

CONSIDER SYSTEMS OF THE FORM

$$X'' = F(X)$$

MULTIPLY BY X' TO GET

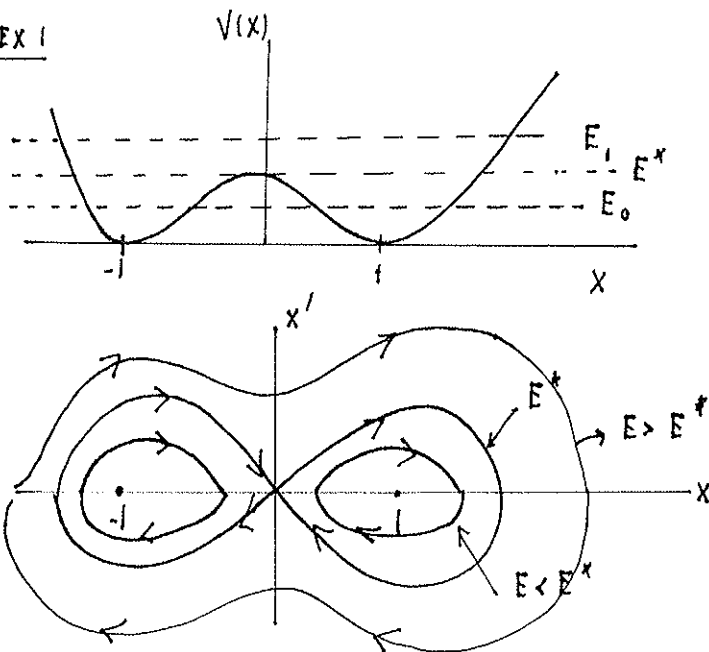
$$X' X'' = X' F(X)$$

LET $V(X) = - \int^X F(\lambda) d\lambda$. THEN, $\frac{dV}{dt} = -F'(x) \frac{dx}{dt}$. $V'(X) = -F(X)$

THUS, $\frac{d}{dt} \left[\frac{1}{2} X'^2 + V(X) \right] = 0$

HENCE $\frac{1}{2} X'^2 + V(X) = E$ E total energy.

EX 1



$$V(X) = \frac{1}{4} (1 - X^2)^2$$

$$V'(X) = -F(X)$$

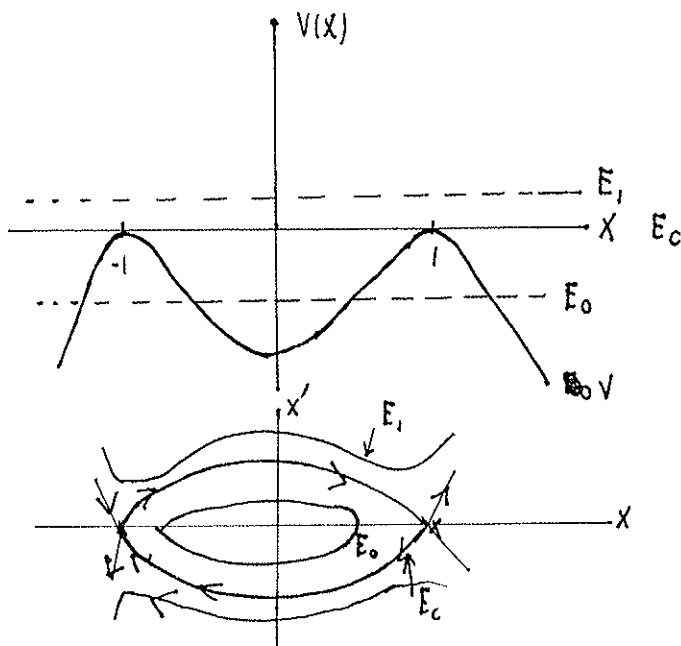
$$\rightarrow F(X) = X(1 - X^2)$$

$$X'' = X(1 - X^2)$$

$$X'' = -X + X^3 = 0$$

EX 2

LET $F(X) = -X + X^3 \rightarrow V(X) = -\frac{1}{4} (1 - X^2)^2$



closed orbits are periodic solutions.

heteroclinic connection $-1 \rightarrow 1$.

NOTE:

(i) symmetry $\dot{X} \rightarrow -\dot{X}$

$$V'(x_E) = 0.$$

WE LET

$$x = x_E + \tilde{x} \quad \text{WITH } \tilde{x} \ll 1 \text{ TO GET}$$

$$\tilde{x}'' = -V'(x_E) + \tilde{x} V''(x_E) + \dots$$

$$\text{THUS} \quad \tilde{x}'' + V''(x_E) \tilde{x} = 0$$

$$\text{NOW} \quad \tilde{x} = e^{\lambda t} \rightarrow \lambda^2 + V''(x_E) = 0$$

(i) IF $V''(x_E) > 0 \rightarrow x_E$ IS LOCAL MINIMA OF V THEN

$$\lambda = \pm i [V''(x_E)]^{1/2}.$$

$$\tilde{x} \sim A \cos[\omega_E t] + B \sin[\omega_E t] \quad \omega_E = \sqrt{V''(x_E)} \text{ frequency}$$

local oscillations x_E is called a center

$\tilde{x}$ is bounded as $t \rightarrow \infty \rightarrow x_E$ is neutrally stable.

(ii) $V''(x_E) < 0 \rightarrow \lambda = \pm \sqrt{-V''(x_E)}$

$$\tilde{x} \sim A e^{\sqrt{-V''(x_E)} t} + B e^{-\sqrt{-V''(x_E)} t}$$

exponential growth and decay. x_E is called a saddle point.

EXAMPLE 1 (SIMPLE PENDULUM)

$$\ddot{\varphi} + \frac{g}{L} \sin \varphi = 0$$

$$V'(\varphi) = \frac{g}{L} \sin \varphi$$



$$V'(\varphi) = \frac{g}{L} \sin \varphi$$

$$V(\varphi) = \frac{g}{L} (1 - \cos \varphi)$$

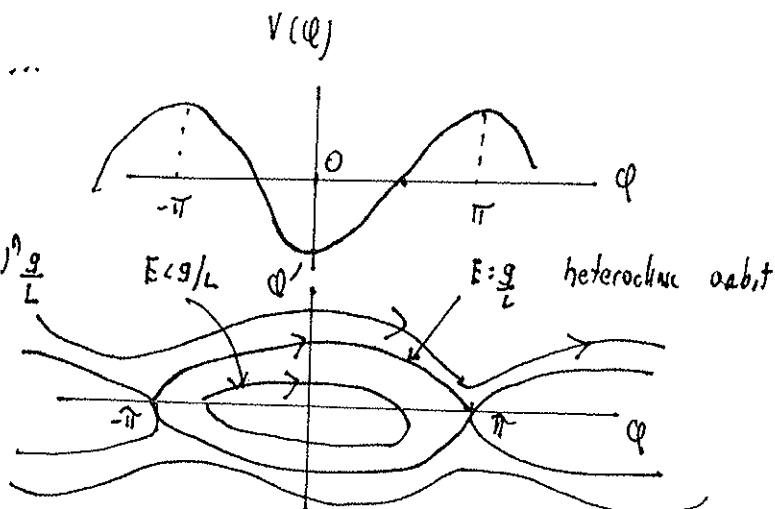
$$V(\varphi) = -\frac{g}{L} \cos \varphi$$

$$\text{THUS} \quad \frac{\varphi'^2}{2} + \frac{g}{L} (-\cos \varphi) = E$$

NOW $V'(\varphi) = 0$ AT $\varphi = 0, \pm\pi, \pm 2\pi, \dots$

$$\text{LET} \quad \frac{\varphi'^2}{2} + \frac{g}{L} (-\cos \varphi) = E$$

$$\text{NOW} \quad V''(\varphi) = \frac{g}{L} \cos \varphi = \frac{g}{L} \cos(n\pi) = \pm (-1)^n \frac{g}{L}$$



NOW SUPPOSE $Q(0) = Q_0$, $Q'(0) = 0$. THEN,

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$$\frac{Q''}{2} = \frac{g}{L} (\cos Q_0 - \cos Q)$$

NOW ON $Q' \neq 0$

$$Q' = -\sqrt{\frac{2g}{L}} [\cos Q_0 - \cos Q]^{1/2}$$

LET T BE THE PERIOD.

$$\frac{dQ}{dt} = -\sqrt{\frac{2g}{L}} (\cos Q_0 - \cos Q)^{1/2} \quad Q(0) = Q_0 \quad Q\left(\frac{T}{2}\right) = -Q_0$$

$$\text{THU} \quad \int_{Q_0}^{-Q_0} \frac{dQ}{[\cos Q - \cos Q_0]^{1/2}} = -\sqrt{\frac{2g}{L}} \int_0^{T/2} dt = -\frac{T}{2} \left(\frac{2g}{L}\right)^{1/2}.$$

$$\text{THU} \quad T = 2 \left(\frac{L}{2g}\right)^{1/2} \int_{-Q_0}^{Q_0} \frac{1}{[\cos Q - \cos Q_0]^{1/2}} dQ.$$

NOW LET $\cos Q = 1 - 2 \sin^2 Q/2$

$$\cos Q_0 = 1 - 2 \sin^2 Q_0/2$$

$$T = 4 \left(\frac{L}{2g}\right)^{1/2} \int_0^{Q_0} \frac{dQ}{\sqrt{2} [\sin^2 Q_0/2 - \sin^2 Q/2]^{1/2}}.$$

LET $\sin(Q/2) = \sin(Q_0/2) \sin \phi$

$$\frac{1}{2} \cos Q/2 dQ = \sin Q_0/2 \cos \phi d\phi$$

$$Q = 0 \rightarrow \phi = 0$$

$$Q = Q_0 \rightarrow \phi = \pi/2$$

$$T = 4 \left(\frac{L}{2g}\right)^{1/2} \int_0^{\pi/2} \frac{2 \sin Q_0/2 \cos \phi d\phi}{\cos Q_0/2 \sqrt{2} [\sin^2 Q_0/2 - \sin^2 Q_0/2 \sin^2 \phi]^{1/2}} d\phi$$

$$T = 4 \left(\frac{L}{2g}\right)^{1/2} \sqrt{2} \int_0^{\pi/2} \frac{\cos \phi d\phi}{\cos Q_0/2 \cos \phi d\phi}$$

$$\text{BUT} \quad \cos Q_0/2 = \sqrt{1 - K^2 \sin^2 \phi} \quad K = \sin Q_0/2.$$

$$\text{THU} \quad T = 4 \left(\frac{L}{g}\right)^{1/2} \int_0^{\pi/2} \left(\sqrt{1 - K^2 \sin^2 \phi} \right)^{-1} d\phi$$

$$K = \sin Q_0/2$$

EXPANSION

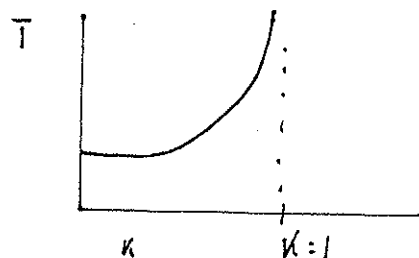
(4)

IN THE SMALL ANGLE APPROXIMATION $K \ll 1$

$$T \approx 4 \left(\frac{L}{g} \right)^{1/2} \int_0^{\pi/2} \left(1 + \frac{K^2}{2} \sin^2 \phi + \dots \right) d\phi$$

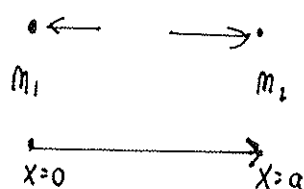
$$\text{so } T \approx 4 \left(\frac{L}{g} \right)^{1/2} \left[\frac{\pi}{2} + \frac{K^2}{2} \frac{\pi}{4} + \dots \right]$$

$$\text{THUS } T \approx 2\pi \left(\frac{L}{g} \right)^{1/2} \left[1 + \frac{K^2}{4} + \dots \right] \quad K = \sin \phi_0/2$$



EXAMPLE 2 PARTICLE BETWEEN TWO MASSES. (attractive)

$$m x'' = \frac{G m_2 m}{(x-a)^2} - \frac{G m_1 m}{x^2}$$



$$K_2 = G m_2$$

$$K_1 = G m_1$$

$$\text{THUS } x'' = \frac{K_2}{(x-a)^2} - \frac{K_1}{x^2}$$

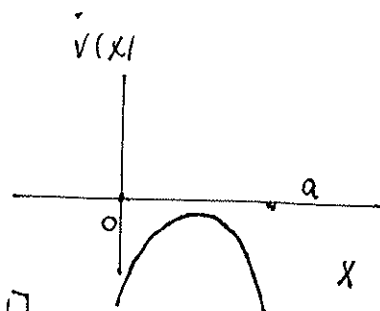
$$\text{NOW EQUILIBRIA IS } \frac{K_2}{K_1} = \left(\frac{x-a}{x} \right)^2 \quad \text{so } \frac{x-a}{x} = - \left(K_2/K_1 \right)^{1/2}$$

$$1 + \left(\frac{K_2}{K_1} \right)^{1/2} = \frac{a}{x}$$

$$\text{THUS } x_e = \frac{a}{1 + \left(K_2/K_1 \right)^{1/2}}$$

$$\text{NOW } V'(x) = \frac{K_1}{x^2} - \frac{K_2}{(x-a)^2}$$

$$V(x) = \frac{K_2}{x-a} - \frac{K_1}{x}$$

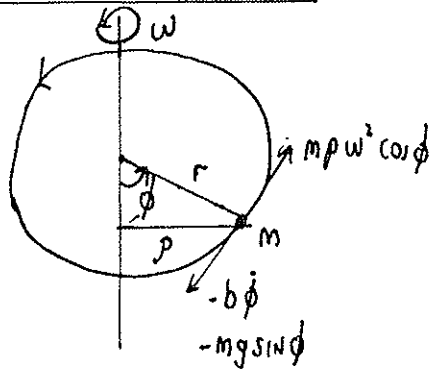


$$\text{NOW } V''(x) = -\frac{2K_1}{x^3} + \frac{2K_2}{(x-a)^3} = \frac{2}{x(x-a)} \left[\frac{K_2 x}{(x-a)^2} - \frac{K_1}{x^2} \right]$$

$$\text{BUT } \frac{K_2}{(x-a)^2} = \frac{K_1}{x^2} \quad \text{so } V''(x)_e = \frac{2}{x_e(x-a)_e} \left[\frac{K_1}{x_e} - \frac{K_1(x_e-a)}{x_e^2} \right] = \frac{2K_1 a}{x_e^3(x-a)_e} < 0$$

THUS x_e IS A MINIMUM $\rightarrow$ INSTABILITY.

EXAMPLE (BEAD ON A WIRE)



$$m r \ddot{\phi} = -b \dot{\phi} - mg \sin \phi + m \rho \omega^2 \cos \phi$$

$$\text{BUT } \rho = r \sin \phi$$

$$\text{SO } m r \ddot{\phi} = -b \dot{\phi} - mg \sin \phi + m r \omega^2 \cos \phi \sin \phi$$

ASSUME NO DAMPING $b = 0$. $t = T \tau$

$$\frac{m r}{T^2} \phi'' = -mg \sin \phi + m r \omega^2 \cos \phi \sin \phi$$

$$\frac{1}{T^2 \omega^2} \phi'' = -\frac{g}{r \omega^2} \sin \phi + \sin \phi \cos \phi$$

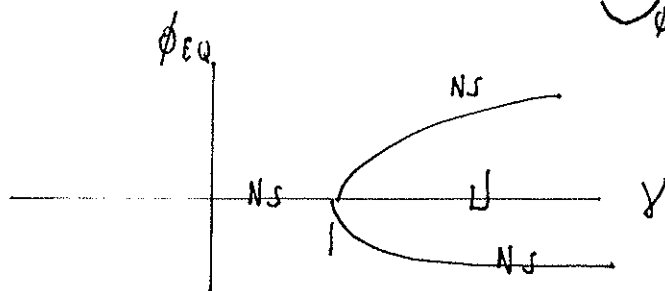
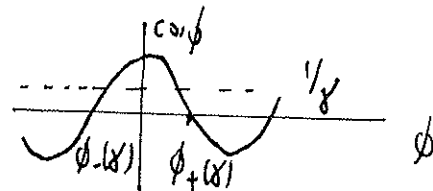
LET $T = 1/\omega$, $\gamma = r \omega^2 / g$.

$$\phi'' = -\frac{1}{\gamma} \sin \phi + \sin \phi \cos \phi$$

$$\text{NOW } F(\phi) = -\frac{1}{\gamma} \sin \phi + \frac{1}{2} \sin(2\phi)$$

EQUILIBRIA AT $\sin \phi = 0 \rightarrow \phi = 0, \pi$

$$\cos \phi = 1/\gamma \rightarrow \phi = \phi_c(\gamma), \gamma \geq 1$$



$$\text{NOW } V(\phi) = \frac{1}{\gamma} \int \sin u \, du - \frac{1}{2} \int \sin(2u) \, du$$

$$V(\phi) = -\frac{1}{\gamma} \cos(\phi) + \frac{1}{4} \cos 2\phi$$

$$V''(\phi) = \frac{1}{\gamma} \cos \phi - \cos 2\phi$$

$$V''(0) = \frac{1}{\gamma} - 1$$

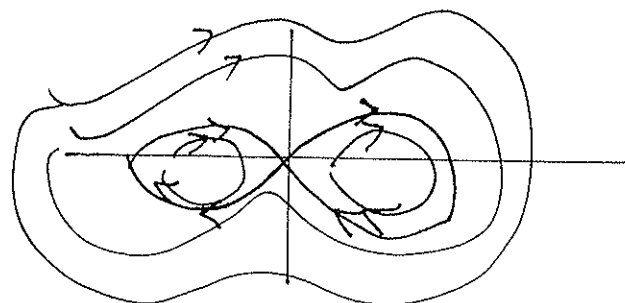
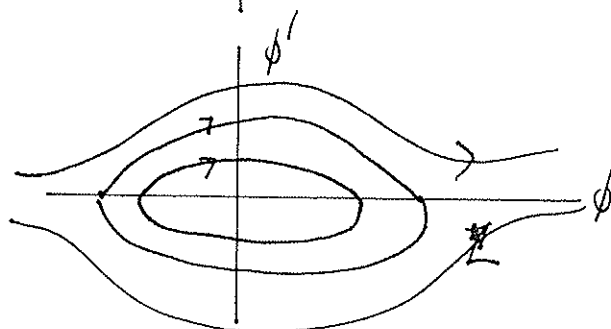
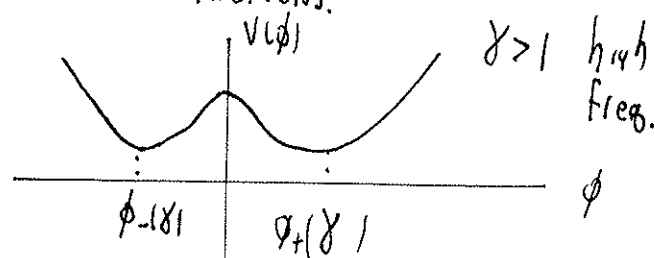
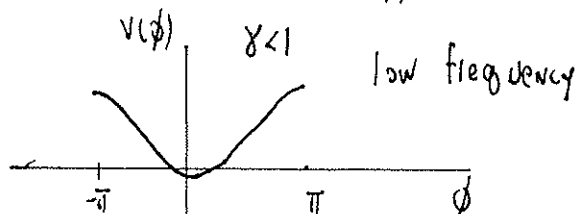
$$V''(0) > 0 \rightarrow \gamma < 1 \quad \text{stable oscillation (low frequency)}$$

$$V''(\pi) = -\frac{1}{\gamma} - 1 < 0$$

$$V''(\pi) < 0 \rightarrow \gamma > 1 \quad \text{exponential decay growth.}$$

now $V''(\phi_+) = \frac{1}{\gamma} \cos \phi_+ - (2 \cos^2 \phi_+ - 1) = \frac{1}{\gamma^2} - 2 \left(\frac{1}{\gamma^2} - \frac{1}{2} \right) = 1 - 1/\gamma^2.$

but if $\gamma > 1$ $V''(\phi_+) < 0 \rightarrow$ stable oscillations.



now if $\gamma > 1$ THEN $\omega = (1 - 1/\gamma^2)^{1/4}$ is frequency.

This implies that

$$T = \omega (1 - 1/\gamma^2)^{1/4}, \text{ SINCE } T = 2\pi/\omega.$$

EXAMPLE (DE LA YED BIFURCATION)

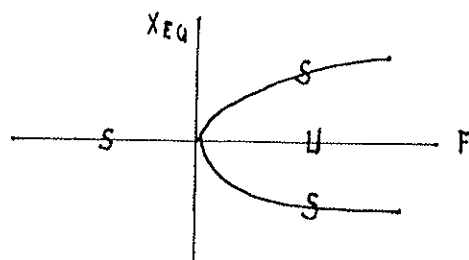
LET $X'' + KX' = -X(2X^2 - F) \quad K > 0$

$$X'(0) = 0, \quad X''(0) = X_0.$$

now DIAGRAM IS; EQUILIBRIA AT $X=0$ AND $X = \pm \sqrt{F/2}$ WHEN $F > 0$.

THE BIFURCATION DIAGRAM IS

(7)



$K=0$ BRANCHES are neutrally stable

$K>0$ Asymptotically stable

let $F = -1 + \epsilon t$ $\epsilon = .0025$ to see a delayed bifurcation.

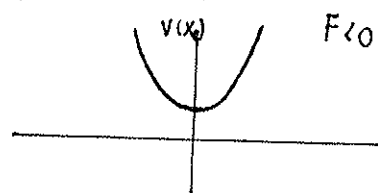
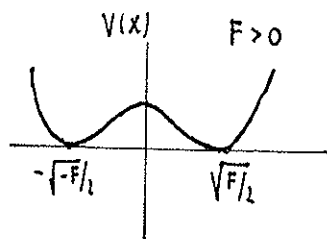
$$x'' + Kx' + V'(x) = 0$$

$$V'(x) = 2x^3 - Fx$$

$$V = \frac{x^4}{2} - \frac{Fx^2}{2} + C$$

THUS WE CAN WRITE

$$V(x) = \frac{1}{2} \left[x^2 - F/2 \right]^2$$



$$x'' + V'(x) = 0$$

$$\text{let } x = x_E + y$$

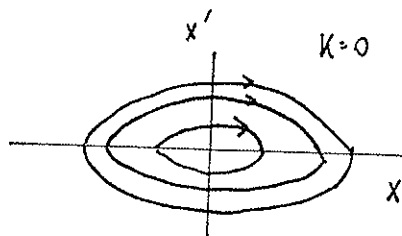
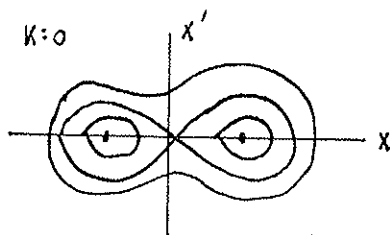
$$\text{stable if } V''(x_E) > 0.$$

$$y' = x''$$

$$y' = -V'(x)$$

$$J = \begin{pmatrix} 0 & 1 \\ -V''(x_E) & 0 \end{pmatrix}$$

$$\text{eigenvalue, } \lambda^2 + V''(x_E) = 0$$



NOW IF $K>0$,

$$x' x'' + x' V'(x) = -K x'^2$$

$$\frac{d}{dt} \left[\frac{1}{2} x'^2 + V(x) \right] = -K x'^2$$

let $E(x, x') = E(t) = \frac{1}{2} x'^2 + V(x)$. THEN $E \geq 0$ FOR ALL t .

$$\text{THIS EQUATION GIVES } \frac{dE}{dt} \leq 0 \quad E(0) = E_0 \geq 0, \quad E(t) \geq 0.$$

THUS $E \rightarrow 0$ AS $t \rightarrow \infty$.

NOW WRITE THIS AS $y = x'$

$$y' + Ky = xF - 2x^3$$

$$\text{so } \begin{cases} x' = y \\ y' = -Ky + xF - 2x^3 \end{cases}$$

WE NOW PLOT FOR $F > 0$.

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LET $X = X_E + \hat{X}$, $Y = Y_E + \hat{Y}$.

THEN $\hat{X}' = \hat{Y}$

$\hat{X}' = J \hat{X}$

$\hat{Y}' = (F - 6X_E^2) \hat{X} - K \hat{Y}$

$J = \begin{pmatrix} 0 & 1 \\ F - 6X_E^2 & -K \end{pmatrix}$

NOW $X_E = 0 \rightarrow$ eigenvalues are $-\lambda(-\lambda - K) - F = 0$

THUS $\lambda^2 + K\lambda - F = 0$ $\lambda_{\pm} = \frac{-K \pm \sqrt{K^2 + 4F}}{2}$

IF $F > 0 \rightarrow \lambda_+ > 0, \lambda_- < 0 \rightarrow X=0$ SADDLE POINT.

NOW $X_E = \pm \sqrt{F/2}$

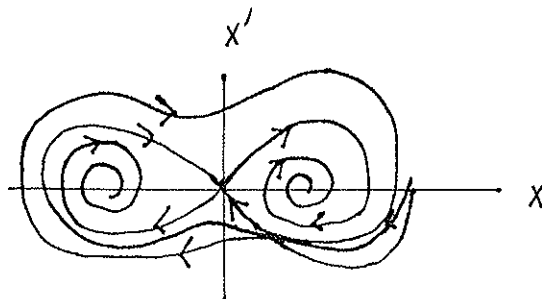
$J = \begin{pmatrix} 0 & 1 \\ -2F & -K \end{pmatrix}$

$-\lambda(-\lambda - K) + 2F = 0$

$\lambda^2 + \lambda K + 2F = 0$

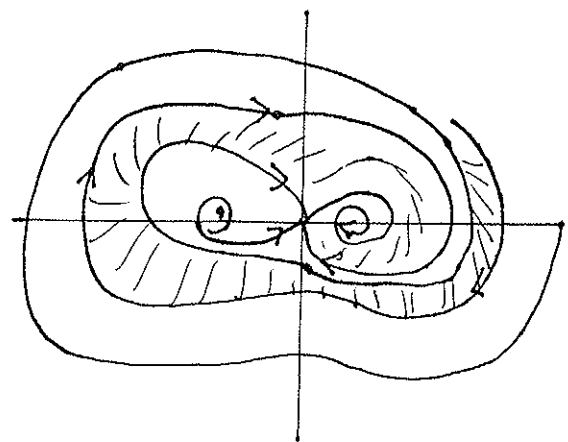
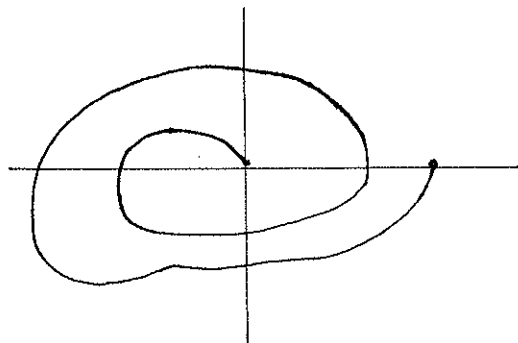
$\lambda_{\pm} = \frac{-K \pm \sqrt{K^2 - 8F}}{2}$

REAL $\lambda_{\pm}^2 < 0 \rightarrow$ stable.



THE BASIN OF $(\sqrt{F/2}, 0)$ IS AS SHOWN

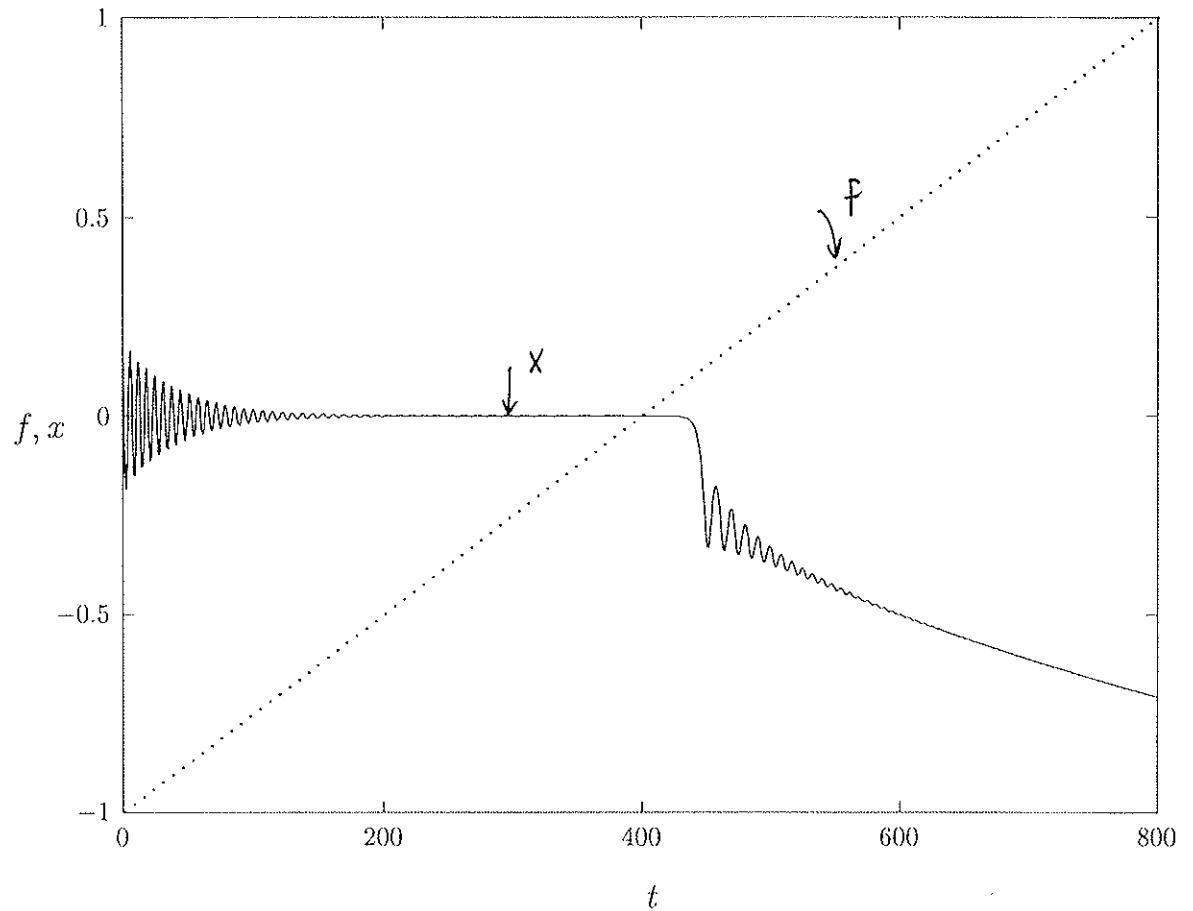
BASIN OF ATTRACTION



$$x'' + \kappa x' + x(2x^2 - f) = 0$$

$$f = -1 + \epsilon t$$

(1)



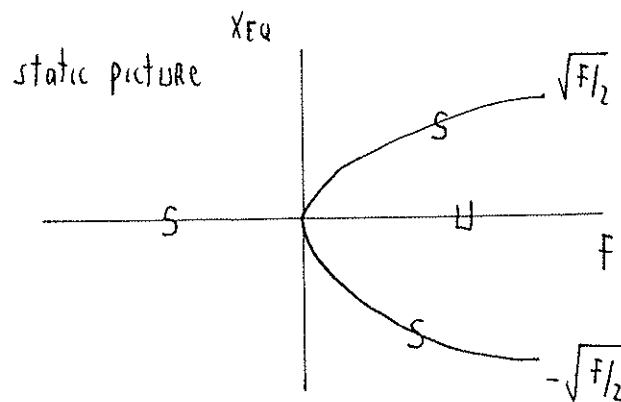
$$\epsilon = .0025$$

$$\kappa = 0.05$$

$$x(0) = 0.20$$

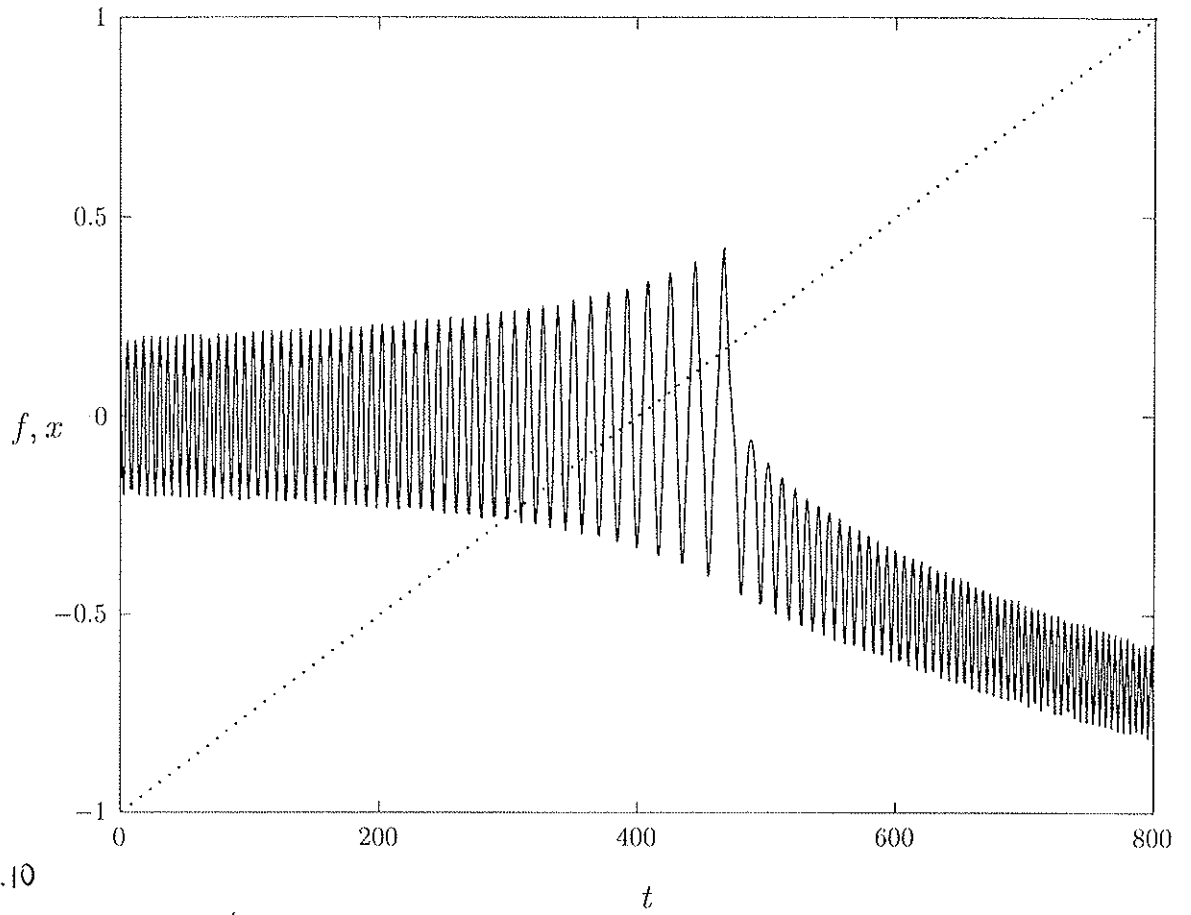
$$x'(0) = 0$$

transition based on static theory should
occur as f crosses through zero.

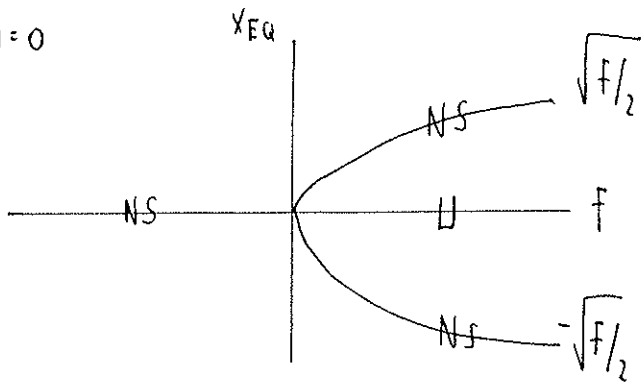


strong stability
since $\kappa > 0$.

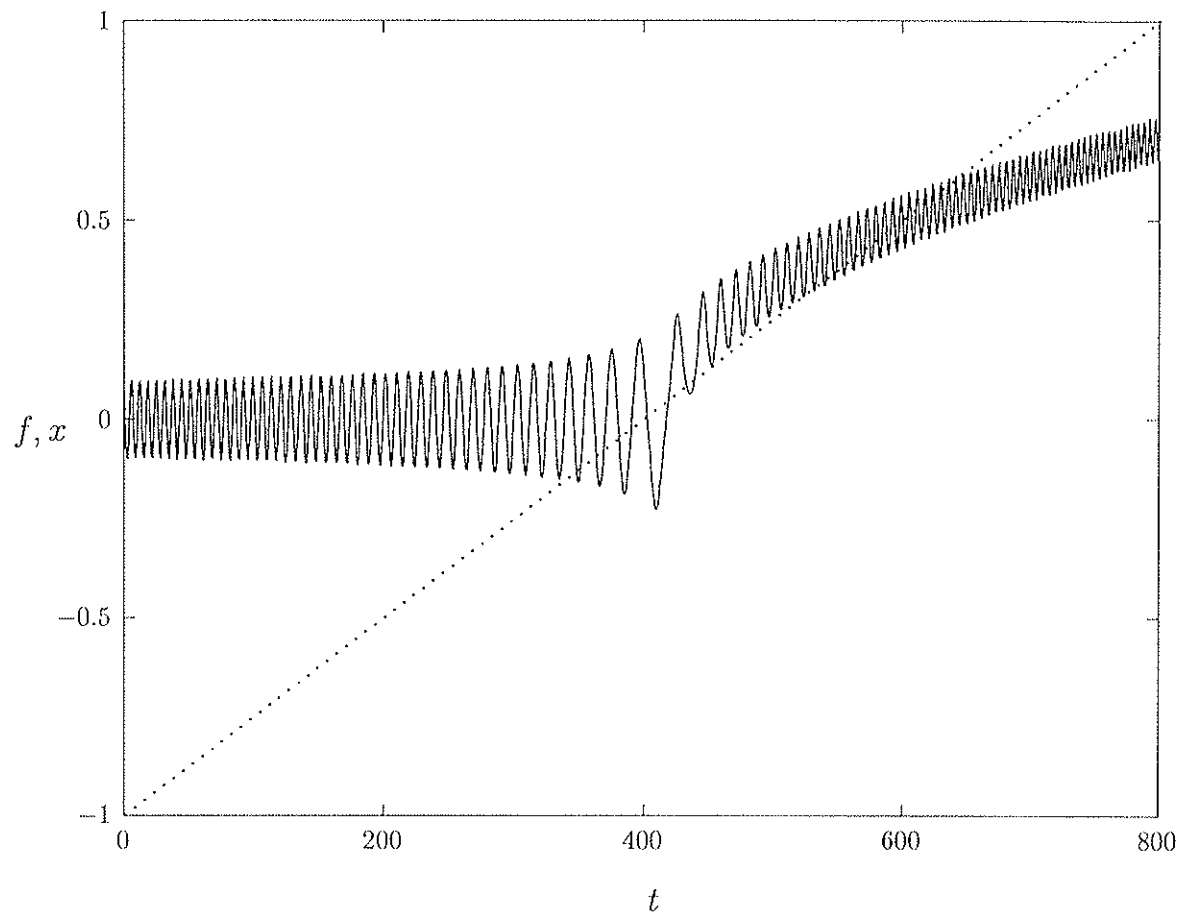
$$x'' + x(2x^2 - f) = 0 \quad f = -1 + \epsilon t$$



$$x(0) = -0.10$$
$$x'(0) = 0$$



neutrally stable
equilibrium branches
since no
damping.



with a different initial condition we get other point or $t \rightarrow \infty$

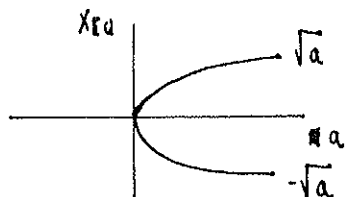
$$x(0) = 0.10$$

$$x'(0) = 0$$

EXAMPLE

$$x'' = a - x^2$$

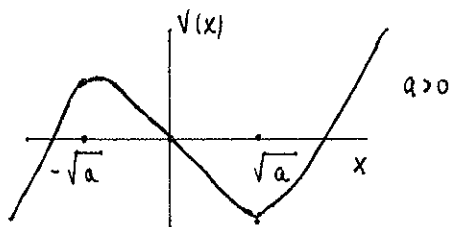
(9)



NOW $V'(x) = x^2 - a$ $V(x) = \frac{x^3}{3} - ax + C$

$$V(x) = \frac{x}{3} (x^2 - 3a)$$

$$V(\pm\sqrt{a}) = \pm \frac{\sqrt{a}}{3} (a - 3a) = \mp \frac{2a}{3} \sqrt{a}$$



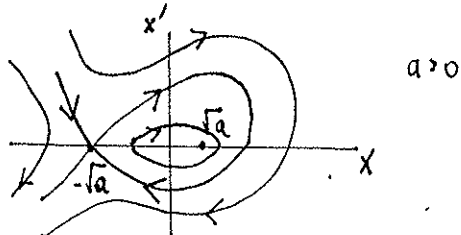
$$V''(x) = 2x$$

$$V''(\sqrt{a}) > 0 \rightarrow \text{center MINIMUM}$$

$$V''(-\sqrt{a}) < 0 \rightarrow \text{saddle point MAXIMUM}$$

$$V(\sqrt{a}) = -\frac{2a}{3}\sqrt{a} \quad V(-\sqrt{a}) = \frac{2a}{3}\sqrt{a}$$

THIS suggests a phase portrait as



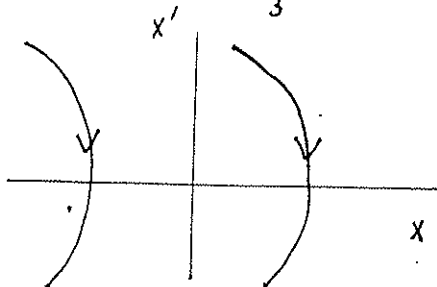
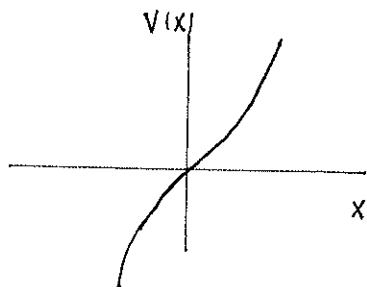
EQUATION OF SEPARATRIX.

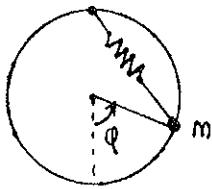
$$\frac{x'^2}{2} + V(x) = V(-\sqrt{a})$$

$$\frac{x'^2}{2} + \frac{x}{3} (x^2 - 3a) = \frac{2a}{3}\sqrt{a}$$

NOW IF $a < 0$ WE GET

$$V(x) = \frac{x}{3} (x^2 - 3a)$$





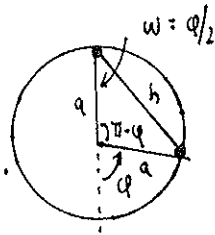
gravity

no damping

spring force: natural length = L_0

now gravity is $mg \sin \varphi$ tangent to the spring.

THU $m a \varphi'' = k [h - L] \sin(\varphi) - mg \sin \varphi$



$$h^2 = a^2 + a^2 - 2a^2 \cos(\pi - \varphi) = 2a^2 [1 + \cos \varphi]$$

$$\cos^2(2\varphi) = \cos^2 \varphi - \sin^2 \varphi = 2\cos^2 \varphi - 1$$

$$\cos^2 \varphi = \frac{1 + \cos 2\varphi}{2} \rightarrow h = \sqrt{2} a \sqrt{2} \cos(\varphi/2)$$

$$h = 2a \cos(\varphi/2)$$

THU $m a \varphi'' = k [2a \cos(\varphi/2) - L] \sin(\varphi/2) - mg \sin \varphi$

let $t = \tau T$.

$$\frac{m a}{T^2} \ddot{\varphi} = k [2a \cos(\varphi/2) - L] \sin(\varphi/2) - mg \sin \varphi$$

$$\frac{a}{g T^2} \ddot{\varphi} = \frac{2a k}{mg} [\cos(\varphi/2) - \gamma] \sin(\varphi/2) - \sin \varphi \quad \gamma = \frac{L}{2a}$$

choose $T = \sqrt{a/g}$.

$$\ddot{\varphi} = b [\cos(\varphi/2) - \gamma] \sin(\varphi/2) - \sin \varphi \quad \gamma = L/2a$$

$$b = 2ak/mg$$

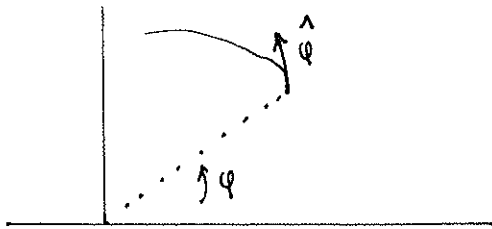
$\gamma > 1$ compression

$\gamma < 1$ tension

EXAMPLE (PLANETARY MOTION)

CENTRAL FORCE

(11)



$$\hat{r} = (\cos \varphi, \sin \varphi)$$

$$\hat{r} \cdot \hat{\varphi} = 0$$

$$\hat{\varphi} = (-\sin \varphi, \cos \varphi)$$

$$\dot{\hat{r}} = \dot{\varphi} \hat{\varphi}$$

$$\dot{\hat{\varphi}} = -\dot{\varphi} \hat{r}$$

NOW $x(t) = R(t) \cos(\varphi(t))$

$$y(t) = R(t) \sin(\varphi(t))$$

$$\dot{x} = \dot{R} \cos \varphi - R \sin \varphi \dot{\varphi}$$

$$\dot{y} = \dot{R} \sin \varphi + R \cos \varphi \dot{\varphi}$$

so $\underline{v} = \dot{R} \hat{r} + R \dot{\varphi} \hat{\varphi}$

$$\underline{\dot{v}} = \ddot{R} \hat{r} + \dot{R} \dot{\hat{r}} + (\dot{R} \dot{\varphi})' \hat{\varphi} + R \dot{\varphi} \dot{\hat{\varphi}} = \ddot{R} \hat{r} + (\dot{R} \dot{\varphi})' \hat{\varphi} + (R \dot{\varphi})' (-\hat{r}) + (R \dot{\varphi})' \hat{\varphi}$$

$$\underline{a} = (\ddot{R} - R \dot{\varphi}^2) \hat{r} + (R \ddot{\varphi} + 2 \dot{R} \dot{\varphi}) \hat{\varphi}$$

NOW SUPPOSE WE HAVE A CENTRAL FORCE $\underline{F} = (F_0, 0) = F_0 \hat{r}$

THEN $m \underline{a} = \underline{F}$

YIELDS,

$$m (\ddot{R} - R \dot{\varphi}^2) = F_0$$

$$R \ddot{\varphi} + 2 \dot{R} \dot{\varphi} = 0 \rightarrow R^2 \ddot{\varphi} + 2 R \dot{R} \dot{\varphi} = 0$$

$$\rightarrow (R^2 \dot{\varphi})' = 0$$

THU $R^2 \dot{\varphi} = L$ $L = \text{angular momentum.}$

$$\ddot{R} - \frac{R L^2}{R^4} = \frac{F_0(R)}{m}$$

DEFINE $V'(R) = -L^2/R^3 - 1/m F_0(R)$

NOW IF $F_0 = -\kappa R^p$ $\kappa > 0$ $p < -1$ (attractive force)

$$V(R) = \frac{L^2}{2R^2} + \frac{\kappa}{m(1+p)} R^{1+p}$$

NOW $V'(R) = -\frac{L^2}{R^3} + \frac{K}{m} R^p$

$V' > 0$ FOR $R > R_0$ $p > -3$

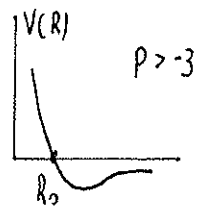
$V' < 0$ FOR $R < R_0$ $p < -3$

NOW $V' = 0$ WHEN $R_0 = \left(\frac{m L^2}{K} \right)^{1/p+3}$

THUS $V''(R_0) = \frac{3L^2}{R_0^4} + \frac{Kp}{m} R_0^{p-1} = \frac{1}{R_0^4} \left(3L^2 + \frac{Kp}{m} R_0^{p+3} \right) = \frac{1}{R_0^4} \left(3L^2 + \frac{Kp}{m} \frac{m L^2}{K} \right)$

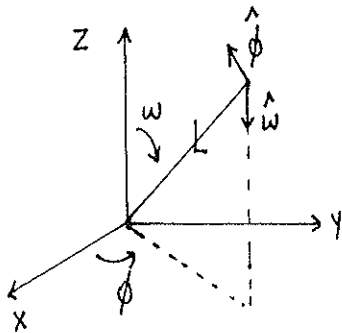
THIS GIVES $V''(R_0) = \frac{L^2}{R_0^4} (3+p)$ IF $p > -3$ $V''(R_0) > 0$ stability.

Frequency $\omega = \frac{L}{R_0^2} (3+p)^{1/2}$



SPHERICAL PENDULUM

RIGID ROD HAS LENGTH L



$$x = L \sin \omega \cos \phi$$

$$y = L \sin \omega \sin \phi$$

$$z = L \cos \omega$$

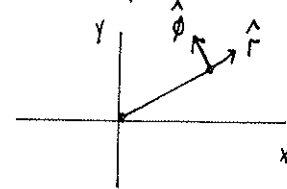
$$\dot{\hat{r}} = \dot{\phi} \hat{\phi}$$

$$\hat{\phi} = (-\sin \phi, \cos \phi, 0)$$

$$\hat{r} = (\cos \phi, \sin \phi, 0)$$

$$\hat{\omega} = (0, 0, 1)$$

IN $\omega = \pi/2$ PLANE



NOW $\underline{\dot{X}} = L \sin \omega \dot{\hat{r}} + L \cos \omega \dot{\hat{\omega}}$

$$\underline{\dot{V}} = L \cos \omega \dot{\omega} \hat{r} + L \sin \omega \dot{\phi} \hat{\phi} - L \sin \omega \dot{\omega} \hat{\omega}$$

$$\underline{\dot{V}} = L \cos \omega \dot{\omega} \hat{r} + L \sin \omega \dot{\phi} \hat{\phi} - L \sin \omega \dot{\omega} \hat{\omega}$$

NOW kinetic energy $K.E. = \frac{1}{2} m |\underline{\dot{V}}|^2 = \frac{1}{2} m L^2 \dot{\omega}^2 + \frac{1}{2} m L^2 \sin^2 \omega \dot{\phi}^2$

$$P.E. = mg L \cos \omega$$

HENCE $L = K.E. - P.E. = \frac{1}{2} m L^2 \dot{\omega}^2 + \frac{1}{2} m L^2 \sin^2 \omega \dot{\phi}^2 - mg L \cos \omega$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\phi}} \right) - \frac{\partial L}{\partial \phi} = 0 \rightarrow \frac{d}{dt} (m L^2 \sin^2 \omega \dot{\phi}) = 0$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\omega}} \right) - \frac{\partial L}{\partial \omega} = 0 \rightarrow \frac{d}{dt} (m L^2 \dot{\omega}) - m L^2 \dot{\phi}^2 \sin \omega \cos \omega - m g L \sin \omega = 0$$

THUS we get 2 equations:

$$\left\{ \begin{array}{l} \ddot{\omega} - \dot{\phi}^2 \sin \omega \cos \omega - \frac{g}{L} \sin \omega = 0 \\ \dot{\phi} = h / \sin^2 \omega \end{array} \right. \quad h = \text{CONSTANT}$$

RE-WRITING: $\ddot{\omega} - h^2 \cot \omega \csc^3 \omega - \frac{g}{L} \sin \omega = 0.$

$$\dot{\phi} = h / \sin^2 \omega$$

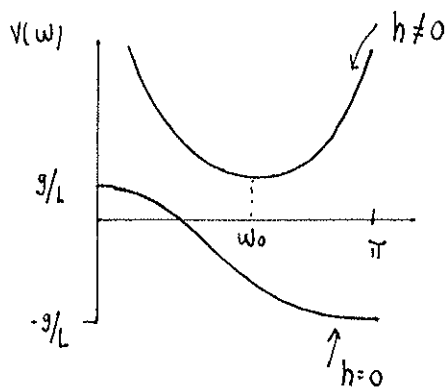
NOW DEFINE THE POTENTIAL BY

$$V'(\omega) = -h^2 \cot \omega \csc^3 \omega - \frac{g}{L} \sin \omega$$

$$V(\omega) = \frac{h^2}{2} \csc^2 \omega + \frac{g}{L} \cos(\omega) \quad 0 \leq \omega \leq \pi$$

HENCE

$$\ddot{\omega} + V'(\omega) = 0.$$



HENCE IF $h=0$ THEN $\omega=0$ UNSTABLE (top of motion)
 $\omega=\pi$ STABLE (bottom of sphere)

NOTICE h IS "angular momentum" IN plane of constant ω . $\rightarrow$ how fast it is whirling around. Thus if $h=0$ (bottom of sphere and top of sphere are stable and unstable respectively).

NOW IF $h>0$ we only have one equilibrium point at ω_0 . IT WILL BE STABLE AND CORRESPONDS TO A WHIRLING MOTION IN THE FIXED LATITUDE PLANE

NOTICE $V'(\omega_0) = 0 \rightarrow \frac{L h^2}{g} = \frac{\sin^4(\omega_0)}{\cos(\omega_0)}$

NOTICE THAT AS $h \rightarrow \infty$ THEN $\omega_0 \rightarrow \pi/2$. WE expect that
 THE inequality $\omega_0 > \pi/2$ (IN southern hemisphere) should hold for any $h > 0$.

NOW SUPPOSE h IS LARGE. Let

$$t = T\tau.$$

$$\frac{1}{T^2} \omega'' - h^2 \cot \omega \csc^3 \omega - \frac{g}{L} \sin \omega = 0$$

$$\frac{1}{T} \phi' = h / \sin^2 \omega$$

$$\omega' = d\omega/d\tau$$

CHOOSE $T = 1/h$. THEN we get the approximate system:

$$\omega'' - \cot \omega \csc^3 \omega = 0$$

$$\phi' = 1 / \sin^2 \omega$$

$$V(\omega) = \frac{1}{2} \csc^2 \omega.$$

