

PROBLEM 1

(i) $Lu = u'' + \lambda u = f(x)$ WITH $u'(0) = u'(L) = 0$ IS SELF-ADJOINT.

THE HOMOGENEOUS PROBLEM $\phi'' + \lambda \phi = 0$, $\phi'(0) = \phi'(L) = 0$ WITH $\lambda \geq 0$

HAS NONTRIVIAL SOLUTION $\phi_0 = 1$ WHEN $\lambda = 0$

$$\phi_n = \cos\left(\frac{n\pi x}{L}\right) \text{ WHEN } \lambda_n = \frac{n^2 \pi^2}{L^2}, n=1,2,\dots$$

THUS WHEN $\lambda = 0$, THE SOLVABILITY CONDITION IS $\int_0^L f(x) dx = 0$

WHEN $\lambda_n = \frac{n^2 \pi^2}{L^2}$, $n=1,2,\dots$ THE SOLVABILITY CONDITION IS $\int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx = 0$.

(ii) SUPPOSE $\lambda > 0$ AND $\lambda \neq \lambda_n$ FOR ANY $n=1,2,\dots$

THE GREEN'S FUNCTION $v(t;x)$ SATISFIES

$$v'' + \lambda v = \delta(t-x) \text{ IN } 0 < t < L$$

$$v'(0) = v'(L) = 0.$$

THE HOMOGENEOUS PROBLEM HAS $v = \text{SPAN} \{ \cos(\sqrt{\lambda}(x-\phi)), \sin(\sqrt{\lambda}(t+\phi)) \}$.

TO MAKE $v' = 0$ AT $t=0, L$, WE HAVE

$$v = \begin{cases} A \cos(\sqrt{\lambda} t), & 0 \leq t \leq x \\ B \cos(\sqrt{\lambda}(t-L)), & \text{ON } x \leq t \leq L \end{cases}$$

NOW CONTINUITY AT $x=t$ HOLDS IF $v = \begin{cases} C \cos(\sqrt{\lambda}(x-L)) \cos(\sqrt{\lambda} t), & 0 \leq t \leq x \\ C \cos(\sqrt{\lambda}(t-L)) \cos(\sqrt{\lambda} x), & x \leq t \leq L \end{cases}$

THE JUMP CONDITION $v'|_{x^+} - v'|_{x^-} = 1$ REQUIRES

$$\sqrt{\lambda} C [-\sin(\sqrt{\lambda}(x-L)) \cos(\sqrt{\lambda} x) + \cos(\sqrt{\lambda}(x-L)) \sin(\sqrt{\lambda} x)] = 1$$

BUT $\sin(A-B) = \sin A \cos B - \cos A \sin B$ SO THAT

$$-\sqrt{\lambda} C \sin(-\sqrt{\lambda} L) = 1 \rightarrow C = \frac{1}{\sqrt{\lambda} \sin(\sqrt{\lambda} L)}$$

HENCE

$$v = \frac{1}{\sqrt{\lambda} \sin(\sqrt{\lambda} L)} \begin{cases} \cos(\sqrt{\lambda}(x-L)) \cos(\sqrt{\lambda} t), & 0 \leq t \leq x \\ \cos(\sqrt{\lambda}(t-L)) \cos(\sqrt{\lambda} x), & x \leq t \leq L \end{cases}$$

THE INTEGRAL REPRESENTATION IS OBTAINED FROM

$$\int_0^L v Lu dt = (v u' - u v') \Big|_0^L + \int_0^L u Lv dt$$

$\leftarrow = 0 \rightarrow$

THUS,

$$u(x) = \int_0^L v(t;x) f(t) dt$$

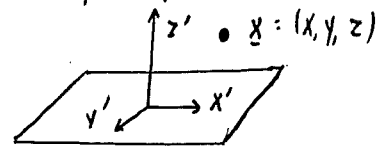
(iii) FOR $\lambda = 0$, THE MODIFIED GREEN'S FUNCTION IS

$$v'' = \delta(t-x) - \frac{1}{L} \text{ ON } 0 < t < L; v'(0) = v'(L) = 0.$$

PROBLEM 2 SUPPOSE THAT

$$\Delta u = 0 \text{ in } -\infty < x < \infty, -\infty < y < \infty, z > 0$$

WITH $u_z(x, y, 0) = f(x, y)$ AND $u \sim C/r$ AS $r = (x^2 + y^2 + z^2)^{1/2} \rightarrow \infty$.



(i) NOW THE G-FUNCTION $G(\underline{x}'; \underline{x})$ SATISFIES

$$\begin{cases} \Delta' G = \delta(\underline{x}' - \underline{x}) & \text{IN } -\infty < x', y' < \infty, z' > 0 \\ G_{z'} = 0 & \text{ON } z' = 0. \end{cases}$$

• $\underline{x}_I = (x, y, -z)$

NOW THE FREE SPACE G-FUNCTION IS $G^f = -\frac{1}{4\pi |\underline{x}' - \underline{x}|}$

BY METHOD OF IMAGES WE NEED AN IMAGE CHARGE AT $\underline{x}_I = (x, y, -z)$

WITH SAME SIGN.

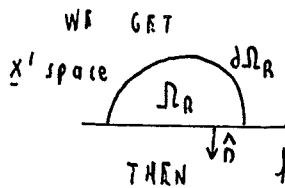
THUS $G(\underline{x}'; \underline{x}) = -\frac{1}{4\pi |\underline{x}' - \underline{x}|} - \frac{1}{4\pi |\underline{x}' - \underline{x}_I|}$

$$|\underline{x}' - \underline{x}| = \sqrt{(x' - x)^2 + (y' - y)^2 + (z' - z)^2}$$

$$|\underline{x}' - \underline{x}_I| = \sqrt{(x' - x)^2 + (y' - y)^2 + (z' + z)^2}$$

THIS SATISFIES $G_{z'} = 0$ ON $z' = 0$.

(ii) NOW USE GREEN'S SECOND IDENTITY BY INTEGRATING OVER A HEMI-SPHERE OF RADIUS R .



WE GET $\Delta' u = 0$ IN $z' \geq 0$
 $u_{z'} = f(x', y')$ ON $z' = 0$

THEN $\lim_{R \rightarrow \infty} \int_{\Omega_R} (G \Delta' u - u \Delta' G) d\underline{x}' = \lim_{R \rightarrow \infty} \int_{\partial \Omega_R} \left(G \frac{\partial u}{\partial n} \Big|_{r=R} - u \frac{\partial G}{\partial n} \Big|_{r=R} \right) ds$
 $+ \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left(G \frac{\partial u}{\partial n} - u \frac{\partial G}{\partial n} \right) \Big|_{z'=0} dx' dy'$

NOW ON $\partial \Omega_R$, $ds = O(R^2)$ AND $G \sim O(1/R)$

WHILE $u \sim O(1/R)$. HENCE $\lim_{R \rightarrow \infty} \int_{\partial \Omega_R} \left(G \frac{\partial u}{\partial n} \Big|_{r=R} - u \frac{\partial G}{\partial n} \Big|_{r=R} \right) ds = 0$.

NOW $\partial G / \partial n = -\partial G / \partial z'$ ON $z' = 0$, AND $\partial u / \partial n = -\partial u / \partial z'$ ON $z' = 0$.

HENCE, $\lim_{R \rightarrow \infty} \int_{\Omega_R} -u \delta(\underline{x}' - \underline{x}) d\underline{x}' = - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} G \frac{\partial u}{\partial z'} \Big|_{z'=0} dx' dy'$, WITH $\frac{\partial u}{\partial z'} \Big|_{z'=0} = f(x', y')$.

THEREFORE, $u(\underline{x}) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} G \Big|_{z'=0} f(x', y') dx' dy'$

NOW $G \Big|_{z'=0} = -\frac{1}{2\pi [(x' - x)^2 + (y' - y)^2 + z^2]^{1/2}}$ FROM PART (i).

WE CONCLUDE THAT $u(\underline{x}) = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{f(x', y')}{[(x' - x)^2 + (y' - y)^2 + z^2]^{1/2}} dx' dy'$.

(iii) NOW LET $\underline{x}' = (x', y', 0)$ AND $\underline{x} = (x, y, z)$.

FOR $|\underline{x}| \rightarrow \infty$ WE USE $|\underline{x}' - \underline{x}|^2 = |\underline{x}|^2 + |\underline{x}'|^2 - 2\underline{x} \cdot \underline{x}'$
 $|\underline{x}' - \underline{x}|^2 = |\underline{x}|^2 \left(1 - \frac{2\underline{x} \cdot \underline{x}'}{|\underline{x}|^2} + \dots \right)$ FOR $|\underline{x}|$ LARGE.

NOW $|\underline{x}' - \underline{x}| \cong |\underline{x}| \left(1 - \frac{2\underline{x} \cdot \underline{x}'}{|\underline{x}|^2} \right)^{1/2}$

SO $\frac{1}{|\underline{x}' - \underline{x}|} \cong \frac{1}{|\underline{x}|} \left(1 - \frac{2\underline{x} \cdot \underline{x}'}{|\underline{x}|^2} \right)^{-1/2} \cong \frac{1}{|\underline{x}|} + \frac{\underline{x} \cdot \underline{x}'}{|\underline{x}|^3} + \dots$

USING $(1+h)^{-1/2} \sim 1 + h/2$.

NOW THIS GIVES

$$u(\underline{x}) = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{f(x', y')}{[|\underline{x}' - \underline{x}|^2]^{3/2}} dx' dy' \cong -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left(\frac{1}{|\underline{x}|} + \frac{\underline{x} \cdot \underline{x}'}{|\underline{x}|^3} \right) f(x', y') dx' dy'$$

THIS GIVES THAT

$$u(\underline{x}) \sim \left(-\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x', y') dx' dy' \right) \frac{1}{|\underline{x}|} + \frac{P \cdot \underline{x}}{|\underline{x}|^3} + \dots$$

WHERE WE IDENTIFY

$$C = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x', y') dx' dy'$$

$$P = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \underline{x}' f(x', y') dx' dy'$$

WHERE $\underline{x}' = (x', y', 0)$.

REMARK WE COULD ALSO FIND C BY DIVERGENCE THEOREM, INTEGRATING

OVER A HEMI-SPHERE Ω_R WE OBTAIN

$$\lim_{R \rightarrow \infty} \int_{\Omega_R} \Delta u d\underline{x} = \lim_{R \rightarrow \infty} \int_{\partial \Omega_R} \nabla u \cdot \hat{n} dS + \lim_{R \rightarrow \infty} \int_{-R}^R \int_{-R}^R \left(-\frac{\partial u}{\partial z'} \right) \Big|_{z'=0} dx' dy'$$

BUT $\Delta u = 0$ AND $\nabla u \cdot \hat{n} = \partial u / \partial r$ WITH $u \sim C/r$ AS $r \rightarrow \infty$.

HENCE $0 = 2\pi \left(-C/R \right) R^2 + \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} (-f(x', y')) dx' dy'$

WE CONCLUDE THAT $C = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x', y') dx' dy'$.

PROBLEM 3

CONSIDER $U_{xx} + U_{yy} + U_{zz} = \delta(x-x_0)\delta(y-y_0)\delta(z-z_0)$ IN $0 < z < \pi$, $-\infty < x, y < \infty$
 $U = 0$ ON $z = 0, \pi$ AND $U \rightarrow 0$ AS $(x^2 + y^2 + z^2)^{1/2} \rightarrow \infty$.

(i) WE EXPAND $U(x, y, z) = \sum_{n=1}^{\infty} C_n(x, y) \sin(nz)$

AS $\{\sin(nz)\}$ ARE THE EIGENFUNCTIONS IN THE z DIRECTION.

BY ORTHOGONALITY $C_n = \frac{2}{\pi} \int_0^{\pi} U \sin(nz) dz$.

MULTIPLY THE PDE BY $\sin(nz)$ AND INTEGRATE.

$$C_{n,xx} + C_{n,yy} + \frac{2}{\pi} \int_0^{\pi} U_{zz} \sin(nz) dz = \frac{2}{\pi} \delta(x-x_0)\delta(y-y_0) \sin(nz_0)$$

INTEGRATING BY PARTS

$$C_{n,xx} + C_{n,yy} + \frac{2}{\pi} \left[\underbrace{U_z \sin(nz)}_{\leftarrow = 0 \rightarrow} - \eta U \cos(nz) \right] \Big|_0^{\pi} = \frac{2}{\pi} \eta^2 \int_0^{\pi} U \sin(nz) dz$$

$$= \frac{2}{\pi} \delta(x-x_0)\delta(y-y_0) \sin(nz_0)$$

THEFORE $\Delta_0 C_n - \eta^2 C_n = \frac{2}{\pi} \delta(x-x_0) \sin(nz_0) \quad (1)$

IS THE PDE FOR $C_n(x, y)$. HERE $\Delta_0 \equiv \partial_{xx} + \partial_{yy}$, $\underline{x} = (x, y)$, $\underline{x}_0 = (x_0, y_0)$

(ii) NOW RECALL THAT $\Delta v^f - \eta^2 v^f = \delta(\underline{x} - \underline{x}_0)$ IN $\mathbb{R}^2$
 HAS SOLUTION $v^f = -\frac{1}{2\pi} H_0(\eta |\underline{x} - \underline{x}_0|)$.

COMPARING WITH (1) WE CONCLUDE THAT

$$C_n(\underline{x}) = \frac{2}{\pi} \sin(nz_0) H_0(\eta |\underline{x} - \underline{x}_0|) \left(-\frac{1}{2\pi} \right) = -\frac{1}{\pi^2} H_0(\eta |\underline{x} - \underline{x}_0|) \sin(nz_0).$$

WHERE $H_0(z)$ IS THE MODIFIED BESSEL FUNCTION

PROBLEM 4

(i) $\delta(x^2 - 5x + 4) = \delta((x-4)(x-1)).$

THUS CONCENTRATION POINTS ARE $x=1, x=4.$

RECALL $\delta(f(x)) = \sum_i \frac{\delta(x-x_i)}{|f'(x_i)|}$

SO THAT IF $f(x) = x^2 - 5x + 4$ THEN $f'(1) = 2 \cdot 1 - 5 = -3, f'(4) = (2x-5)|_{x=4} = 3.$

THUS $\delta(x^2 - 5x + 4) = \frac{1}{3} \delta(x-4) + \frac{1}{3} \delta(x-1).$

(ii) DEFINE $f_\sigma(x) = \frac{\sigma}{(x-1)^2 + \sigma^2}$. THEN $\lim_{\substack{\sigma \rightarrow 0 \\ x \neq 1}} f_\sigma(x) = 0, \lim_{\substack{\sigma \rightarrow 0 \\ x=1}} f_\sigma(x) = \infty.$

NOW WE INTEGRATE $\int_{-\infty}^{\infty} \frac{\sigma}{(x-1)^2 + \sigma^2} dx = \frac{\sigma^2}{\sigma^2} \int_{-\infty}^{\infty} \frac{1}{y^2 + 1} dy = \text{ARCTAN } y \Big|_{-\infty}^{\infty} = \frac{\pi}{2} - \left(-\frac{\pi}{2}\right) = \pi.$

let $x-1 = \sigma y$

WE CONCLUDE THAT $\lim_{\substack{\sigma \rightarrow 0 \\ x \neq 1}} f_\sigma(x) = \pi \delta(x-1).$

(iii) NOW $u_{rr} + \frac{2}{r} u_r - u = f(r)$ ON $1 \leq r \leq 2, u(1) = u(2) = 0.$

MULTIPLY BY r^2 TO OBTAIN.

$$r^2 u_{rr} + 2 u_r - r^2 u = r^2 f(r)$$

$$\rightarrow (r^2 u_r)_r - r^2 u = r^2 f(r)$$

$$Lu \equiv [(r^2 u_r)_r - r^2 u] \text{ IS SELF-ADJOINT.}$$

(iv) RECALL IF $\Delta u = \delta(\underline{x} - \underline{x}_0)$ IN 2-D, THEN $u \sim \frac{1}{2\pi} \log |\underline{x} - \underline{x}_0|.$

NOW TAKE GRADIENT, IT COMMUTES WITH LAPLACIAN SO THAT

$$\Delta(\nabla u) = \nabla_x \delta(\underline{x} - \underline{x}_0).$$

NOW DOT PRODUCT WITH $\underline{p}$ $\Delta(\nabla u) \cdot \underline{p} = \nabla_x \delta(\underline{x} - \underline{x}_0) \cdot \underline{p}.$

SINCE $\underline{p}$ CONSTANT VECTOR $\Delta(\nabla u \cdot \underline{p}) = \nabla_x \delta(\underline{x} - \underline{x}_0) \cdot \underline{p}.$

THUS $\underline{v} = \nabla_x u \cdot \underline{p} = \frac{1}{2\pi} \nabla_x \log |\underline{x} - \underline{x}_0| \cdot \underline{p} = \frac{1}{2\pi} \frac{(\underline{x} - \underline{x}_0) \cdot \underline{p}}{|\underline{x} - \underline{x}_0|^2}.$

(NOTE: $\nabla_x \log |\underline{x} - \underline{x}_0| = \frac{1}{2} \nabla_x \log((x-x_0)^2 + (y-y_0)^2) = \frac{(x-x_0, y-y_0)}{|\underline{x} - \underline{x}_0|^2} = \frac{\underline{x} - \underline{x}_0}{|\underline{x} - \underline{x}_0|^2}.$