

Lecture 12: Heat equation on a circular ring - full Fourier Series

(Compiled 3 March 2014)

In this lecture we use separation of variables to solve the heat equation subject on a thin circular ring with *periodic boundary conditions*. In this case we reduce the problem to expanding the initial condition function $f(x)$ in an infinite series of both cosine and sine functions, which we refer to as the Full Range Fourier Series.

Key Concepts: Heat Equation; Periodic Boundary Conditions; separation of variables; Full Fourier Series.

12.1 Heat equation on a circular Ring - Full Fourier Series

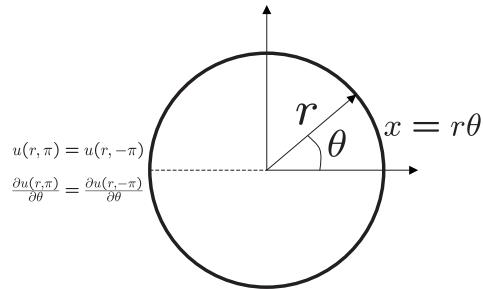


FIGURE 1. Consider a thin conducting ring with thermal conductivity α^2 that has a given initial temperature distribution

Physical Interpretation: Consider a thin circular wire in which there is no radial temperature dependence, i.e., $u(r, \theta) = u(\theta)$ so that $\frac{\partial u}{\partial r} = 0$. In this case the polar Laplacian reduces to

$$\begin{aligned}\Delta u &= \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} \\ &= \frac{\partial^2 u}{\partial(r\theta)^2}\end{aligned}\tag{12.1}$$

and if we let $x = r\theta$ then $\frac{\partial^2 u}{\partial(r\theta)^2} = u_{xx}$. In this case the heat distribution in the ring is determined by the following initial value problem with periodic boundary conditions

$$u_t = \alpha^2 u_{xx} \quad (12.2)$$

$$\text{BC: } \left. \begin{array}{l} u(-L, t) = u(L, t) \\ \frac{\partial u}{\partial x}(-L, t) = \frac{\partial u}{\partial x}(L, t) \end{array} \right\} \text{ Periodic BC}$$

$$\text{IC: } u(x, 0) = f(x)$$

Assume $u(x, t) = X(x)T(t)$.

$$\text{As before: } \frac{X''(x)}{X(x)} = \frac{\dot{T}(t)}{\alpha^2 T(t)} = -\lambda^2.$$

$$\text{IVP: } \frac{\dot{T}(t)}{\alpha^2 T(t)} = -\lambda^2 \Rightarrow T(t) = ce^{-\lambda^2 t}.$$

$$r = \frac{2L}{2\pi} = \frac{L}{\pi} = \text{ Constant.}$$

$$\text{BVP: } \left. \begin{array}{l} X'' + \lambda^2 X = 0 \\ X(-L) = X(L) \\ X'(-L) = X'(L) \end{array} \right\} \begin{array}{l} \text{Eigenvalue Problem} \\ \text{look for } \lambda \text{ such that} \\ \text{nontrivial } x \text{ can be found.} \end{array}$$

$$X(x) = A \cos(\lambda x) + B \sin(\lambda x)$$

$$X(-L) = A \cos(\lambda L) - B \sin(\lambda L) = A \cos(\lambda L) + B \sin(\lambda L) = X(L)$$

$$\text{therefore } 2B \sin(\lambda L) = 0$$

$$X'(x) = -A\lambda \sin(\lambda x) + B\lambda \cos(\lambda x) \quad (12.3)$$

$$X'(-L) = +A\lambda \sin(\lambda L) + B\lambda \cos(\lambda L) = -A\lambda \sin(\lambda L) + B\lambda \cos(\lambda L) = X'(L)$$

$$\text{therefore } 2A\lambda \sin(\lambda L) = 0$$

$$\text{Therefore } \lambda_n L = (n\pi) \quad n = 0, 1, \dots$$

Solutions to (12.2) that satisfy the BC are thus of the form

$$u_n(x, t) = e^{-(\frac{n\pi}{L})^2 \alpha^2 t} \left\{ A_n \cos\left(\frac{n\pi x}{L}\right) + B_n \sin\left(\frac{n\pi x}{L}\right) \right\}. \quad (12.4)$$

Superposition of all these solutions yields the general solution

$$u(x, t) = A_0 + \sum_{n=1}^{\infty} \left\{ A_n \cos\left(\frac{n\pi x}{L}\right) + B_n \sin\left(\frac{n\pi x}{L}\right) \right\} e^{-(\frac{n\pi}{L})^2 \alpha^2 t}. \quad (12.5)$$

In order to match the IC we have

$$f(x) = u(x, 0) = A_0 + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) + B_n \sin\left(\frac{n\pi x}{L}\right). \quad (12.6)$$

As before we obtain expressions for the A_n and B_n by projecting $f(x)$ onto $\sin\left(\frac{n\pi x}{L}\right)$ and $\cos\left(\frac{n\pi x}{L}\right)$.

$$\begin{aligned} \int_{-L}^L f(x) \left\{ \begin{array}{l} \sin\left(\frac{m\pi x}{L}\right) \\ \cos\left(\frac{m\pi x}{L}\right) \end{array} \right\} dx &= A_0 \int_{-L}^L \left\{ \begin{array}{l} \sin\left(\frac{m\pi x}{L}\right) \\ \cos\left(\frac{m\pi x}{L}\right) \end{array} \right\} dx \\ &\quad + \sum_{n=1}^{\infty} A_n \int_{-L}^L \cos\left(\frac{n\pi x}{L}\right) \left\{ \begin{array}{l} \sin\left(\frac{m\pi x}{L}\right) \\ \cos\left(\frac{m\pi x}{L}\right) \end{array} \right\} dx \\ &\quad + \sum_{n=1}^{\infty} B_n \int_{-L}^L \sin\left(\frac{n\pi x}{L}\right) \left\{ \begin{array}{l} \sin\left(\frac{m\pi x}{L}\right) \\ \cos\left(\frac{m\pi x}{L}\right) \end{array} \right\} dx. \end{aligned} \quad (12.7)$$

As in the previous example we use the orthogonality relations:

$$\begin{aligned} \int_{-L}^L \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) dx &= L\delta_{mn} \\ \int_{-L}^L \cos\left(\frac{m\pi x}{L}\right) \cos\left(\frac{n\pi x}{L}\right) dx &= L\delta_{mn} \quad m \text{ and } n \neq 0 \\ &= 2L \quad m = n = 0 \\ \int_{-L}^L \sin\left(\frac{m\pi x}{L}\right) \cos\left(\frac{n\pi x}{L}\right) dx &= 0 \quad \forall m, n. \end{aligned} \quad (12.8)$$

Plugging these orthogonality conditions into (1.6) we obtain

$$\left. \begin{aligned} A_0 &= \frac{1}{2L} \int_{-L}^L f(x) dx = \text{average value of } f(x) \text{ on } [-L, L] \\ A_n &= \frac{1}{L} \int_{-L}^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx \text{ and } B_n = \frac{1}{L} \int_{-L}^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx. \end{aligned} \right\} \quad (12.9)$$

Observations:

- (1) (12.6) and (12.9) represent the full Fourier Series Expansion for $f(x)$ on the interval $[-L, L]$.
- (2) By defining $a_n = \frac{1}{L} \int_{-L}^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx = \begin{cases} 2A_0 & \text{and } b_n = B_n \\ A_n & \end{cases}$ the Fourier Series (12.6) is often written in the form

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi x}{L}\right) + b_n \sin\left(\frac{n\pi x}{L}\right). \quad (12.10)$$