \end{cases}$$

and note $\beta \in (1/2, 1)$. Now we may use Theorem 4.1 and Proposition 5.1 to get the following theorem. Although similar to the omitted proof of Theorem 6.2 in [20], there are some important adjustments, and so this time the argument is included.

THEOREM 5.2. Assume d=2 or 3. For every $\varepsilon_0 \in (0,1)$, there is a $c_{5,2}(\varepsilon_0) > 0$ such that for any Borel set, A, of $\{x \in \mathbb{R}^d : \varepsilon_0 \leq |x| \leq \varepsilon_0^{-1}\}$,

$$\mathbb{P}_{\delta_0}(\partial \mathcal{R} \cap A \neq \emptyset) \ge c_{5.2}(\varepsilon_0)C(g_\beta)(A).$$

In particular, for any Borel subset A of \mathbb{R}^d , $C(g_\beta)(A) > 0$ implies that $\mathbb{P}_{\delta_0}(\partial \mathcal{R} \cap A \neq \emptyset) > 0$.

PROOF. Fix $\varepsilon_0 \in (0,1)$. Let Γ be a compact subset of $\{x \in \mathbb{R}^d : \varepsilon_0 \le |x| \le \varepsilon_0^{-1}\}$ such that (without loss of generality) $C(\Gamma) = C(g_\beta)(\Gamma) > 0$. If $I(\Gamma) = I(g_\beta)(\Gamma)$, we may choose $\{x_i^N : 1 \le i \le N\} \subset \Gamma$ so that (suppressing the superscript N) as $N \to \infty$ (see [25]),

(5.1)
$$I_N \equiv \frac{1}{N(N-1)} \sum_{i} \sum_{j \neq i} |x_i - x_j|^{-\beta} \to I(\Gamma) = 1/C(\Gamma).$$

Let

$$\hat{\Gamma} = \{x_i^N : 1 \le i \le N, N \ge 1\} \subset \Gamma.$$

For any fixed $\varepsilon \in (0, \varepsilon_0)$, we define $\partial \mathcal{R}(\varepsilon, \Gamma) \subset \Gamma$ to be

$$\partial \mathcal{R}(\varepsilon, \Gamma) := \{ x \in \hat{\Gamma} : K_1 \varepsilon^2 \le X_{G_{\varepsilon}^x}(1) \le K_2 \varepsilon^2, X_{G_{\varepsilon/2}^x}(1) = 0 \},$$

where $0 < K_1(\varepsilon_0) < K_2(\varepsilon_0) < \infty$ are as in Theorem 4.1. Note the exit measure $X_{G_{\varepsilon}^x}$ is constructed up to a null set for each (x,ε) , so the reader should note that we will only be considering at most countably many values of (x,ε) at a time as is the case in the above definition for a fixed ε . Let $\lambda > 0$ be as in Proposition 5.1. There is some $\delta = \delta(\lambda, \varepsilon_0) \in (0, e^{-1})$ such that

(5.2)
$$\lambda \frac{X_{G_{\varepsilon}^{x}}(1)}{\varepsilon^{2}} \exp\left(-\lambda \frac{X_{G_{\varepsilon}^{x}}(1)}{\varepsilon^{2}}\right) \geq \delta, \ \forall x \in \partial \mathcal{R}(\varepsilon, \Gamma).$$

Therefore by translation invariance, inclusion-exclusion, Theorem 4.1, (5.2), and Proposition 5.1, for $\varepsilon \in (0, \varepsilon_0/R_{4.1})$,

$$\mathbb{P}_{\delta_{0}}(\partial \mathcal{R}(\varepsilon, \Gamma) \neq \emptyset) \geq \mathbb{P}_{\delta_{0}}\left(\bigcup_{i=1}^{N} \{x_{i} \in \partial \mathcal{R}(\varepsilon, \Gamma)\}\right) \\
\geq \sum_{i=1}^{N} \mathbb{P}_{\delta_{0}}(x_{i} \in \partial \mathcal{R}(\varepsilon, \Gamma)) - \sum_{i} \sum_{j \neq i} \mathbb{P}_{\delta_{0}}(x_{i}, x_{j} \in \partial \mathcal{R}(\varepsilon, \Gamma)) \\
\geq Nc_{4.1}(\varepsilon_{0})\varepsilon^{p-2} - \sum_{i} \sum_{j \neq i} \delta^{-2} \mathbb{E}_{\delta_{0}}\left(\prod_{k=i,j} \lambda \frac{X_{G_{\varepsilon}^{x_{k}}}(1)}{\varepsilon^{2}} \exp\left(-\lambda \frac{X_{G_{\varepsilon}^{x_{k}}}(1)}{\varepsilon^{2}}\right)\right) \\
\geq Nc_{4.1}\varepsilon^{p-2} - c_{5.1}\delta^{-2}\varepsilon^{2(p-2)} \sum_{i} \sum_{j \neq i} (1 + |x_{i} - x_{j}|^{2-p}) \\
\geq c_{4.1}N\varepsilon^{p-2} - C(\varepsilon_{0})(N\varepsilon^{p-2})^{2}I_{N}.$$

Now choose $\varepsilon_N \to 0$ such that $N\varepsilon_N^{p-2} = c_{4.1}/(2C(\varepsilon_0)I_N)$. Therefore, for some $c(\varepsilon_0) > 0$,

$$\mathbb{P}_{\delta_0}(\partial \mathcal{R}(\varepsilon_N, \Gamma) \neq \emptyset) \geq \frac{c_{4.1}}{2C(\varepsilon_0)I_N} \frac{c_{4.1}}{2} \to c(\varepsilon_0)C(\Gamma), \text{ as } N \to \infty.$$

This implies

$$\mathbb{P}_{\delta_0}(\partial \mathcal{R}(\varepsilon_N, \Gamma) \neq \emptyset$$
, infinitely often) $\geq c(\varepsilon_0)C(\Gamma)$.

Assume now that

$$\omega \in \{\partial \mathcal{R}(\varepsilon_N, \Gamma) \neq \emptyset, \text{ infinitely often}\}.$$

So we may choose $\{z_N\} \subset \hat{\Gamma} \subset \Gamma$ such that $z_N \in \partial \mathcal{R}(\varepsilon_N, \Gamma)$, where we have suppressed the further subsequence of ε_N in our notation. The definition of $\partial \mathcal{R}(\varepsilon_N, \Gamma)$ gives $X_{G_{\varepsilon_N}^{z_N}}(1) > 0$ and $X_{G_{\varepsilon_N/2}^{z_N}}(1) = 0$. By Proposition 2.3(b)(ii) and translation invariance, $X_{G_{\varepsilon_N/2}^{z_N}}(1) = 0$ implies

(5.3)
$$\mathcal{R} \cap B(z_N, \varepsilon_N/3) = \emptyset.$$

By (2.3),
$$X_{G_{\varepsilon_N}^{z_N}}(1) > 0$$
 implies

(5.4)
$$\mathcal{R} \cap \partial G_{\varepsilon_N}^{z_N}$$
 is non-empty.

Combining (5.3) and (5.4) with an elementary argument in point set topology, we can choose $y_N \in \partial \mathcal{R}$ such that $\varepsilon_N/4 \leq |y_N - z_N| \leq \varepsilon_N$. The

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compactness of Γ implies there is some $x \in \Gamma$ such that $z_{N_k} \to x$ as $N_k \to \infty$. Therefore $y_{N_k} \to x$ and $x \in \partial \mathcal{R}$ since $\partial \mathcal{R}$ is closed, which gives $x \in \Gamma \cap \partial \mathcal{R} \neq \emptyset$, and so the proof is complete for $A = \Gamma$ compact. Use the inner regularity of capacity to extend the inequality to any Borel subset of $\{x \in \mathbb{R}^d : \varepsilon_0 \leq |x| \leq \varepsilon_0^{-1}\}$, and then conclude for any Borel subset A of \mathbb{R}^d that $C(g_\beta)(A) > 0$ implies that $\mathbb{P}_{\delta_0}(\partial \mathcal{R} \cap A \neq \emptyset) > 0$.

PROPOSITION 5.3. For each non-empty open set B in \mathbb{R}^d ,

- (a) $\mathbb{P}_{\delta_0}(\dim(\partial \mathcal{R} \cap B) \ge d_f) > 0$
- (b) $\mathbb{N}_0(\dim(\partial \mathcal{R} \cap B) \ge d_f) := p_{5,3}(B) > 0$

PROOF. (a) is derived from Theorem 5.2 by taking A to be the range of an appropriate independent Lévy process, exactly as in the proof of Proposition 6.5 of [20]. (b) then follows easily from (a) by making trivial changes to the cluster decomposition proof of Corollary 6.6 in [20].

To help upgrade the lower bound in part (a) of the above to probability one, we need to extend (a) to more general initial conditions through a scaling argument.

LEMMA 5.4. There is a constant $q_{5.4} > 0$ so that if $X'_0 \in M_F(\mathbb{R}^d)$ is supported on $\{|x| = r\}$ and $\delta = X'_0(1)$ satisfies $0 < \delta \le r^2$, then

$$\mathbb{P}_{X_0'}\Big(\dim\Big(\partial\mathcal{R}\cap B\Big(0,r-\frac{\sqrt{\delta}}{2}\Big)\Big)\geq d_f\Big)\geq q_{5.4}.$$

PROOF. Define $X_0^{(\delta)}(A) = \delta^{-1} X_0'(\sqrt{\delta}A)$, so that $X_0^{(\delta)}$ is supported on $\{|x| = r/\sqrt{\delta}\}$ and has total mass one. By scaling properties of SBM (see, e.g., Ex. II.5.5 in [21]) and scale invariance of Hausdorff dimension, we may conclude that

(5.5)
$$\mathbb{P}_{X_0'}(\dim(\partial \mathcal{R} \cap B\left(0, r - \frac{\sqrt{\delta}}{2}\right) \ge d_f) = \mathbb{P}_{X_0^{(\delta)}}(\dim(\partial \mathcal{R} \cap B\left(0, \frac{r}{\sqrt{\delta}} - \frac{1}{2}\right) \ge d_f).$$

Now work in our standard set-up for SBM with initial law $X_0^{(\delta)}$ so that (by (2.1)), $X_t = \sum_{j \in J} X_t(W_j) = \int X_t(W) \Xi(dW)$ for all t > 0, where Ξ is a

Poisson point process with intensity $\mathbb{N}_{X_0^{(\delta)}}$. For $r \geq \sqrt{\delta}$ define

$$\tau_{\rho}(W_{j}) = \inf\{t \geq 0 : |\hat{W}_{j}(t)| \leq \rho\},\$$

$$U_{\rho}(W_{j}) = \inf\{t \geq 0 : |\hat{W}_{j}(t) - \hat{W}_{j}(0)| \geq \rho\},\$$
and
$$N_{1} = \sum_{j \in J} 1(\tau_{(r/\sqrt{\delta}) - (1/2)}(W_{j}) < \infty) := \#(I_{r,\delta}).$$

Here as usual inf $\emptyset = \infty$. Then N_1 is Poisson with mean

(5.6)
$$m_{r,\delta} := \mathbb{N}_{X_0^{(\delta)}}(\tau_{(r/\sqrt{\delta})-(1/2)} < \infty) \le \mathbb{N}_{X_0^{(\delta)}}(U_{1/2}(W) < \infty)$$

= $\mathbb{N}_0(U_{1/2}(W) < \infty) := \overline{m} < \infty$,

where $X_0^{(\delta)}(1) = 1$ and translation invariance are used in the equality, and the finiteness of \bar{m} follows from Theorem 1 of [15].

Let $\mathcal{R}(W_j) = \{\hat{W}_j(t) : t \leq \sigma(W_j)\}$ (recall (2.2)) be the range of the jth excursion, so that

$$\mathcal{R} \cap B\left(0, \frac{r}{\sqrt{\delta}} - \frac{1}{2}\right) = \bigcup_{j \in J, \tau_{(r/\sqrt{\delta}) - (1/2)}(W_j) < \infty} \left(\mathcal{R}(W_j) \cap B\left(0, \frac{r}{\sqrt{\delta}} - \frac{1}{2}\right)\right).$$

We may assume (by additional randomization) that conditional on $I_{r,\delta}$, we have $\{W_j: j \in I_{r,\delta}\}$ are iid with law $\mathbb{N}_{X_0^{(\delta)}}(W \in \cdot | \tau_{(r/\sqrt{\delta})-(1/2)} < \infty)$. Therefore the right-hand side of (5.5) is at least

$$(5.7) \quad \mathbb{P}_{X_0^{(\delta)}}(N_1 = 1) \mathbb{N}_{X_0^{(\delta)}} \left(\dim \left(\partial \mathcal{R} \cap B(0, \frac{r}{\sqrt{\delta}} - \frac{1}{2}) \right) \ge d_f \middle| \tau_{\frac{r}{\sqrt{\delta}} - \frac{1}{2}} < \infty \right)$$

$$= \frac{m_{r,\delta} e^{-m_{r,\delta}}}{m_{r,\delta}} \mathbb{N}_{x_0} \left(\dim \left(\partial \mathcal{R} \cap B(0, \frac{r}{\sqrt{\delta}} - \frac{1}{2}) \right) \ge d_f \right),$$

where $x_0 = (\frac{r}{\sqrt{\delta}})e_1$ and e_1 is the first unit basis vector. We also have used the fact that spherical symmetry shows we could have taken any x_0 on the sphere of radius $r/\sqrt{\delta}$. Now again use translation invariance and spherical symmetry to see that the right side of (5.7) equals

$$e^{-m_{r,\delta}} \mathbb{N}_0 \left(\dim \left(\partial \mathcal{R} \cap B(x_0, |x_0| - \frac{1}{2}) \right) \ge d_f \right)$$

$$\ge e^{-m_{r,\delta}} \mathbb{N}_0 \left(\dim \left(\partial \mathcal{R} \cap B(e_1, 1/2) \right) \ge d_f \right)$$

$$\ge e^{-\overline{m}} \mathbb{N}_0 \left(\dim \left(\partial \mathcal{R} \cap B(e_1, 1/2) \right) \ge d_f \right)$$

$$\ge e^{-\overline{m}} p_{5.3} \left(B(e_1, 1/2) \right) > 0,$$
(5.8)

where the first inequality follows by $B(e_1, 1/2) \subset B(x_0, |x_0| - 1/2)$, the second inequality uses (5.6) and the first inequality in the last line holds by Proposition 5.3(b). We have shown that the right-hand side of (5.8) is a lower bound for (5.5), and so have proved the lemma with $q_{5.4} = e^{-\overline{m}}p_{5.3}(B(e_1, 1/2))$.

6. Exit Measures and Continuous State Branching Processes. Throughout this section we assume that $r_0 > 0$ and the initial state X_0 satisfy

$$(6.1) B_{2r_0} \subset \operatorname{Supp}(X_0)^c.$$

To finish the proofs of Propositions 1.6, 1.7 we need to establish some properties of the total exit measure process $X_{Gr_0-r}(1), 0 \le r < r_0$. We will show in Proposition 6.2 that, for any $r_0 > 0$, the "time changed" process $Z_t = X_{G_{r_0e^{-t}}}(1)/(r_0e^{-t})^2, t \ge 0$, is a time homogeneous continuous state branching process (CSBP) and thus has no negative jumps.

A non-negative function $\lambda \mapsto u(\lambda)$ on $[0, \infty)$ is completely concave iff for every y > 0 $\lambda \to \exp(-yu(\lambda))$ is the Laplace transform of a probability measure on the half-line. (See (4.1) in [23] for a Lévy-Khintchine representation of such functions). We recall the definition of a continuous state branching process from Section 4 of [23].

Definition. A (finite) continuous state branching process (CSBP) Z is a time-homogeneous $[0, \infty)$ -valued Markov process with no fixed time discontinuities (if $t_n \to t$, then $Z(t_n) \to Z(t)$ a.s.), and such that there is a family of completely concave functions $\{u(s, \cdot) : s > 0\}$ satisfying

(6.2)
$$E(\exp(-\lambda Z(t_2))|Z(s), s \le t_1)$$

$$= \exp(-Z(t_1)u(t_2 - t_1, \lambda)) \text{ a.s. } \forall t_2 > t_1 \ge 0, \lambda \ge 0.$$

The associated family $\{u(t,\cdot): t>0\}$ is the log-Laplace transform of Z. Recall that $U^{\lambda,R}(x)=U^{\lambda,R}(|x|)$ is the unique continuous map on $\{|x|\geq R\}$ which is C^2 on G_R and satisfies

(6.3)
$$\Delta U = U^2 \text{ on } G_R \text{ and } U = \lambda \text{ on } \partial G_R.$$

A simple application of the comparison principle (e.g., Chapter V, Lemma 7 of [17]), using the last part of (1.27), gives

(6.4)
$$U^{\lambda,R}(x) \le \lambda \quad \forall |x| \ge R.$$

Define

(6.5)
$$u(t,\lambda) = e^{2t} U^{\lambda,1}(e^t) \text{ for } t \ge 0.$$

Notation. We define $Y(r) = X_{G_{r_0-r}}$ and $\mathcal{E}_r = \mathcal{E}_{G_{r_0-r}} \vee \{\mathbb{N}_{X_0} - \text{null sets}\}$ for $0 \le r < r_0$, and for $t \ge 0$ set

$$Z(t) = X_{G_{r_0e^{-t}}}(1)\frac{e^{2t}}{r_0^2} = Y(r_0(1-e^{-t}))(1)e^{2t}r_0^{-2}; \quad \mathcal{G}_t = \mathcal{E}_{r_0(1-e^{-t})} = \mathcal{E}_{G_{r_0e^{-t}}}.$$

It is not hard to show that \mathcal{E}_r is non-decreasing in r (the corresponding result for half-spaces is noted prior to (7.2) of [20] and the observation made there applies to balls as well.) By Proposition 2.3 of [16], Y is (\mathcal{E}_r) -adapted and Z is (\mathcal{G}_t) -adapted. Let $\mathcal{E}_r^+ = \mathcal{E}_{r+}$ denote the associated right-continuous filtration. In addition to \mathbb{N}_{X_0} , we will also work under the probability $Q_{X_0}(\cdot) = \mathbb{N}_{X_0}(\cdot|Y_0(1) > 0)$, where (6.1) ensures that $\mathbb{N}_{X_0}(Y_0(1) > 0) < \infty$. Note that

(6.6) for any r.v.
$$Z \ge 0$$
, and any $r \ge 0$, $Q_{X_0}(Z|\mathcal{E}_r) = \mathbb{N}_{X_0}(Z|\mathcal{E}_r) Q_{X_0}$ -a.s.

because $\{Y_0(1) > 0\} \in \mathcal{E}_0$. When conditioning on \mathcal{E}_r under Q_{X_0} , we are adding the slightly larger class of Q_{X_0} -null sets to \mathcal{E}_r , but will not record this distinction in our notation. Below we will apply the definition of (CSBP) under the σ -finite measure \mathbb{N}_{X_0} as well as Q_{X_0} . We write Q_{x_0} for $Q_{\delta_{x_0}}$ as usual.

LEMMA 6.1. (a) If $0 \le t_1 < t_2$ and $\lambda \ge 0$, then

(i)
$$Q_{X_0}\left(e^{-\lambda Z_{t_2}}\middle|\mathcal{G}_{t_1}\right) = \mathbb{N}_{X_0}\left(e^{-\lambda Z_{t_2}}\middle|\mathcal{G}_{t_1}\right) = \exp(-Z_{t_1}u(t_2 - t_1, \lambda)).$$

(ii)

$$Q_{X_0} \left(\left(e^{-\lambda Z_{t_2}} - e^{-\lambda Z_{t_1}} \right)^2 \right) = Q_{X_0} \left(\exp(-Z_{t_1} u(t_2 - t_1, 2\lambda)) - 2 \exp(-\lambda Z_{t_1} - Z_{t_1} u(t_2 - t_1, \lambda)) + \exp(-2\lambda Z_{t_1}) \right),$$
(6.7)

and similarly for \mathbb{N}_{X_0} .

- (b) For all t > 0, $\lambda \mapsto u(t, \lambda)$ is completely concave.
- (c) $(Z_t, t \geq 0)$ is a (time-homogeneous) (\mathcal{G}_t) -Markov process under Q_{X_0} or \mathbb{N}_{X_0} .

PROOF. (a) (6.6) shows that for $\lambda \geq 0$, the left-hand side of (i) equals the middle expression, which by Proposition 2.2(a)(ii) and then (1.25) equals

$$\mathbb{E}_{X_{G_{r_0e}-t_1}}(\exp(-\lambda e^{2t_2}r_0^{-2}X_{G_{r_0e}-t_2}(1)))$$

$$= \exp\left(-\int U^{\lambda e^{2t_2}r_0^{-2},r_0e^{-t_2}}(x)X_{G_{r_0e}-t_1}(dx)\right)$$

$$= \exp\left(-U^{\lambda e^{2t_2}r_0^{-2},r_0e^{-t_2}}(r_0e^{-t_1})X_{G_{r_0e}-t_1}(1)\right)$$

$$= \exp(-u(t_2-t_1,\lambda)Z_{t_1}),$$

where scaling (i.e., (1.24)) is used in the last line. This gives (i). It is then easy to derive (ii) by expanding out the square, conditioning on \mathcal{G}_{t_1} and finally using (i).

(b) Let $y_0 > 0$ and t > 0. Let m_r be the uniform distribution on $\{|x| = r\}$ and set $W = e^{2t} r_0^{-2} X_{G_{r_0e^{-t}}}(1)$. Apply (1.25) and then scaling ((1.24)) to see that for all $\lambda \geq 0$,

$$\mathbb{E}_{y_0 r_0^2 m_{r_0}}(\exp(-\lambda W)) = \exp(-y_0 r_0^2 U^{\lambda e^{2t} r_0^{-2}, r_0 e^{-t}}(r_0))$$

$$= \exp(-y_0 r_0^2 r_0^{-2} e^{2t} U^{\lambda, 1}(e^t))$$

$$= \exp(-y_0 u(t, \lambda)).$$

(c) This is immediate from (a)(i), (b) (to define the family of laws $\{P_x : x \ge 0\}$), and a monotone class argument.

Proposition 1.8 in the Introduction follows from part (c) of our next result.

PROPOSITION 6.2. (a) Y is an inhomogeneous (\mathcal{E}_r) -Markov process under \mathbb{N}_{X_0} or Q_{X_0} . That is, for $\psi: M_F(\mathbb{R}^d) \to [0,\infty)$ Borel measurable and $0 \le r_1 < r_2$,

$$Q_{X_0}(\psi(Y(r_2))|\mathcal{E}_{r_1}) = \mathbb{N}_{X_0}(\psi(Y(r_2))|\mathcal{E}_{r_1}) = \mathbb{E}_{Y(r_1)}(\psi(Y(r_2)))$$
 a.s.

(b) If $0 \le r_1 < r_2 < r_0$, then the total mass, $Y_r(1)$, of Y_r satisfies

(6.8)
$$\mathbb{N}_{X_0}(Y_{r_2}(1)|\mathcal{E}_{r_1}) = \begin{cases} Y_{r_1}(1) & \text{if } d = 2\\ \frac{r_0 - r_2}{r_0 - r_1} Y_{r_1}(1) & \text{if } d = 3. \end{cases}$$

Under \mathbb{N}_{X_0} or Q_{X_0} , $Y_r(1)$ has a cadlag version on $[0, r_0)$ which is an (\mathcal{E}_r^+) -supermartingale (an (\mathcal{E}_r^+) -martingale if d=2), satisfies (6.8) with $\mathcal{E}_{r_1}^+$ in place of \mathcal{E}_{r_1} , and has only non-negative jumps a.e.

(c) Under \mathbb{N}_{X_0} or Q_{X_0} , $Z(t), t \geq 0$ has a caddag version with only non-negative jumps which is a CSBP with log Laplace transform given by $\{u(t,\cdot): t>0\}$ in (6.5).

PROOF. (a) This is immediate from Proposition 2.2(a)(ii) and (6.6). (b,c) Let B denote a d-dimensional Brownian motion starting at x under P_x^B and

$$\tau_r = \inf\{t \ge 0 : |B_t| \le r\} \quad (\inf \emptyset = \infty).$$

Recalling (6.1), Proposition 3 in Chapter V of [17] shows that for $0 \le r < r_0$, (6.9)

$$\mathbb{N}_{X_0}(Y_r(1)) = \int P_x^B(\tau_{r_0-r} < \infty) dX_0(x) = \begin{cases} X_0(1) & \text{if } d = 2\\ \int \frac{r_0-r}{|x|} dX_0(x) & \text{if } d = 3. \end{cases}$$

Return now to the probability \mathbb{P}_{X_0} , and use (2.4) and the above to see that

(6.10)
$$\mathbb{E}_{X_0}(Y_r(1)) = \mathbb{N}_{X_0}(Y_r(1)) = \begin{cases} X_0(1) & \text{if } d = 2\\ \int \frac{r_0 - r}{|x|} dX_0(x) & \text{if } d = 3. \end{cases}$$

Although we have assumed $B_{2r_0} \subset \operatorname{Supp}(X_0)^c$, both (6.9) and (6.10) will apply if $\operatorname{Supp}(X_0) \subset G_{r_0-r}$. This allows us to apply (6.10), with $X_0 = Y_{r_1}$ and $r = r_2$, and (a) to derive (6.8).

Turning to the second part of (b) and (c) we first work with Z. Let $t_n \uparrow t > 0$ ($t_n < t$) and set $r_n = r_0(1 - e^{-t_n}) \uparrow r_0(1 - e^{-t}) = r \in (0, r_0)$. By (6.8) and supermartingale convergence, $\{Y_{r_n}\}$ converges \mathbb{N}_{X_0} -a.e. to a limit we denote by $Y_{r_n}(1)$ for now. (The σ -finiteness of \mathbb{N}_{X_0} is not an issue here, but the reader who prefers probabilities may work with Q_{X_0} and note that on the complementary set, $\{Y_0(1) = 0\}$, $Y_{r_n}(1) = 0$ \mathbb{N}_{X_0} -a.e. by (6.8) with $r_1 = 0$ there. Henceforth we will not make such arguments.) It follows that

(6.11)
$$Z_{t_n} \to e^{2t} r_0^{-2} Y_{r-}(1) := Z_{t-} \quad \mathbb{N}_{X_0} - \text{a.e.}$$

By (6.7).

$$Q_{X_0}\left(\left(e^{-\lambda Z_t} - e^{-\lambda Z_{t_n}}\right)^2\right)$$

$$= Q_{X_0}\left(\exp(-Z_{t_n}u(t - t_n, 2\lambda)) - 2\exp(-(\lambda + u(t - t_n, \lambda))Z_{t_n})\right)$$

$$+ \exp(-2\lambda Z_{t_n})$$

$$\to Q_{X_0}\left(\exp(-2\lambda Z_{t_n}) - 2\exp(-2\lambda Z_{t_n}) + \exp(-2\lambda Z_{t_n})\right) \text{ as } n \to \infty$$

$$= 0,$$

where Dominated Convergence is used in the above convergence. This and (6.11) show that $Z_{t_n} \to Z_t \ Q_{X_0}$ -a.s. The fact, noted above, that $Y_0(1) = 0$ implies $Z_{t_n} = Z_t = 0 \ \mathbb{N}_{X_0}$ -a.e. allows us to upgrade this to

(6.12)
$$Z_{t_n} \to Z_t \quad \mathbb{N}_{X_0} - a.e. \quad \text{if } t_n \uparrow t > 0.$$

A simpler argument, now using reverse supermartingale convergence, shows that

(6.13)
$$Z_{t_n} \to Z_t \quad \mathbb{N}_{X_0} - a.e. \quad \text{if } t_n \downarrow t \geq 0.$$

(6.12) and (6.13) imply $Y_r(1)$ is continuous in measure on $[0, r_0)$. Therefore by (6.8) there is a cadlag version of $(Y_r(1), r \in [0, r_0))$ under \mathbb{N}_{X_0} (we do not change the notation) which is an (\mathcal{E}_r^+) -supermartingale (martingale if d=2) satisfying (6.8) with $\mathcal{E}_{r_1}^+$ in place of \mathcal{E}_{r_1} . This gives a cadlag version of Z which satisfies the (\mathcal{G}_{t+}) version of Lemma 6.1(a)(i), and so is (\mathcal{G}_{t+}) -Markov under \mathbb{N}_{X_0} or Q_{X_0} , just as for Lemma 6.1(c). Clearly (6.12) and (6.13) imply that $Z_{t-} = Z_t \ \mathbb{N}_{X_0}$ -a.e., and so Z has no fixed time discontinuities. It follows from the above and Lemma 6.1(b) that $(Z_t, t \geq 0)$ is a (CSBP) with log Laplace transform $\{u(t, \cdot) : t > 0\}$ under \mathbb{N}_{X_0} or Q_{X_0} . A theorem of Lamperti (see. e.g. p. 1044 of [23]) shows that $(Z_t, t \geq 0)$ has only non-negative jumps a.e. and so the same applies to $(Y_r(1), r \in [0, r_0])$.

REMARK 6.3. Although in this work we only use the above results, we briefly discuss the processes Z. and Y.(1) in the general context of CSBP's. By Proposition 6.2(c) above and Theorem 4 of [23] there is a Lévy measure $\tilde{\pi}$ on $[0,\infty)$ satisfying $\int \ell^2 \wedge 1 \, d\tilde{\pi}(\ell) < \infty$ and constants $\tilde{a} \in \mathbb{R}, b \geq 0$, such that if

(6.14)
$$\Psi(u) = \tilde{a}u - bu^2 + \int_0^\infty (1 - e^{-u\ell} - u\ell e^{-\ell}) \, d\tilde{\pi}(\ell), \quad u \ge 0,$$

then $t \mapsto u(t,\lambda)$ is the unique solution of

(6.15)
$$\frac{du(t,\lambda)}{dt} = \Psi(u(t,\lambda)), \quad u(0,\lambda) = \lambda.$$

Z is often called a Ψ -CSBP. (6.14) implies Ψ is concave on $[0,\infty)$ and differentiable on $(0,\infty)$. If

$$\lambda_d = 2(4-d),$$

then a short calculation using (6.5) and (6.3) gives (primes denote derivatives with respect to t)

(6.16)
$$u''(t,\lambda) = (6-d)u' + u(u-\lambda_d), \ t \ge 0.$$

Differentiating both sides of (6.15) and using (6.16) on the resulting left-hand side, leads to the first order ode for Ψ ,

$$\Psi'\Psi(u(t,\lambda)) = (6-d)\Psi(u(t,\lambda)) + u(t,\lambda)(u(t,\lambda) - \lambda_d), \quad \Psi(0) = 0.$$

Letting $t \to 0$ and varying λ we conclude that Ψ is a solution of the ode

(6.17)
$$\Psi'\Psi(u) = (6-d)\Psi(u) + u(u-\lambda_d), \ u > 0, \quad \Psi(0) = 0.$$

By using this equation to analyze the behaviour of Ψ near ∞ it is easy to see that in (6.14), b=0. The concavity of Ψ implies $\lim_{u\downarrow 0} \frac{\Psi(u)}{u}=\lim_{u\downarrow 0} \Psi'(u)\in (-\infty,+\infty]$. If we divide both sides of (6.17) by u and let $u\downarrow 0$ we conclude this limit, $\Psi'(0)$ is in fact finite and satisfies

$$\Psi'(0)^2 = (6 - d)\Psi'(0) - \lambda_d,$$

that is, $\Psi'(0) = 2$ if d = 2, and $\Psi'(0) = 1$ or 2 if d = 3. It is not hard to see using (6.8) that, in fact, $\Psi'(0) = 1$ if d = 3. The fact that this derivative is finite, already implies that $\int_0^\infty \ell d\tilde{\pi}(\ell) < \infty$ and (6.14) can be rewritten as

$$(6.18) \quad \Psi(u) = a_d u + \int_0^\infty (1 - e^{-u\ell} - u\ell) d\pi(\ell), \quad \int_0^\infty \ell \wedge \ell^2 d\pi(\ell) < \infty,$$

where now $a_d = \Psi'(0) = 4 - d$, by the above. The ode (6.17) can be used to study the tail behaviour of Ψ , and hence π , via Tauberian theorems. For example it is not hard to show that for some explicit $c_{6.19} > 0$,

(6.19)
$$\lim_{\varepsilon \downarrow 0} \varepsilon^{3/2} \pi([\varepsilon, \infty)) = c_{6.19}.$$

The process of total mass of the exit measure from B_r (as opposed to G_{r_0-r}) is studied in [10] as an inhomogeneous CSBP. The setting there is for general branching mechanisms, but the ideas used above and in defining Z appear to be novel. It would be of interest to study the detailed behaviour of the measure-valued process $r \to X_{G_{r_0-r}}$.

In [20] we instead worked with the exit measure from half spaces $H_r = \{x: x_1 < r\}$, where the total mass process is a Ψ -CSBP with $\Psi(u) = \frac{\sqrt{6}}{3}u^{3/2}$ (see [10] and Proposition 4.1 of [20]; it is interesting to note that the same 3/2 stable branching mechanism appears in [18] for subordinate trees). The CSBP analysis in [20] was simpler due to this explicit 3/2-stable Ψ , but halfplanes were clumsier and led to less precise results. See the discussion at the end of the Introduction. From the perspective of small jumps, the balls look like half-spaces and so (6.19)) is to be expected. We hasten to add, however, that $\Psi(u)$ has behaviour near u=0 (corresponding to π at ∞) which is quite different and dimension dependent. Again using (6.17) one can show that $\forall \varepsilon > 0$,

(6.20)
$$u^{1+\varepsilon} \ll a_d u - \Psi(u) \ll u \quad \text{as } u \downarrow 0 \text{ if } d = 2,$$
$$u^2 \ll a_d u - \Psi(u) \ll u^{2-\varepsilon} \text{ as } u \downarrow 0 \text{ if } d = 3.$$

7. Proof of Propositions 1.6, 1.7. We use the notation from Section 6. In particular X_0 and $r_0 > 0$ are as in (6.1), $Y_r = X_{G_{r_0-r}}$ for $0 \le r < r_0$, and $Q_{X_0}(\cdot) = \mathbb{N}_{X_0}(\cdot|Y_0(1) > 0)$.

In what follows we always will work with the cadlag versions of $Y_r(1)$, and hence Z_t , constructed in Proposition 6.2(b) above. We let W denote a generic snake under \mathbb{N}_{X_0} or Q_{X_0} with the associated "tip process" $\hat{W}(t)$ and excursion length σ . Define

$$T_0(W) = \inf\{r \in [0, r_0) : Y_r(1) = 0\} \in [0, r_0], \text{ where } \inf \emptyset = r_0,$$

and

(7.1)
$$\hat{T}_0(W) = \inf\{|\hat{W}(t)| : 0 \le t \le \sigma\} = \inf\{|x| : x \in \mathcal{R}\},\$$

the final equality holding \mathbb{N}_{X_0} -a.e. by (2.2). Clearly we have

$$Q_{X_0}(\cdot) = \mathbb{N}_{X_0}(\cdot | T_0 > 0).$$

LEMMA 7.1. The sets $\{T_0 > 0\}$ and $\{\hat{T}_0 < r_0\}$ coincide \mathbb{N}_{X_0} -a.e., and on this set, $\hat{T}_0 = r_0 - T_0 \, \mathbb{N}_{X_0}$ -a.e.

PROOF. For every rational q in $[0, T_0)$, $X_{G_{r_0-q}}(1) > 0$ implies $\partial G_{r_0-q} \cap \mathcal{R}$ is non-empty (by (2.3)) and so by (2.2) $\hat{T}_0 \leq r_0 - q$. This proves that

(7.2)
$$\hat{T}_0 \le r_0 - T_0 \quad \mathbb{N}_{X_0}$$
-a.e. on $\{T_0 > 0\}$.

Conversely assume $r_0 > T_0$ and choose rationals q, q' so that $T_0 < q' < q < r_0$. Then $X_{G_{r_0-q'}}(1) = 0$ and the special Markov property (Proposition 2.2(b)) at $R_1 = r_0 - q'$ shows that $\mathbb{N}_{X_0}(\mathcal{R} \cap B_{r_0-q} \neq \emptyset | \mathcal{E}_{q'}) = 0$ a.e on $\{T_0 < q'\}$. This proves that

(7.3)
$$\hat{T}_0 \ge r_0 - T_0 \quad \mathbb{N}_{X_0}$$
-a.e. on $\{T_0 < r_0\}$.

The above is trivial if $T_0 = r_0$ and so we have shown (by (7.2) and (7.3))

$$\hat{T}_0 = r_0 - T_0$$
 \mathbb{N}_{X_0} -a.e on $\{T_0 > 0\}$.

Finally, note that (7.2) shows $T_0 > 0$ implies $\hat{T}_0 < r_0$, and (7.3) shows $T_0 = 0$ implies $\hat{T}_0 \ge r_0$, which in turn shows $\hat{T}_0 < r_0$ implies $T_0 > 0$ (all up to \mathbb{N}_{X_0} null sets). This proves the a.e. equality of $\{T_0 > 0\}$ and $\{\hat{T}_0 < r_0\}$, and completes the proof.

LEMMA 7.2. (a) For $0 < r < r_0$,

$$\mathbb{N}_{X_0}(0 < T_0 \le r) = \mathbb{N}_{X_0}(r_0 - r \le \hat{T}_0 < r_0)$$

$$= \mathbb{N}_{X_0} \Big(1(X_{G_{r_0}}(1) > 0) \exp[-X_{G_{r_0}}(1)(r_0 - r)^2 U^{\infty,1}(r_0/(r_0 - r))] \Big).$$

(b)
$$\mathbb{N}_{X_0}(T_0 \in dr) \ll dr$$
 on $\{0 < r < r_0\}$ and $\mathbb{N}_{X_0}(\hat{T}_0 \in dr) \ll dr$ on $\{0 < r < 2r_0\}$.

PROOF. (a) Using (1.25) and scaling ((1.24) with $\lambda = \infty$), we have for $0 < r < r_0$,

$$\mathbb{P}_{X_{G_{r_0}}}(X_{G_{r_0-r}}(1)=0) = \exp(-X_{G_{r_0}}(U^{\infty,r_0-r}))$$

$$= \exp(-X_{G_{r_0}}(1)(r_0-r)^{-2}U^{\infty,1}(r_0/(r_0-r))).$$

The special Markov property (Proposition 2.2(a)(ii)) shows for $0 < r < r_0$,

$$\mathbb{N}_{X_0}(0 < T_0 \le r) = \mathbb{N}_{X_0}(1(T_0 > 0)\mathbb{N}_{X_0}(X_{G_{r_0-r}}(1) = 0|\mathcal{E}_0))
= \mathbb{N}_{X_0}(1(T_0 > 0)\mathbb{P}_{X_{G_{r_0}}}(X_{G_{r_0-r}}(1) = 0))
= \mathbb{N}_{X_0}\Big(1(X_{G_{r_0}}(1) > 0)\exp[-X_{G_{r_0}}(1)(r_0 - r)^{-2}U^{\infty,1}(r_0/(r_0 - r))]\Big),$$

where (7.4) has been used in the last line. This, together with Lemma 7.1, gives (a).

(b) The right-hand side of (a) is continuously differentiable in $r \in (0, r_0)$ because $U^{\infty,1}$ is C^2 on G_1 (recall (1.27)). Here we note that it is easy to justify differentiation inside the integral since $\mathbb{N}_{X_0}(X_{G_{r_0}}(1) > 0) < \infty$ (recall (6.1)), $\mathbb{N}_{X_0}(X_{G_{r_0}}(1)) < \infty$ (recall (6.9)), and $(U^{\infty,1})'(r)$ is bounded on compacts away from $\{r \leq 1\}$. This gives the first part of (b). Lemma 7.1 now implies the absolute continuity of $\mathbb{N}_{X_0}(\hat{T}_0 \in dr)$ on $\{0 < r < r_0\}$. But (6.1) allows us to replace r_0 with αr_0 for any $1 < \alpha < 2$ in the above reasoning and so conclude that $\mathbb{N}_{X_0}(\hat{T}_0 \in dr)$ is absolutely continuous on $\{0 < r < 2r_0\}$.

Proof of Proposition 1.6 assuming Proposition 1.7. By translation invariance we may assume $x_1 = 0$. Fix r_0, r_1 and X_0 as in our hypotheses. We must show that

(7.5)
$$X_{G_{r_1}}(1) = 0$$
 and $X_{G_{r_0}}(1) > 0$ imply $\dim(B_{r_0} \cap \partial \mathcal{R}) \ge d_f \mathbb{P}_{X_0} - a.s.$

Measurability issues are easily handled using Lemma 2.1 and will henceforth be ignored. We work under \mathbb{P}_{X_0} in the standard set-up and so from (2.4) have for $0 < r \le r_0$ and $J_0 = \{j \in J : \hat{T}_0(W_j) \le r_0\}$,

$$X_{G_r} = \sum_{j \in J} X_{G_r}(W_j) = \sum_{j \in J} X_{G_r}(W_j) 1(\hat{T}_0(W_j) \le r) = \sum_{j \in J_0} X_{G_r}(W_j) 1(\hat{T}_0(W_j) \le r).$$

Here we used the fact that $\hat{T}_0(W_j) > r$ implies $X_{G_r}(W_j) = 0$ (e.g. by (2.3) and (2.2)). Recall from (2.2) that the range of the *j*th excursion W_j is

$$\mathcal{R}_j := \mathcal{R}(W_j) = \{ \hat{W}_j(t) : t \le \sigma(W_j) \}.$$

It follows easily from (2.1) (see (2.19) in [20]) that for $x \in \overline{B_{r_0}}$,

$$L^x = \sum_{j \in J_0} L^x(W_j),$$

and therefore,

(7.7)

$$\mathcal{R} \cap B_{r_0} = \bigcup_{j \in J_0} (\mathcal{R}_j \cap B_{r_0}) \text{ and so } 1_{\mathcal{R} \cap B_{r_0}}(x) = 1 \Big(\sum_{j \in J_0} 1_{\mathcal{R}(W_j) \cap B_{r_0}}(x) > 0 \Big).$$

We will frequently use the elementary topological result

(7.8)
$$B_{r_0} \cap \partial F = B_{r_0} \cap \partial (B_{r_0} \cap F) = B_{r_0} \cap \partial (\overline{B_{r_0}} \cap F)$$
 for any closed set F .

It follows easily from Lemma 7.1 that (7.9)

$$T_0 = r_0 - \bigwedge_{i \in J_0} \hat{T}_0(W_i)$$
 on $\{T_0 > 0\} = \{\bigwedge_{i \in J_0} \hat{T}_0(W_i) < r_0\} \mathbb{P}_{X_0} - \text{a.s.}$.

In view of the absolute continuity properties of \hat{T}_0 under \mathbb{N}_0 from Lemma 7.2 we see from the above that if $N_0 = |J_0|$, a Poisson mean $\mathbb{N}_{X_0}(\hat{T}_0 \leq r_0)$ random variable, then

(7.10)
$$T_0 > 0 \text{ iff } J_0 \neq \emptyset \text{ iff } N_0 > 0 \mathbb{P}_{X_0} - \text{a.s.}.$$

By enlarging our probability space and randomizing the above Poisson points we may assume that there is an iid sequence $\{\widetilde{W}_j: j \in \mathbb{N}\}$, independent of the Poisson variable $N_0 = |J_0|$ with mean $\mathbb{N}_{X_0}(\hat{T}_0 \leq r_0)$, and with common law

(7.11)
$$\mathbb{N}_{X_0}(\cdot | \hat{T}_0 \le r_0) = \mathbb{N}_{X_0}(\cdot | \hat{T}_0 < r_0),$$

(the last equality by Lemma 7.2) and so that

(7.12)
$$\sum_{j \in J_0} \delta_{W_j} = \sum_{j=1}^{N_0} \delta_{\widetilde{W}_j}.$$

Let $\widehat{\widehat{W}}_j$ denote the tip of the jth excursion and define

$$\hat{T}^j = \hat{T}_0(\widetilde{W}_j) < r_0$$
 (a.s. by (7.11)),

$$\widetilde{\mathcal{R}}_{i} = \mathcal{R}(\widetilde{W}_{i}) = \{\widehat{\widetilde{W}}_{i}(t) : t \leq \sigma(\widetilde{W}_{i})\}.$$

Note that $X_{G_{r_1}}(1) = 0$ implies $T_0 < r_0 - r_1$ a.s. and so, in view of (7.9), (7.13)

$$X_{G_{r_1}}(1) = 0$$
 and $Y_0(1) = X_{G_{r_0}}(1) > 0$ imply $\hat{T}^j \ge r_1 \ \forall \ j \le N_0 \ \mathbb{P}_{X_0}$ - a.s.

The independence of the \hat{T}^j 's and fact they have no positive atoms by Lemma 7.2 imply

(7.14)

$$\mathbb{P}_{X_0}(\hat{T}^j = \hat{T}^{j'} \text{ for some } 1 \le j \ne j' \le N_0, X_{G_{r_1}}(1) = 0, X_{G_{r_0}}(1) > 0) = 0.$$

So on $\{X_{G_{r_1}}(1) = 0, X_{G_{r_0}}(1) > 0\}$ there is an a.s. unique $\tilde{j} \leq N_0$ s.t. $\hat{T}^{\tilde{j}} = \min\{\hat{T}^j : j \leq N_0\}$. (7.14) and (7.13) imply that (if an empty minimum is r_0) $\mathbb{P}_{X_0} - \text{a.s.}$,

$$\tilde{T} := \min\{\hat{T}^j : j \neq \tilde{j}, j \leq N_0\} > \hat{T}^{\tilde{j}} \geq r_1$$

$$\text{on } \{X_{G_{r_1}}(1) = 0, X_{G_{r_0}}(1) > 0\} \subset \{N_0 \geq 1\}.$$

Hence \widetilde{T} is the largest radius $r \leq r_0$ so that a single excursion \widetilde{W}_j enters B_r (it exists on $\{X_{G_{r_1}}(1) = 0, X_{G_{r_0}}(1) > 0\}$).

By the definition of \tilde{j} and \tilde{T} we have from (7.15) and (7.7),

$$B_{\tilde{T}} \cap \mathcal{R} = B_{\tilde{T}} \cap \tilde{\mathcal{R}}_{\tilde{j}}$$
 on $\{X_{G_{r_1}}(1) = 0, X_{G_{r_0}}(1) > 0\}$ a.s.

Therefore, using the above and (7.8) we obtain

$$\mathbb{P}_{X_{0}}(\dim(B_{r_{0}} \cap \partial \mathcal{R}) \geq d_{f}, X_{G_{r_{1}}}(1) = 0, X_{G_{r_{0}}}(1) > 0)$$

$$\geq \mathbb{P}_{X_{0}}(\dim(B_{\tilde{T}} \cap \partial \mathcal{R}) \geq d_{f}, X_{G_{r_{1}}}(1) = 0, X_{G_{r_{0}}}(1) > 0)$$

$$= \mathbb{P}_{X_{0}}(\dim(B_{\tilde{T}} \cap \partial \tilde{\mathcal{R}}_{\tilde{j}}) \geq d_{f}, X_{G_{r_{1}}}(1) = 0, X_{G_{r_{0}}}(1) > 0)$$

$$\geq \mathbb{P}_{X_{0}}(\{N_{0} \geq 1\} \cap (\bigcap_{i \leq N_{0}} \{\dim(B_{r} \cap \partial \tilde{\mathcal{R}}_{i}) \geq d_{f} \ \forall r > \hat{T}^{j} \geq r_{1}\})).$$

In the last line we have used $N_0 \ge 1$ iff $X_{G_{r_0}}(1) > 0$ (by (7.10)), and on this set, $\hat{T}^j \ge r_1$ for all $j \le N_0$ implies $T_0 \le r_0 - r_1$ (by (7.9)) and so $X_{G_{r_1}}(1) = 0$. We also use the fact (from (7.15)) that if $\{0 < T_0 \le r_0 - r_1\}$ then $\tilde{T} > \hat{T}^{\tilde{j}}$. The independence of the \tilde{W}_i 's and their joint independence

then $\tilde{T} > \hat{T}^{\tilde{j}}$. The independence of the \tilde{W}_j 's and their joint independence from N_0 together with their common law in (7.11) imply that (7.16) equals

$$(7.17) \ \mathbb{E}_{X_0} \Big(1(N_0 \ge 1) \prod_{j=1}^{N_0} \mathbb{N}_{X_0} (\dim(\partial \mathcal{R} \cap B_r) \ge d_f \ \forall r > \hat{T}_0 \ge r_1 | \hat{T}_0 < r_0) \Big).$$

By Proposition 1.7 and Lemma 7.1 each of the terms in the above product equals $\mathbb{N}_{X_0}(\hat{T}_0 \geq r_1 | \hat{T}_0 < r_0)$ and so (7.17) equals

$$\mathbb{E}_{X_0}(1(N_0 \ge 1)1(\wedge_{j=1}^{N_0} \hat{T}_0(\tilde{W}_j) \ge r_1))$$

$$(7.18) \qquad = \mathbb{P}_{X_0}(0 < T_0 \le r_0 - r_1) = \mathbb{P}_{X_0}(X_{G_{r_1}}(1) = 0, X_{G_{r_0}}(1) > 0).$$

In the first equality we used (7.9) and (7.10). We have proved the left-hand side of (7.16) exceeds the above, and we conclude that

$$X_{G_{r_0}}(1) = 0$$
 and $X_{G_{r_0}}(1) > 0$ imply $\dim(B_{r_0} \cap \partial \mathcal{R}) \ge d_f \mathbb{P}_{X_0} - \text{a.s.}$,

thus proving (7.5).

Recall again that we always work with the cadlag version of $Y_r(1)$ from Proposition 6.2(b) which only has non-negative jumps and is an (\mathcal{E}_r^+) supermartingale. Define a sequence of (\mathcal{E}_r^+) -stopping times by

$$T_{n^{-1}} = \inf\{r \le r_0 : Y_r(1) \le 1/n\} \quad (\inf \emptyset = r_0).$$

Then

(7.19) on
$$\{0 < T_0\}$$
 (and so Q_{X_0} -a.s.) $T_{n-1} \uparrow T_0$ and $T_{n-1} < T_0$,

where the last inequality holds since $Y_r(1)$ has no negative jumps. So under Q_{X_0} , T_0 is a predictable stopping time which is announced by $\{T_{n^{-1}}\}$ and so (see (12.9)(ii) in Chapter VI of [22])

$$\mathcal{E}_{T_0-}^+ = \vee_n \mathcal{E}_{T_n}^+.$$

Let $D_r = \{\dim(B_r \cap \partial \mathcal{R}) \geq d_f\}$ for $0 < r \leq r_0$. We assume \mathcal{E}_r^+ is augmented by Q_{X_0} -null sets throughout this Section.

To finish the proof of Proposition 1.7 we need:

LEMMA 7.3. If $X_0 = \delta_{x_0}$ where $|x_0| \ge 2r_0$, then

$$(7.21) D_{r_0} \in \mathcal{E}_{T_0-}^+.$$

For the proof of Proposition 1.7 below it would suffice to show that $D_{r_0} \cap \{T_0 < r_0\} \in \mathcal{E}_{T_0-}^+$, and this latter result should be intuitively obvious, as we now explain. With Lemma 7.1 in mind, we see that $\mathcal{E}_{T_0-}^+$ includes information generated by the excursions of W outside of its minimum radius. If this minimum radius is positive (as is the case on $\{T_0 < r_0\}$) it is intuitively clear that this includes all the information generated by W. Even without intersecting with $\{T_0 < r_0\}$, however, none of the mass that hits the origin will survive for any length of time and so again all of W will have been observed. This last point stems from the fact that points are polar for Brownian motion in more than one dimension and be more formally justified using a mean measure result for the integral of the snake (Proposition 2 in Ch. IV of [17] with p=1).

Proof of Lemma 7.3. Following the derivation of (7.11) in [20], one sees that in order to get (7.21), it is enough to show (cf. (7.18) of [20])

(7.22)
$$\mathbb{N}_{x_0} \left(\int_0^\infty 1(\inf_{v \le \zeta_u} |W_u(v)| = \hat{T}_0) du \right) = 0.$$

Next follow the proof of (7.18) in [20] using the historical process and its Palm measure formula, to bound the left-hand side of (7.22) by (cf. (7.22) of [20])

$$\int_0^\infty E_{x_0}^B \left(\exp\left(-\int_0^s \frac{2(4-d)}{|B_t - m_s|^2} dt \right) \right) ds.$$

Here B denotes a d-dimensional Brownian motion starting at x_0 under $P_{x_0}^B$ and $m_s = \inf_{s' \leq s} |B_{s'}|$. A simple application of Lévy's modulus for B shows that $\int_0^s \frac{2(4-d)}{|B_t-m_s|^2} dt$ is infinite a.s. and so proves (7.22) as required. More details may be found in Supplementary Material [13] where the actual definition of \mathcal{E}_r is even used.

Proof of Proposition 1.7. Clearly it suffices to fix $x_0 \in \operatorname{Supp}(X_0)$ and prove the result with \mathbb{N}_{x_0} in place of \mathbb{N}_{X_0} . By translation invariance we may assume $x_1 = 0$, and so $|x_0| \geq 2r_0$. Fix $0 < r_1 < r_0$. Assume $0 \leq r < r_0$ and $n \in \mathbb{N}$ is large enough so that $r + n^{-1} < r_0$. By Lemma 2.1(b) there is a universally measurable map $\psi : \mathcal{K} \to [0, 1]$ such that

$$1(D_{r_0-r-n^{-1}}) = 1(\dim(\partial(\overline{B}_{r_0-r-n^{-1}} \cap \mathcal{R}) \cap B_{r_0-r-n^{-1}}) \ge d_f) \quad \text{(by (7.8))}$$

$$(7.23) \qquad = \psi(\overline{B}_{r_0-r-n^{-1}} \cap \mathcal{R}).$$

Recall that conditional expectations with respect to \mathcal{E}_r , under \mathbb{N}_{x_0} and Q_{x_0} , agree Q_{x_0} -a.s., and note that Proposition 2.2(b) can be trivially extended to universally measurable maps. Therefore up to Q_{x_0} -null sets, on $\{4n^{-2} \leq Y_r(1) \leq (r_0 - r)^2\} (\in \mathcal{E}_r)$ we have

$$Q_{x_0}(D_{r_0}|\mathcal{E}_r) \ge Q_{x_0}(D_{r_0-r-n^{-1}}|\mathcal{E}_r)$$

$$= \mathbb{P}_{Y_r}(D_{r_0-r-n^{-1}}) \quad \text{(by (7.23) and Proposition 2.2(b))}$$

$$\ge \mathbb{P}_{Y_r}\left(\dim\left(\partial\mathcal{R}\cap B_{r_0-r-(\sqrt{Y_r(1)}/2)}\right) \ge d_f\right)$$

$$\ge q_{5.4},$$

where Lemma 5.4 and the assumed bounds on $Y_r(1)$ are used in the last inequality, and the assumed lower bound on $Y_r(1)$ is used in the next to last

inequality. Let $n \to \infty$ and take limits from above in $r \in \mathbb{Q}_+$ (recall $Y_r(1)$ is cadlag) to conclude that

(7.24)

$$M_r := Q_{x_0}(D_{r_0}|\mathcal{E}_r^+) \ge q_{5.4} \text{ on } \{0 < Y_r(1) < (r_0 - r)^2\} \ \forall r \in \mathbb{Q} \cap (0, r_0) \ Q_{x_0} \text{-a.s.}.$$

Here M_r is a cadlag version of the bounded martingale on the left-hand side. Using right-continuity one can strengthen (7.24) to (7.25)

$$M_r = Q_{x_0}(D_{r_0}|\mathcal{E}_r^+) \ge q_{5.4} \text{ on } \{0 < Y_r(1) < (r_0 - r)^2\} \ \forall r \in (0, r_0) \ Q_{x_0} - \text{a.s.}.$$

On $\{0 < T_0 \le r_0 - r_1\}$ we have from (7.19) and the lack of negative jumps for $Y_r(1)$, (7.26)

for n large,
$$T_{n-1} \in (0, r_0 - r_1)$$
 and $Y_{T_{n-1}}(1) = n^{-1} < (r_0 - T_{1/n})^2 Q_{x_0} - \text{a.s.}$.

By Corollary (17.10) in Chapter VI of [22], (7.25), and (7.26), we have Q_{x_0} -a.s. on $\{0 < T_0 \le r_0 - r_1\} \in \mathcal{E}_{T_0-}^+$,

(7.27)
$$Q_{x_0}(D_{r_0}|\mathcal{E}_{T_0-}^+) = \lim_{n \to \infty} M(T_{n^{-1}}) \ge q_{5.4}.$$

Multiplying the above by $1(\{0 < T_0 \le r_0 - r_1\})$, we see from Lemma 7.3 that

$$1(D_{r_0} \cap \{0 < T_0 \le r_0 - r_1\}) \ge q_{5.4}1(\{0 < T_0 \le r_0 - r_1\})$$
 Q_{x_0} - a.s., and therefore by Lemma 7.1,

$$r_1 \le \hat{T}_0 < r_0 \text{ implies } \dim(B_{r_0} \cap \partial \mathcal{R}) \ge d_f \quad Q_{x_0} - \text{a.s..}$$

This remains true if we replace r_0 by any $r \in (r_1, r_0]$ since we still have $B_{2r} \subset \operatorname{Supp}(X_0)^c$. Therefore we may fix ω outside a Q_{x_0} -null set so that for any $r \in (r_1, r_0] \cap \mathbb{Q}$, $r_1 \leq \hat{T}_0 < r$ implies $\dim(B_r \cap \partial \mathcal{R}) \geq d_f$. By monotonicity of the conclusion in r this means that $\{r_1 \leq \hat{T}_0 < r_0\}$ implies $\dim(B_r \cap \partial \mathcal{R}) \geq d_f$ for all $r > \hat{T}_0$. This gives Proposition 1.7 under Q_{x_0} . The result under \mathbb{N}_{x_0} is now immediate from the definition of Q_{x_0} , and $\{Y_0(1) > 0\} = \{\hat{T}_0 < r_0\} \mathbb{N}_{x_0}$ -a.e. (by Lemma 7.1).

APPENDIX A: APPENDIX

A.1. Proof of Lemma 2.1. (a) Let $K^{(\varepsilon)} = \{x : \rho(x,K) < \varepsilon\}$. If $K_0 \in \mathcal{K}$ and $0 < r \le 1$ are fixed it suffices to show that $\{K \in \mathcal{K} : \rho(K \cap \overline{B}_R, K_0) < r\}$ is Borel. If $r_n \uparrow r$, this set is equal to

$$\{K: K \cap \overline{B}_R \subset K_0^{(r)}\} \cap \{K: K_0 \subset (K \cap \overline{B}_R)^{(r)}\}$$

=\{K: K \cap \overline{B}_R \cap K_0^{(r)}\} \cap \left(\bigcup_{n=1}^\infty\{K: K_0 \subseteq (K \cap \overline{B}_R)^{r_n}\}\right) := S_1 \cap \left(\bigcup_{n=1}^\infty S_2^n\right).

It is then not hard to show that S_1 is open in \mathcal{K} and S_2^n is closed in \mathcal{K} . (b) This easily reduces to showing that for any fixed rationals $q \in (0, \alpha)$ and $r \in (0, R)$, the following describes an universally measurable subset of K's in \mathcal{K} :

For any natural number N there is a finite number of open balls B^1, \ldots, B^M centered at points in \mathbb{Q}^d and with rational radii

$$r_1, \ldots, r_M > 0$$
 satisfying $\sum_{i=1}^M r_i^q < N^{-1}$ so that $\partial K \cap \overline{B}_r \subset \bigcup_{i=1}^M B^i$.

So fixing B^i and r as above, it suffices to show

$$A_1 = \{K \in \mathcal{K} : \partial K \cap \overline{B}_r \subset \bigcup_{i=1}^M B^i\}^c$$
 is an analytic set in \mathcal{K} ,

because this implies A_1 , and hence A_1^c , is a universally measurable set in \mathcal{K} . Let $K_0 = \left(\bigcup_{i=1}^M B^i\right)^c \cap \overline{B}_r \in \mathcal{K}$, $H_0 = \{(x,K) \in \mathbb{R}^d \times \mathcal{K} : x \in K\}$, and for $n \geq 1$, set $H_n = \{(x,K) \in \mathbb{R}^d \times \mathcal{K} : B_{n-1}(x) \not\subset K\}$. Then

$$A_{1} = \{K \in \mathcal{K} : K_{0} \cap \partial K \neq \emptyset\}$$

$$= \{K \in \mathcal{K} : \exists x \in K_{0} \text{ s.t. } x \in K \text{ and } \forall n \in \mathbb{N} \ B_{n^{-1}}(x) \cap K^{c} \neq \emptyset\}$$

$$= \{K \in \mathcal{K} : \exists x \in \mathbb{R}^{d} \text{ s.t. } (x, K) \in (K_{0} \times \mathcal{K}) \cap \left(\bigcap_{n=0}^{\infty} H_{n}\right)\}.$$

Using the well-known fact that the projection of a Borel subset of $K_0 \times \mathcal{K}$ onto \mathcal{K} is an analytic subset of \mathcal{K} (see eg. Theorem 13 in Ch. III of [5] and note the argument goes through with \mathbb{R}^d in place of \mathbb{R}), it then suffices to show each H_n is Borel. H_0 is the countable intersection of the open sets $H_0^M = \{(x, K) : x \in K^{(1/M)}\}$. Moreover it is not hard to see that H_n is open for $n \geq 1$, and we are done.

A.2. Proof of Lemma 4.5. For the proof we will use the following lemma of Marc Yor (see Proposition 2.5 of [20] and [26]). Recall that for $\gamma \in \mathbb{R}$, (ρ_t) denotes a γ -dimensional Bessel process starting from r > 0 under $P_r^{(\gamma)}$, (\mathcal{F}_t) is the filtration generated by ρ and $\tau_R = \inf\{t \geq 0 : \rho_t \leq R\}$.

LEMMA A.1. Let $\lambda \geq 0$, $\mu \in \mathbb{R}$, r > 0 and $\nu = \sqrt{\lambda^2 + \mu^2}$. If $\Phi_t \geq 0$ is \mathcal{F}_{t} -adapted, then for all R < r, we have

$$E_r^{(2+2\mu)}\Big(\Phi_{t\wedge\tau_R}\exp\Big(-\frac{\lambda^2}{2}\int_0^{t\wedge\tau_R}\frac{1}{\rho_s^2}ds\Big)\Big) = r^{\nu-\mu}E_r^{(2+2\nu)}\Big((\rho_{t\wedge\tau_R})^{-\nu+\mu}\Phi_{t\wedge\tau_R}\Big).$$

Now we are ready to give the

Proof of Lemma 4.5. We use Fatou's lemma and then Lemma A.1 to get that for $a \ge 0$,

$$\begin{split} &E_{x}\Big(1_{(\tau_{R}<\infty)} \exp\Big(\int_{0}^{\tau_{R}} \frac{a}{|B_{s}|^{q}} ds\Big) \exp\Big(-\int_{0}^{\tau_{R}} \frac{2(4-d)-\zeta/2}{|B_{s}|^{2}} ds\Big)\Big) \\ &\leq \liminf_{t\to\infty} E_{|x|}^{(2+2\mu)} \Big(1_{(\tau_{R}\leq\tau_{R}\wedge t)} \exp\Big(\int_{0}^{\tau_{R}\wedge t} \frac{a}{\rho_{s}^{q}} ds\Big) \exp\Big(-\int_{0}^{\tau_{R}\wedge t} \frac{2(4-d)-\zeta/2}{\rho_{s}^{2}} ds\Big)\Big) \\ &= |x|^{\nu_{\zeta}-\mu} \liminf_{t\to\infty} E_{|x|}^{(2+2\nu_{\zeta})} \Big(1_{(\tau_{R}\leq\tau_{R}\wedge t)} \exp\Big(\int_{0}^{\tau_{R}\wedge t} \frac{a}{\rho_{s}^{q}} ds\Big) \rho_{t\wedge\tau_{R}}^{\mu-\nu_{\zeta}}\Big) \\ &= (R/|x|)^{\mu-\nu_{\zeta}} E_{|x|}^{(2+2\nu_{\zeta})} \Big(1_{(\tau_{R}<\infty)} \exp\Big(\int_{0}^{\tau_{R}} \frac{a}{\rho_{s}^{q}} ds\Big)\Big) \text{ (since } \rho_{t\wedge\tau_{R}} = R \text{ on } \{\tau_{R}\leq t\}) \\ &= (R/|x|)^{p_{\zeta}} E_{|x|}^{(2+2\nu_{\zeta})} \Big(\exp\Big(\int_{0}^{\tau_{R}} \frac{a}{\rho_{s}^{q}} ds\Big) \Big| \tau_{R}<\infty\Big), \end{split}$$

where in next to the last line we use monotone convergence for $a \geq 0$, and in the last line the hitting probabilities for Bessel processes (e.g. (48.3) and (48.5) in Ch. V of [22]) as well as $p_{\zeta} = \mu + \nu_{\zeta}$. Note that for a < 0, by bounded convergence, we get equality in the second line above (with $\lim \inf_{t \to \infty}$ replaced by $\lim_{t \to \infty}$) and thus proceeding as above we get, by using bounded convergence again in the next to the last line, that (4.10) holds for a < 0.

It remains to verify the lower bound in (4.10), for $a \ge 0$. Fix T > 0. Then we have

$$E_{x}\left(1(\tau_{R} < \infty) \exp\left(\int_{0}^{\tau_{R}} \frac{a}{|B_{s}|^{q}} ds\right) \exp\left(-\int_{0}^{\tau_{R}} \frac{2(4-d)-\zeta/2}{|B_{s}|^{2}} ds\right)\right)$$

$$\geq E_{|x|}^{(2+2\mu)} \left(1_{(\tau_{R} < \infty)} \exp\left(\int_{0}^{\tau_{R} \wedge T} \frac{a}{\rho_{s}^{q}} ds\right) \exp\left(-\int_{0}^{\tau_{R}} \frac{2(4-d)-(\zeta/2)}{\rho_{s}^{2}} ds\right)\right)$$

$$= |x|^{\nu_{\zeta}-\mu} \lim_{t \to \infty} E_{|x|}^{(2+2\nu_{\zeta})} \left(1_{(\tau_{R} \le \tau_{R} \wedge t)} \exp\left(\int_{0}^{\tau_{R} \wedge t \wedge T} \frac{a}{\rho_{s}^{q}} ds\right) R^{\mu-\nu_{\zeta}}\right)$$

$$= (R/|x|)^{\mu-\nu_{\zeta}} E_{|x|}^{(2+2\nu_{\zeta})} \left(1_{(\tau_{R} < \infty)} \exp\left(\int_{0}^{\tau_{R} \wedge T} \frac{a}{\rho_{s}^{q}} ds\right)\right)$$

$$= (R/|x|)^{p_{\zeta}} E_{|x|}^{(2+2\nu_{\zeta})} \left(\exp\left(\int_{0}^{\tau_{R} \wedge T} \frac{a}{\rho_{s}^{q}} ds\right) |\tau_{R} < \infty\right),$$

where in the first equality we used bounded convergence and Lemma A.1, in the second equality bounded convergence again, and in the last equality the hitting probabilities for Bessel processes. Now let $T \to \infty$, to get the required lower bound, and we are done.

SUPPLEMENTARY MATERIAL

Supplementary material to the paper "On the topological boundary of the range of super-Brownian motion'

(doi: COMPLETED BY THE TYPESETTER; .pdf). The supplementary material contains the proof of Proposition 5.1 and gives more details of the proof of Lemma 7.3.

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