

PATTERN FORMATION

(P1)

$$(1) \begin{cases} u_t = D_1 \Delta u + \lambda f(u, v) \\ v_t = D_2 \Delta v + \lambda g(u, v) \end{cases} \quad (\lambda > 0) \quad \lambda \text{ corresponds to changing domain length } \lambda \uparrow \text{ domain is longer.}$$
$$u_x = v_x = 0 \text{ at } x = 0, 1$$

NOW ASSUME THAT (u_e, v_e) IS A STEADY STATE SOLUTION.

DEFINE $J_0 = \begin{pmatrix} f_u^e & f_v^e \\ g_u^e & g_v^e \end{pmatrix}$

THEN IF $\det J_0 > 0$, $\text{trace } J_0 < 0$ WE HAVE THAT (u_e, v_e) IS STABLE WRT ODE DYNAMICS, OR WRT SPATIALLY HOMOGENEOUS PERTURBATIONS.

NOW FOR (1) WE LINEARIZE

$$\begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} u_e \\ v_e \end{pmatrix} + \begin{pmatrix} \zeta \\ \eta \end{pmatrix} e^{\sigma t} e^{i q x}$$

WE OBTAIN THAT

$$(1.5) \quad \sigma \begin{pmatrix} \zeta \\ \eta \end{pmatrix} = \begin{pmatrix} -D_1 q^2 + \lambda f_u^e & \lambda f_v^e \\ \lambda g_u^e & -D_2 q^2 + \lambda g_v^e \end{pmatrix} \begin{pmatrix} \zeta \\ \eta \end{pmatrix}$$

THUS WE GET

$$(2) \quad \sigma^2 - [\lambda \text{trace } J_0 - (D_1 + D_2) q^2] \sigma + Q(q) = 0$$

$$\text{WHERE } Q(q) = D_1 D_2 q^4 - q^2 \lambda [D_1 g_v^e + D_2 f_u^e] + \lambda^2 \det J_0.$$

THEREFORE WE CONCLUDE THAT $\text{RE}[\sigma] < 0 \quad \forall q$ IFF

$$(i) \quad \lambda \text{trace } J_0 - (D_1 + D_2) q^2 < 0$$

$$(ii) \quad Q(q) > 0. \quad (\text{NOTICE } \sigma_+ \sigma_- = Q).$$

NOTICE THAT SINCE $\text{trace } J_0 < 0$ THEN (i) IS ALWAYS SATISFIED.

THUS $\text{RE}[\sigma] > 0$ FOR SOME q WHEN $Q(q) < 0$.

NOW IT IS CONVENIENT TO INTRODUCE $\hat{\lambda}$ by $\lambda = g^2 \hat{\lambda}$.

THEN $Q(g) = g^4 F(\hat{\lambda})$

WITH $f(\hat{\lambda}) = \hat{\lambda}^2 \det J_0 - \hat{\lambda} [D_1 g_v^e + D_2 f_u^e] + D_1 D_2$.

NOW WE FACTOR

$$f(\hat{\lambda}) = \det J_0 \left[(\hat{\lambda} - \hat{\lambda}^*)^2 + \frac{D_1 D_2}{\det J_0} - \left(\frac{D_1 g_v^e + D_2 f_u^e}{2 \det J_0} \right)^2 \right]$$

WHERE $\hat{\lambda}^* = \frac{D_1 g_v^e + D_2 f_u^e}{2 \det J_0}$.

NOW WE HAVE THAT $f(\hat{\lambda}) < 0$ FOR λ ON RANGE

$$\hat{\lambda}_- = \sqrt{\left(\frac{D_1 g_v^e + D_2 f_u^e}{2 \det J_0} \right)^2 - \frac{D_1 D_2}{\det J_0}} < \hat{\lambda} - \hat{\lambda}^* < \sqrt{\left(\frac{D_1 g_v^e + D_2 f_u^e}{2 \det J_0} \right)^2 - \frac{D_1 D_2}{\det J_0}} = \hat{\lambda}_+$$

WHEN $D_1 g_v^e + D_2 f_u^e > 0$

AND $(D_1 g_v^e + D_2 f_u^e)^2 > 4 D_1 D_2 \det J_0$.

THUS WE HAVE $f(\hat{\lambda}) < 0$ ON $\hat{\lambda}_- < \hat{\lambda} - \hat{\lambda}^* < \hat{\lambda}_+$

WHEN THE CONDITION

$$(3) \quad (D_1 g_v^e + D_2 f_u^e) > 2 \sqrt{D_1 D_2 \det J_0} \quad \text{HOLDS.}$$

RECALLING THAT $Q(g) = g^4 F(\hat{\lambda})$ WE THEN HAVE $Q < 0$

ON $g^2(\hat{\lambda}_- + \hat{\lambda}_+) < \lambda < g^2(\hat{\lambda}_- + \hat{\lambda}_+)$

WHEN (3) HOLDS, THEN (1) IS SAID TO BE TURING UNSTABLE.

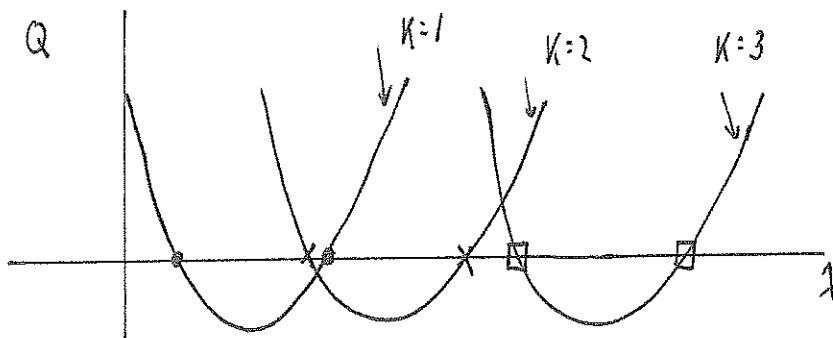
NOW WITH NEUMANN BOUNDARY CONDITION ON A DOMAIN OF LENGTH l WE HAVE THAT

$$Q = Q_K = K\pi \quad \text{WITH} \quad K = 1, 2, 3, \dots$$

WE DO NOT WORRY ABOUT $K=0$ SINCE THIS MODE IS STABLE.

IN THIS WAY WE OBTAIN THE QUALITATIVE PLOT OF Q VERSUS l

FOR $Q = Q_1, Q = Q_2, \dots$ UNDER CONDITION (3)



• points are $Q_1^2 (\hat{\lambda}_x + \hat{\lambda}_x)$
 x points are $Q_2^2 (\hat{\lambda}_x + \hat{\lambda}_x)$
 □ points are $Q_3^2 (\hat{\lambda}_x + \hat{\lambda}_x)$

THIS WE SEE THAT l IS INCREASED THE MODE WITH $K=1$
 I.E. $(0, \pi x)$ GOES UNSTABLE FIRST. AS $l \uparrow$ THE HIGHER MODES
 ARE UNSTABLE. NOW ANALYZE WHEN (3) HOLDS

REMARKS (i) NOW IF $D_1 = D_2$ THEN (3) CANNOT HOLD.

PROOF RECALL WE ASSUMED $\text{trace } J_0 < 0 \rightarrow f_u^e + g_v^e < 0$.

THIS (3) IS IMPOSSIBLE WHEN $D_1 = D_2$.

(ii) IF $D_1 \neq D_2$ THEN FOR $\text{trace } J_0 < 0$, $\det J_0 > 0$ AND (3)

TO HOLD, WE MUST HAVE

(I) ONE OF f_u^e, g_v^e IS POSITIVE.

(II) J_0 CAN'T BE A SYMMETRIC MATRIX.

PROOF IF $D_1 \neq D_2$ THEN IF BOTH f_u^e, g_v^e WERE NEGATIVE, THEN

(3) IS CLEARLY IMPOSSIBLE. THUS AT LEAST ONE OF f_u^e, g_v^e MUST BE POSITIVE. BUT SINCE $\text{trace } J_0 = f_u^e + g_v^e < 0 \rightarrow$ ONLY ONE OF f_u^e, g_v^e IS POSITIVE. TAKING WLOG

$$f_u^e > 0, \quad g_v^e < 0$$

THEN (3) HOLDS ONLY IF D_2/D_1 IS SUFFICIENTLY LARGE.

THE U IS THE ACTIVATOR ($f_u^e > 0 \rightarrow$ AUTOCATALYTIC TERM) WHILE V IS THE INHIBITOR ($g_v^e < 0 \rightarrow$ INHIBITS ITS OWN PRODUCTION).

THE CONDITION D_2/D_1 LARGE ENOUGH MEANS THAT WE MUST HAVE SHORT RANGE ACTIVATION AND LARGE RANGE INHIBITION TO GENERATE A TURING INSTABILITY. THUS PROVES (I).

FOR THE SECOND STATEMENT WE HAVE $\det J_0 = f_u^e g_v^e - f_v^e g_u^e > 0$ IS ASSUMED. BUT SINCE f_u^e, g_v^e HAVE DIFFERENT SIGNS, $f_u^e g_v^e < 0$ WHICH REQUIRES THAT $f_v^e g_u^e < 0$, I.E. f_v^e, g_u^e DIFFERENT SIGNS. THUS

$$J_0 = \begin{pmatrix} f_u^e & f_v^e \\ g_u^e & g_v^e \end{pmatrix} \text{ IS NOT A SYMMETRIC MATRIX.}$$

(iii) FOR $f_u^e > 0, g_v^e < 0$ AND $D_2/D_1 \gg 1$ THEN (3) REDUCES

APPROXIMATELY TO

$$g_v^e + \frac{D_2}{D_1} f_u^e > 2 \sqrt{\frac{D_2}{D_1} \det J_0} \Rightarrow \frac{D_2}{D_1} (f_u^e)^2 > 4 \det J_0.$$

$$\text{OR } D_2/D_1 \gtrsim \frac{4 \det J_0}{(f_u^e)^2}.$$

FOR OTHER SYSTEMS THE BIFURCATION PARAMETER b MAY APPEAR IN THE KINETIC ONLY

$$U_t = D_1 U_{xx} + F(U, V, b)$$

$$V_t = D_2 V_{xx} + g(U, V, b)$$

POSSIBLY ALSO THE SPATIALLY HOMOGENEOUS STATE CAN HAVE A HOPF BIFURCATION AFTER RELAXING CONDITIONS ON F AND g .

EXAMPLE CONSIDER THE Gierer-Meinhardt MODEL

$$U_t = \varepsilon^2 U_{xx} - U + U^p / \sqrt{\gamma}$$

$$p > 1, \gamma > 0, m > 0, s \geq 0,$$

$$\tau V_t = D V_{xx} - V + U^m / \sqrt{s}$$

$$D > 0, \tau > 0.$$

EITHER ON INFINITE LINE OR WITH NEUMANN BC.

THE STEADY-STATE IS $(U_e, V_e) = (1, 1)$.

WE HAVE $F(U, V) = -U + U^p / \sqrt{\gamma}$, $g(U, V) = \frac{1}{\tau} (-V + U^m / \sqrt{s})$

WE CALCULATE (1) $\begin{cases} F_U^e = p-1 & F_V^e = -\gamma \\ g_U^e = m/\tau & g_V^e = -(1+s)/\tau \end{cases}$

THUS PUT $\begin{pmatrix} U \\ V \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} + e^{i\theta x + \sigma t} \begin{pmatrix} \zeta \\ \eta \end{pmatrix}$

WE GET UPON LINEARIZING

$$\begin{pmatrix} -\varepsilon^2 \theta^2 + (p-1) & -\gamma \\ m/\tau & -\frac{D\theta^2 + (1+s)}{\tau} \end{pmatrix} \begin{pmatrix} \zeta \\ \eta \end{pmatrix} = \sigma \begin{pmatrix} \zeta \\ \eta \end{pmatrix}$$

$$\text{set } \det \left(\begin{pmatrix} \dots \end{pmatrix} - \sigma I \right) = 0$$

NOW WE MULTIPLY TO FIND CHARACTERISTIC POLYNOMIAL

$$\begin{pmatrix} -\varepsilon^2 q^2 + (p-1) - \sigma & -\tau \\ m/\tau & -\frac{[Dq^2 + (1+s)]}{\tau} - \sigma \end{pmatrix} \begin{pmatrix} \zeta \\ \eta \end{pmatrix} = 0$$

$$\sigma^2 - \frac{[Dq^2 + (1+s)]}{\tau} [-\varepsilon^2 q^2 + (p-1)] + \sigma [\varepsilon^2 q^2 - (p-1)] + \sigma \frac{[Dq^2 + (1+s)]}{\tau} + \frac{\tau m}{\tau} = 0.$$

DEFINE $J_0 = \begin{pmatrix} p-1 & -\tau \\ m/\tau & -(1+s)/\tau \end{pmatrix}$

$$\text{trace } J_0 = p-1 - (1+s)/\tau$$

$$\det J_0 = -(p-1) \frac{(1+s)}{\tau} + \frac{\tau m}{\tau}$$

WE OBTAIN

$$(2) \quad \left\{ \begin{array}{l} \sigma^2 - \left[\text{trace } J_0 - \frac{(D + \varepsilon^2)}{\tau} q^2 \right] \sigma + Q(q) = 0 \\ \text{WITH } Q(q) = \varepsilon^2 \frac{D}{\tau} q^4 - q^2 \left[-\varepsilon^2 (1+s)/\tau + D/\tau (p-1) \right] + \det J_0 \end{array} \right.$$

WE ASSUME THAT $\det J_0 > 0$, $\text{trace } J_0 < 0$,

$$\text{THUS } \det J_0 > 0 \rightarrow \frac{\tau m}{p-1} - (1+s) > 0$$

$$\text{trace } J_0 < 0 \rightarrow 0 < \tau < \frac{(1+s)}{p-1}$$

NOTICE THAT IF WE LABEL

$$J = \frac{\tau m}{p-1} - (1+s)$$

THEN $\det J_0 = \frac{(p-1)}{\tau} J.$

(3)

WE CONCLUDE THAT $\text{Re } \sigma < 0$ IFF

(P7)

$$\left. \begin{aligned} \text{trace } J_0 - \left(\frac{D}{\tau} + \varepsilon^2 \right) q^2 < 0 \\ \text{AND } Q(q) > 0. \end{aligned} \right\} \quad (4)$$

SINCE THE FIRST CONDITION ALWAYS HOLDS, WE HAVE INSTABILITY
FOR SOME WAVELENGTH q IFF

$$Q(q) < 0 \quad \text{FOR SOME } q_- < q < q_+.$$

NOW WE CONSIDER

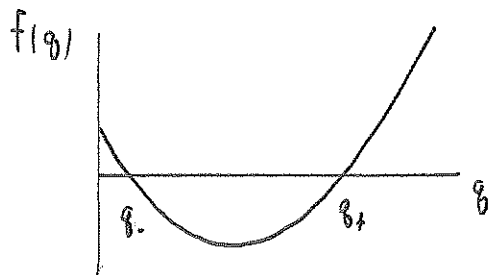
$$Q(q) = \frac{1}{\tau} F(q) \quad (5)$$

$$\text{WITH } F(q) = \varepsilon^2 D q^4 - q^2 [D(p-1) - \varepsilon^2(1+p)] + (p-1)$$

WE NOW LOOK FOR ROOTS OF $F(q) = 0$.

NOTE: $F(0) = (p-1) > 0$, $F \rightarrow +\infty$ AS $q \rightarrow \infty$. HOWEVER,

IF $\varepsilon < 0$, WE HAVE $F(q) < 0$ FOR $0(1) < q < 0(1/\varepsilon)$.



TO FIND THE EDGES OF THE INSTABILITY BAND FIND APPROXIMATE
ROOTS OF (5) WHEN $\varepsilon \ll 1$.

UPPER EDGE BALANCE 1ST AND 2ND TERM IN F IN (5). $\varepsilon^2 q^4 \sim q^2$

$$\text{THUS PUT } q = \frac{1}{\varepsilon} q_{+0} \text{ SO } F \sim \frac{1}{\varepsilon^2} [D q_{+0}^4 - q_{+0}^2 [D(p-1) \dots]]$$

$$\text{SO } q_{+0} \approx \sqrt{p-1}$$

NOW TO FIND LOWER EDGE BALANCE 2ND AND 3RD TERM IN

(P8)

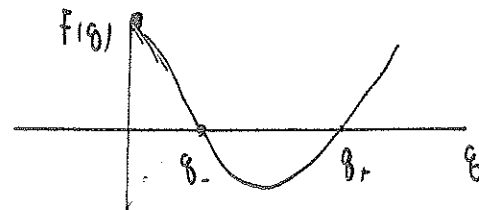
TO OBTAIN

$$-q_-^2 [D(p-1) - \dots] + (p-1)J + \text{small} \approx 0$$

THU $q_- = \sqrt{\frac{J}{D}}$

WE CONCLUDE THAT IN THE SINGULARLY PERTURBED LIMIT $\epsilon \rightarrow 0$, THE GM MODEL IS TURNING UNSTABLE TO A WIDE BAND OF MODES WITH

$$q_- < q < q_+$$



WITH $q_- \sim \sqrt{\frac{J}{D}}$ AS $\epsilon \rightarrow 0$ $J \equiv \frac{m\Gamma}{p-1} - (J+1)$

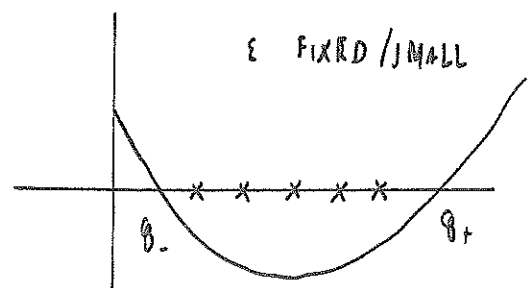
$q_+ \sim \sqrt{\frac{p-1}{\epsilon}}$ AS $\epsilon \rightarrow 0$

ON A FINITE DOMAIN OF LENGTH 1 WITH NEUMANN BC

$u_x = v_x = 0$ AT $x=0, 1$, WE HAVE

$$q_k = k\pi, \quad k=1, 2, 3, \dots$$

THU FOR $\epsilon \rightarrow 0$, $\exists$ MANY k FOR WHICH q_k LIES IN THE INTERVAL $q_- < q < q_+$

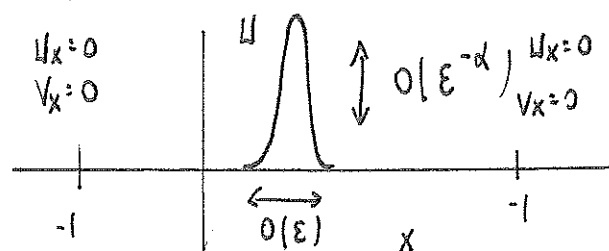


THE QUESTION IS WHAT IS THE "ATTRACTOR" THAT RESULTS FROM THE WIDE BAND OF UNSTABLE MODELS. IT TURNS OUT TO BE LOCALIZED SPIKE SOLUTIONS. FOR INSTANCE CONSIDER

$$U_t = \varepsilon^2 U_{xx} - U + U^3/V^2$$

$$\tau V_t = D V_{xx} - V + U^3$$

CAN WE CONSTRUCT A SOLUTION WHERE U IS LOCALIZED IN THE FORM OF SPIKE



WE PUT $U = \bar{U} \varepsilon^{-\alpha}$ AND $V = \bar{V} \varepsilon^{-\beta}$. WE WANT THAT

THE U EQUATION IS INVARIANT: THUS $[U^2]/[V^2] = 1$ SO $\alpha = \beta$.

NOW IN V -EQUATION WE INTEGRATE FROM $-1 < x < 1$ TO GET

$$\int_{-1}^1 V dx = \int_{-1}^1 U^3 dx \quad \text{BUT SINCE } U = O(\varepsilon^{-\alpha}) \text{ IN A SPIKE OF EXTENT } O(\varepsilon)$$

IN x , WE MUST HAVE $-\beta = 1 - 3\alpha$ SO IF $\alpha = \beta \rightarrow \alpha = 1/2, \beta = 1/2$.

THUS WE RESCALE $U = \varepsilon^{1/2} \bar{U}, V = \varepsilon^{1/2} \bar{V}$ TO OBTAIN

$$\varepsilon^{-1/2} \bar{U}_t = \varepsilon^{-1/2} (\varepsilon^2 \bar{U}_{xx} - \bar{U} + \bar{U}^3/\bar{V}^2)$$

$$\varepsilon^{-1/2} \tau \bar{V}_t = D \varepsilon^{-1/2} \bar{V}_{xx} - \varepsilon^{-1/2} \bar{V} + \varepsilon^{-3/2} \bar{U}^3$$

THIS YIELDS THE RESCALED SYSTEM

$$(*) \quad \begin{cases} \bar{U}_t = \varepsilon^2 \bar{U}_{xx} - \bar{U} + \bar{U}^3/\bar{V}^2 \\ \tau \bar{V}_t = D \bar{V}_{xx} - \bar{V} + \varepsilon^{-1} \bar{U}^3 \end{cases}$$

IN INNER REGION FOR STEADY-STATE, WE PUT $y = x/\varepsilon$ AND GET $\bar{V} = \bar{V}_0$

THEN

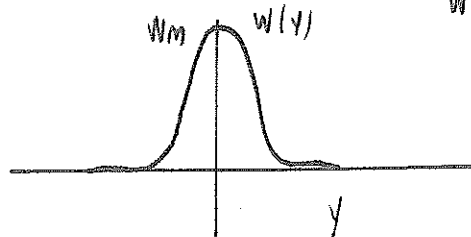
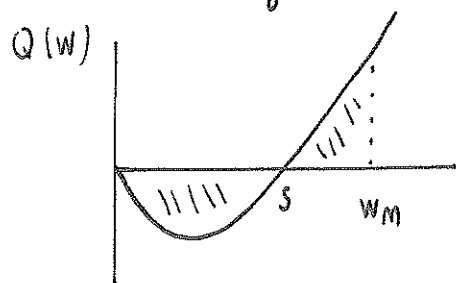
$$\bar{U}_{yy} - \bar{U} + \bar{U}^3/\bar{V}_0^2 = 0$$

A HOMOCLINIC SOLUTION $w(y)$ EXISTS FOR

$$w'' + Q(w) = 0$$

WHEN $Q(0) = 0$, $Q'(0) < 0$, $Q(s) = 0$ FOR SOME $s > 0$, AND $\exists w_m > s$

SUCH THAT $\int_0^{w_m} Q(w) dw = 0$ WITH $Q(w_m) > 0$. UNIQUE IF $w'(0) = 0$.
 $w(0) = w_m$



A SPECIAL CLASS OF SUCH HOMOCLINIC IS WHERE $Q(w) = -w + w^p$
 WITH $p > 1$.

WE MULTIPLY BY w' TO GET $w^{1/2} + V(w) = 0$ ON $0 < y < \infty$.

WITH $V(w) = \int_0^w Q(s) ds$. NOW $w'(0) = 0$ IF $w = w_m$.

WE CAN THEN WRITE $w' = -[-2V(w)]^{1/2}$ WHICH IS SEPARABLE
 AND CAN BE INTEGRATED TO GET $w(y)$ IMPLICITLY.

WE NOW FOCUS ON

$$(1) \quad \left\{ \begin{array}{l} w'' - w + w^p = 0 \quad \text{WITH } p > 1. \\ w(0) > 0, \quad w'(0) = 0, \quad w \rightarrow 0 \text{ AS } |y| \rightarrow \infty. \end{array} \right.$$

THE EXACT HOMOCLINIC IS

$$(2) \quad w(y) = \left(\frac{p+1}{2} \right)^{1/p-1} \left(\operatorname{sech} \left[\frac{(p-1)y}{2} \right] \right)^{2/p-1}$$

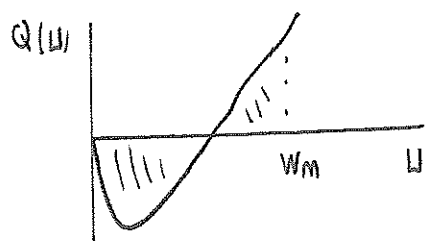
IF $p = 2$ WE HAVE $w = \frac{3}{2} \operatorname{sech}^2(y/2)$

IF $p = 3$ $w = \sqrt{2} \operatorname{sech} y$.

NOW FOR THE EXISTENCE OF SPIKES WE CAN MODIFY PROPERTIES

(H2)

OF $Q(U)$ SLIGHTLY.



$$\text{let } Q(U) = U \log U$$

$$\text{so THAT } Q \rightarrow 0 \text{ AS } U \rightarrow 0^+$$

$$Q'(0) \rightarrow -\infty \text{ AS } U \rightarrow 0^+$$

THE EXACT SOLUTION TO

$$W'' + W \log W = 0, \quad W \rightarrow 0 \text{ as } |y| \rightarrow \infty$$

$$W(0) = W_m, \quad W'(0) = 0$$

$$\text{IS } W = e^{1/2} e^{-y^2/4} \text{ WHICH HAS FASTER THAN}$$

$$\text{SIMPLE EXPONENTIAL DECAY } O(e^{-\sqrt{y}}) \text{ AS } y \rightarrow \infty.$$

REMARK

IN 2-D WE HAVE TO CONSIDER FOR $p > 1$,

$$\Delta_p W = W + W^p = 0$$

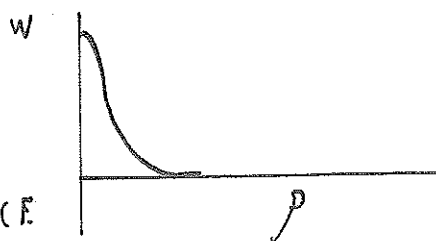
$$\text{WITH } \Delta_p = W'' + \frac{1}{p} W'$$

WE WANT A RADIALY SYMMETRIC SOLUTION $W(\rho)$ WITH

$$W \rightarrow 0 \text{ AS } \rho \rightarrow \infty$$

$$W = O(\rho^{-1/2} e^{-\rho}) \text{ AS } \rho \rightarrow \infty$$

$$W(0) > 0, \quad W'(0) = 0.$$



WE WILL DISCUSS THE RIGOROUS PROPERTIES REGARDING THE EXISTENCE

OF SUCH SOLUTIONS IN 2-D, AND UNIQUENESS.

RATHER CURIOUSLY, $\exists$ EXACT SOLUTION TO

$$\Delta_p W + W \log W = 0 \quad \text{IN } p > 0$$

$$W'(0) = 0, \quad W(0) > 0, \quad W \rightarrow 0 \text{ AS } \rho = |y| \rightarrow \infty.$$

ONE CAN VERIFY

$$W = e^{1 - \rho^2/4}, \text{ WHICH IS A GAUSSIAN.}$$

NOW WE RETURN TO 1-D AND CONSIDER SPECTRAL PROPERTIES OF THE LINEARIZED OPERATOR $L_0 \Phi \equiv \Phi'' + Q'(W) \Phi$. NOTICE L_0 IS SELF-ADJOINT FOR $Q(W) = -W + W^P$ WE HAVE

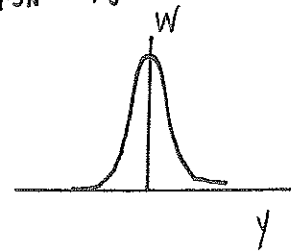
$$(3) \quad \left\{ \begin{array}{l} L_0 \Phi \equiv \Phi'' - \Phi + p W^{p-1} \Phi = \gamma \Phi \quad \text{ON } -\infty < y < \infty \\ \Phi \rightarrow 0 \quad \text{AS } |y| \rightarrow \infty. \end{array} \right.$$

WHERE $W(y)$ GIVEN IN (2) IS UNIQUE SOLUTION TO

$$W'' - W + W^P = 0, \quad -\infty < y < \infty$$

$$W \rightarrow 0 \quad \text{AS } |y| \rightarrow \infty$$

$$W'(0) = 0, \quad W(0) > 0.$$



THE SPECTRUM OF (3) HAS BOTH A DISCRETE AND A CONTINUOUS SPECTRUM.

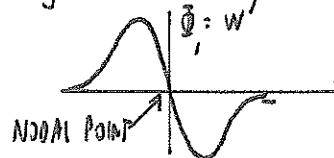
SIMPLE PROPERTIES

(i) CONTINUOUS SPECTRUM IS IN $\gamma < -1$ REAL.

(ii) ANY DISCRETE EIGENVALUE MUST BE REAL SINCE L_0 SELF-ADJOINT.

(iii) $L_0 W' = 0$ SO $\gamma_1 = 0$ IS EIGENVALUE FOR $\Phi_1 = W'$

THIS IS TRANSLATION MODE

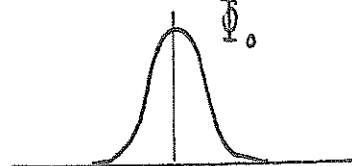


(iv) SINCE $\gamma_1 = 0 \rightarrow$ THE EIGENFUNCTION HAS Φ_1 HAS ONE NODAL POINT.

SINCE THE PRINCIPAL OR LARGEST EIGENVALUE MUST HAVE NO NODAL POINTS, I.E. IS OF ONE SIGN, AND IS A SIMPLE EIGENVALUE,

THEN $\exists \gamma_0 > 0$ AND $\Phi_0 > 0$ FOR WHICH

$$L_0 \Phi_0 = \gamma_0 \Phi_0$$



I.E. $\exists$ ONE UNITABLE EIGENVALUE.

(iv) WE CONCLUDE THAT THE HOMOCLINIC SOLUTION TO

(44)

$$u_t = u_{yy} - u + u^p$$

IS UNSTABLE. i.e. LET W BE HOMOCLINIC. LINEARIZE

BY WRITING $u = w + e^{\nu t} \bar{\phi}(y)$

TO OBTAIN $L_0 \bar{\phi} = \bar{\phi}'' - \bar{\phi} + p w^{p-1} \bar{\phi} = \nu \bar{\phi}$.

BUT $\exists \nu_0 > 0, \rightarrow$ UNSTABLE.

A MORE PRECISE RESULT FOR THE DISCRETE SPECTRUM OF L_0 IS GIVEN IN THE FOLLOWING PROPOSITION

PROPOSITION [DOELMAN ET. AL] LET J BE A POSITIVE INTEGER

SATISFYING $J < \frac{p+1}{p-1} \leq J+1$.

THEN (3) HAS $J+1$ DISCRETE EIGENVALUES ν_j FOR $j=0, \dots, J$ GIVEN BY

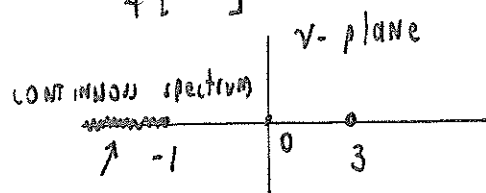
$$\nu_j = \frac{1}{4} \left[(p+1) - j(p-1) \right]^2 - 1.$$

THE PROOF OF THIS IS DIFFICULT AND IS BASED ON OBSERVING THAT A DIFFERENTIAL OPERATOR WITH sech^2 POTENTIAL CAN BE CONVERTED TO A PROBLEM WITH HYPERGEOMETRIC FUNCTION.

EXAMPLE LET $p=3$, THEN $J < 2 \leq J+1$ so $J=1$. THEN $\exists$ TWO DISCRETE EIGENVALUES

$$\nu_0 = \frac{1}{4} [4]^2 - 1 = 3,$$

$$\nu_1 = \frac{1}{4} [4] - 1 = 0.$$



$$w = \sqrt{2} \text{sech } y$$

EXAMPLELET $p = 2$ SO THAT $J < 3 \leq J+1$ SO $J = 2$.

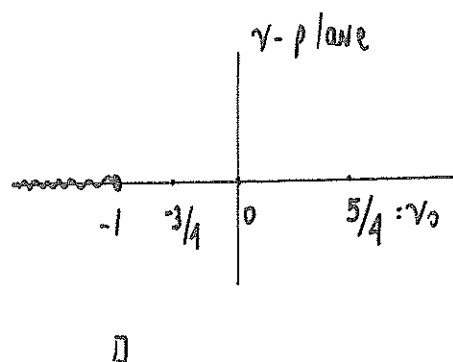
(45)

THUS $J = 3$ DISCRETE EIGENVALUES

$$\gamma_0 = \frac{1}{4} (3)^2 - 1 = 5/4$$

$$\gamma_1 = 0$$

$$\gamma_2 = 1/4 (3 \cdot 2)^2 - 1 = -3/4$$



THE NEXT RESULT GIVES SOME PROPERTIES OF L_0 AND INTERVALS OF W WHICH WE WILL USE LATER.

LEMMA THE FOLLOWING PROPERTIES HOLD IN 1-D FOR L_0 IN (3).

$$(i) \quad L_0 W = (p-1) W^p$$

$$(ii) \quad \text{IF } p = 3 \text{ THEN } L_0 (W^3) = 3 |W^3|.$$

$$(iii) \quad L_0 \left(\frac{W}{p-1} + \frac{1}{2} \gamma W' \right) = W$$

$$(iv) \quad \frac{\int_{-\infty}^{\infty} W^2 dy}{\int_{-\infty}^{\infty} W'^2 dy} = \frac{2(p+1)}{(p-1)} \quad \text{AND} \quad \frac{\int_{-\infty}^{\infty} W^{p+1} dy}{\int_{-\infty}^{\infty} W'^2 dy} = \frac{p+3}{p-1}$$

PROOF

$$(i) \quad L_0 W = (W)'' - (W) + p W^{p-1} (W) = W'' - W + p W^p \quad \text{BUT } W'' - W = -W^p$$

$$\text{SO THAT } L_0 W = -W^p + p W^p = (p-1) W^p.$$

$$(ii) \quad \text{IF } p = 3 \text{ WE CALCULATE}$$

$$L_0 (W^3) = (W^3)'' - (W^3) + 3 W^2 |W^3|$$

$$= (2 W W')' - W^3 + 3 W^4$$

$$= 2 W'^2 + 2 W W'' - W^3 + 3 W^4.$$

BUT $W'' = W - W^3$ AND FROM A FIRST INTEGRAL $W'^2/2 = W^2/2 - W^4/4$. (H6)

WE SUBSTITUTE

$$L_0(W') = 2W \left(\underline{W - W^3} \right)' + 2 \left(\underline{W^2 - W^4/2} \right)' - \underline{W^2} + 3W^4' = 3W^2. \quad \square$$

WE SHOWED EARLIER THAT $\gamma_0 = 3$, THUS $\Phi_0 = W^2 = 2 \operatorname{sech}^2(\gamma)$.

$$\begin{aligned} \text{(iii)} \quad L_0 \left(\frac{W}{p-1} + \frac{1}{2} \gamma W' \right) &= \left(\frac{W}{p-1} + \frac{1}{2} \gamma W' \right)'' - \left(\frac{W}{p-1} + \frac{1}{2} \gamma W' \right) + p W^{p-1} \left(\frac{W}{p-1} + \frac{1}{2} \gamma W' \right) \\ &= \frac{W''}{p-1} + W'' + \frac{1}{2} \gamma W''' - \frac{W}{p-1} - \frac{1}{2} \gamma W' + \frac{p W^p}{p-1} + \frac{1}{2} p \gamma W^{p-1} W' \\ &= \frac{p}{p-1} \left(\underline{W - W^p} \right) + \frac{1}{2} \gamma \left(\underline{W' - p W^{p-1} W'} \right) - \frac{W}{p-1} - \frac{1}{2} \gamma W' + \frac{p W^p}{p-1} + \frac{1}{2} p \gamma \underline{W^{p-1} W'} \end{aligned}$$

ONCE WE USE $W'' = -W + W^p$ AND $W''' = -W' + p W^{p-1} W'$.

CANCELLING THE UNDERLINED TERMS GIVES

$$L_0 \left(\frac{W}{p-1} + \frac{1}{2} \gamma W' \right) = \frac{p}{p-1} W - \frac{W}{p-1} = W. \quad \square$$

(iv) TO CALCULATE THE INTEGRALS

$$W'' - W + W^p = 0.$$

MULTIPLY BY W' AND INTEGRATE

$$W'^2/2 = W^2/2 - W^{p+1}/(p+1)$$

$$\text{THIS GIVES} \quad \int_{-\infty}^{\infty} W'^2 dy = \int_{-\infty}^{\infty} W^2 dy = \frac{2}{p+1} \int_{-\infty}^{\infty} W^{p+1} dy.$$

$$\text{LET} \quad I_1 = \frac{\int_{-\infty}^{\infty} W^2 dy}{\int_{-\infty}^{\infty} W'^2 dy} \quad I_2 = \frac{\int_{-\infty}^{\infty} W^{p+1} dy}{\int_{-\infty}^{\infty} W'^2 dy}.$$

$$\text{THEN} \quad I = I_1 - \frac{2}{p+1} I_2. \quad (X)$$

TO GET A SECOND EQUATION MULTIPLY BY W AND USE INTEGRATION BY PARTS

$$\int_{-\infty}^{\infty} W W'' dy = - \int_{-\infty}^{\infty} W'^2 dy = \int_{-\infty}^{\infty} W^2 dy - \int_{-\infty}^{\infty} W^{p+1} dy$$

THEY YIELD $-1 = I_1 = I_2$ (XX)

SOLVING (X) AND (XX) FOR I_1 AND I_2 WE GET

$$I_2 = \frac{2(p+1)}{(p-1)}$$

$$I_1 = \frac{p+3}{p-1}$$

WE NOW BRIEFLY REMARK ON PROPERTIES OF HETEROCLINIC SOLUTIONS

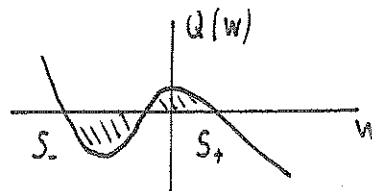
HETEROCLINIC SOLUTIONS

CONSIDER $w'' + Q(w) = 0$

$w \rightarrow S_+$ AS $y \rightarrow +\infty$, $w \rightarrow S_-$ AS $y \rightarrow -\infty$, WITH $S_+ > S_-$

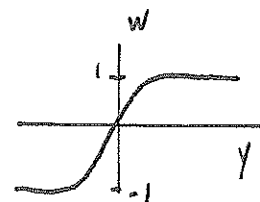
WE PUT $w'(0) > 0$, $w(0) = 0$

WE REQUIRE $Q(S_+) = 0$, $Q'(S_+) < 0$ AND



$\int_{S_-}^{S_+} Q(w) dw = 0$ FOR EXISTENCE OF HETEROCLINIC

NOW IF $Q(w) = 2(w - w^3)$ THEN $S_{\pm} = \pm 1$ AND $w = \tanh y$

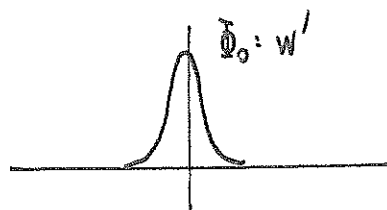


THE LINEARIZED OPERATOR IS

$$L_0 \Phi = \Phi'' + \Phi - 3w^2 \Phi = \gamma \Phi$$

WE CLAIM $\gamma_0 = 0$, $\Phi_0 = w'$ (WHICH HAS NO NODAL POINTS)

IS THE FIRST EIGENPAIR



THUS FOR THE PDE

$$u_t = u_{yy} + 2(u - u^3) \text{ THEN } u_e = w(y)$$

IS NEUTRALLY STABLE.