

# SINGULAR PERTURBATIONS I

WE FIRST CONSIDER SINGULAR PERTURBATIONS FOR THE ROOTS OF POLYNOMIALS. FOR INSTANCE CONSIDER

$$\varepsilon x^4 - x^2 - x + 2 = 0.$$

FIND A 2-TERM EXPANSION FOR EACH OF THE ROOTS AS  $\varepsilon \rightarrow 0$ .

$$\varepsilon x^4 - (x^2 + x - 2) = \varepsilon x^4 - (x+2)(x-1) = 0.$$

WE WILL HAVE ONE ROOT NEAR  $x=1$  AND ANOTHER NEAR  $x=-2$ .

ROOT NEAR  $x=1$ :  $x = 1 + \varepsilon x_1 + \dots$  SO THAT

$$\varepsilon (1 + \varepsilon x_1)^4 - (3 + \varepsilon x_1)(\varepsilon x_1) = \varepsilon (1 - 3x_1) + O(\varepsilon^2) = 0$$

$$\text{so } x_1 = 1/3 \quad x = 1 + \varepsilon/3 + \dots$$

ROOT NEAR  $x=-2$  LET  $x = -2 + \varepsilon x_1 + \dots$

$$\varepsilon (-2 + \varepsilon x_1)^4 - \varepsilon x_1 (-3 + \varepsilon x_1) = \varepsilon (3x_1 - 16) + O(\varepsilon^2) = 0$$

$$\text{so } x_1 = 16/3 \quad x = -2 + 16\varepsilon/3 + \dots$$

FOR THE OTHER TWO ROOTS LET  $x = \varepsilon^{-p} y$   $p > 0$ .

$$\text{THEN } \varepsilon^{1-4p} y^4 - (\varepsilon^{-2p} y^2 + \varepsilon^{-p} y - 2) = 0.$$

$$\text{WE BALANCE } \varepsilon^{1-4p} = \varepsilon^{-2p} \quad \text{so } 1-4p = -2p \rightarrow p = \frac{1}{2}.$$

$$\text{THEN } y^4 - (y^2 + \varepsilon^{1/2} y - 2\varepsilon) = 0.$$

$$\text{so } y^4 - y^2 = 0 \quad \text{OR} \quad y^2 = 1 \rightarrow y = \pm 1.$$

$$\bullet \text{ WE LET } y = 1 + \varepsilon^{1/2} y_1 + \dots$$

$$(1 + \varepsilon^{1/2} y_1)^4 - (1 + \varepsilon^{1/2} y_1)^2 - \varepsilon^{1/2} (1 + \dots) + O(\varepsilon) = 0$$

$$\rightarrow 4 \varepsilon^{1/2} y_1 - 2 \varepsilon^{1/2} y_1 - \varepsilon^{1/2} = 0$$

$$2 y_1 - 1 = 0 \quad y_1 = 1/2.$$

$$\text{so } y = 1 + \varepsilon^{1/2}/2 \rightarrow x = \varepsilon^{-1/2} + 1/2 + \dots$$

$$\bullet \text{ WE LET } y = -1 + \varepsilon^{1/2} y_1 + \dots$$

$$(-1 + \varepsilon^{1/2} y_1)^4 - (-1 + \varepsilon^{1/2} y_1)^2 - \varepsilon^{1/2} (-1 + \varepsilon^{1/2} y_1) + O(\varepsilon) = 0$$

$$-4 \varepsilon^{1/2} y_1 - (-2 \varepsilon^{1/2} y_1) + \varepsilon^{1/2} = 0$$

$$-2 \varepsilon^{1/2} y_1 + \varepsilon^{1/2} = 0 \quad y_1 = 1/2$$

$$\text{so } y = -1 + \varepsilon^{1/2}/2 + \dots$$

$$x = -\varepsilon^{-1/2} + 1/2 + \dots$$

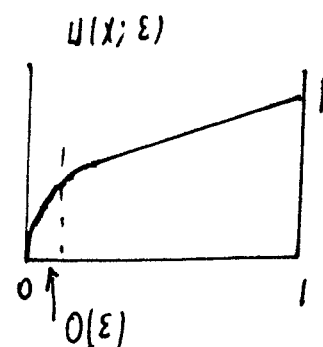
THE ROOTS ARE:

$$\left\{ \begin{array}{l} x = 1 + \varepsilon/3 + \dots \\ x = -2 + 16\varepsilon/3 + \dots \\ x_{\pm} = \pm \varepsilon^{-1/2} + 1/2 + \dots \end{array} \right.$$

EXAMPLE 1 FIND THE EXACT AND A BOUNDARY LAYER SOLUTION FOR

$$\varepsilon u'' + u' = \frac{1}{2}, \quad 0 \leq x \leq 1 \quad \varepsilon \rightarrow 0^+$$

$$u(0) = 0, \quad u(1) = 1$$



THE EXACT SOLUTION IS READILY FOUND TO BE

$$u(x; \varepsilon) = \frac{1 - e^{-x/\varepsilon}}{2(1 - e^{-1/\varepsilon})} + \frac{x}{2}$$

NOTICE THAT  $\lim_{\varepsilon \rightarrow 0} u(x; \varepsilon) = \frac{1}{2} + \frac{x}{2}$  OUTER LIMIT

$x \neq 0$

$\lim_{\varepsilon \rightarrow 0} u(x; \varepsilon) = \frac{1 - e^{-\gamma}}{2}$  INNER LIMIT.

$\gamma = x/\varepsilon$  FIXED

WE NOW DEVELOP A BOUNDARY LAYER THEORY TO CONSTRUCT THE SOLUTION.

AWAY FROM  $x = 0$  WE EXPAND

$$u_{out} = u_0 + \varepsilon u_1 + \varepsilon^2 u_2 + \dots$$

COLLECTING POWERS OF  $\varepsilon$  WE GET

$$u_0' = \frac{1}{2} \rightarrow u_0 = \frac{x + C_0}{2}$$

$$u_1' = -u_0'' \rightarrow u_1 = C_1$$

WE CAN'T SATISFY BOTH BOUNDARY CONDITIONS AT  $x = 0$ , AND  $x = 1$ . ASSUME THAT THERE IS NO BL NEAR  $x = 1$  AND SET  $u_{out}(1) = 1$ . THIS GIVES  $u_0(1) = 1$  AND  $u_1(1) = 0$ . WE OBTAIN  $C_1 = 0$  AND  $C_0 = 1$

SO THAT  $u_0 = \frac{(1+x)}{2}, \quad u_1 = 0.$

THUS  $u_{out} = (1+x)/2 + O(\varepsilon^2).$

NOW SINCE  $u_{out}(0) = 1/2 \neq 0$  WE MUST INSERT A BL NEAR  $x=0$  (4)  
WHERE  $u$  CHANGES RAPIDLY AND  $\varepsilon u''$  IS NO LONGER SMALL.

WE RESCALE NEAR  $x=0$  BY LETTING  $y = x/\varepsilon^p$  AND  $v(y) = u(\varepsilon^p y)$ .  
THEN  $u_x = v_y \varepsilon^{-p}$  AND SO

$$\varepsilon^{1-2p} v'' + \varepsilon^{-p} v' = 1/2, \quad y \geq 0$$

$$v(0) = 0$$

THE ONLY CONSISTENT BALANCE IS TO TAKE  $p=1$  AND SO

$$v'' + v' = \varepsilon/2 \quad \text{ON } 0 \leq y < \infty$$

$$v(0) = 0$$

WE EXPAND  $v = v_0 + \varepsilon v_1 + \dots$  AND OBTAIN

$$v_0'' + v_0' = 0 \quad y \geq 0$$

$$v_0(0) = 0$$

$$\text{THE SOLUTION IS } v_0(y) = A_0(1 - e^{-y})$$

THE INNER EXPANSION IS  $v = A_0(1 - e^{-y}) + O(\varepsilon) + \dots$

VAN-DYKE (1964) MATCHING PRINCIPLE:

NOW THE MATCHING STATEMENT IS THAT THE OUTER EXPANSION  
WHEN EXPRESSED IN INNER VARIABLES (AND LETTING  $\varepsilon \rightarrow 0$ ) MUST  
AGREE ASYMPTOTICALLY WITH THE INNER SOLUTION WHEN  
EXPRESSED IN OUTER VARIABLES (AND LETTING  $\varepsilon \rightarrow 0$ ).

IN THIS SIMPLE CASE THE MATCHING CONDITION IS

$$\frac{1}{2} = \lim_{x \rightarrow 0} u_{out} = \lim_{y \rightarrow \infty} v(y) = A_0 \rightarrow A_0 = 1/2$$

$$u_{out}(x) = (1+x)/2, \quad u_{inn}(x) = \frac{1}{2}(1 - e^{-x/\varepsilon}) \quad x/\varepsilon = 0(1)$$

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A COMPOSITE EXPANSION IS OBTAINED BY ADDING  $U_{out}(X) + U_{inn}(X)$  AND SUBTRACTING THE COMMON PART  $1/2$  USED IN MATCHING.

THIS YIELDS 
$$U_{comp} = \frac{(1+X)}{2} + \frac{1}{2}(1 - e^{-X/\varepsilon}) - \frac{1}{2} = \frac{(1+X)}{2} - \frac{1}{2}e^{-X/\varepsilon}.$$

REMARK ANOTHER WAY TO MATCH INNER AND OUTER SOLUTIONS IS FROM THE CONCEPT OF INTERMEDIATE LIMITS. THE INNER AND OUTER EXPANSIONS MUST AGREE ASYMPTOTICALLY ON ALL INTERMEDIATE SCALES.

IF WE DEFINE  $X_\Lambda = X/\Lambda(\varepsilon)$   $\varepsilon \ll \Lambda(\varepsilon) \ll 1$  AND  $X_\Lambda$  FIXED

THEN 
$$Y = X_\Lambda \Lambda(\varepsilon)/\varepsilon \rightarrow \infty \text{ AS } \varepsilon \rightarrow 0 \text{ SINCE } \Lambda(\varepsilon) \gg \varepsilon$$

$$X = X_\Lambda \Lambda(\varepsilon) \rightarrow 0 \text{ AS } \varepsilon \rightarrow 0$$

THE INTERMEDIATE MATCHING PRINCIPLE OF KAPLUN (1964) STATES THAT THE INNER AND OUTER SOLUTIONS SHOULD AGREE WHEN WRITTEN IN TERM OF  $X_\Lambda$ .

EXAMPLE CONSTRUCT THE SOLUTION TO

$$\varepsilon U'' + (X^2 - 2)U = -1 \quad U(-1) = U(1) = 0, \quad -1 \leq X \leq 1.$$

IN THE OUTER REGION WE WRITE

$$U_{out} = U_0 + \varepsilon U_1 + \dots$$

THIS YIELDS 
$$(2 - X^2)U_0 = 1$$

$$(2 - X^2)U_1 = -U_0''$$

THUS 
$$U_0 = 1/(2 - X^2) \quad \text{AND SO} \quad U_{out} = \frac{1}{2 - X^2} + O(\varepsilon)$$

NOW AS  $X \rightarrow \pm 1$  THE  $U_0 \rightarrow \pm 1$  WHICH DOES NOT SATISFY THE BOUNDARY CONDITION. WE NEED A BL NEAR  $X = \pm 1$

CONSIDER THE REGION NEAR  $x = -1$ .

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$$\text{LET } y = (x+1)/\varepsilon^p \quad \text{AND} \quad v(y) = u(-1 + \varepsilon^p y).$$

$$\text{THEN} \quad \varepsilon^{1-2p} v'' - (2 - (1 + \varepsilon^p y)^2) v = -1 \quad 0 \leq y < \infty$$

$$\text{TO BALANCE TERMS TAKE } p = 1/2. \rightarrow v'' - (2 - (1 + 2\varepsilon^{1/2} y + \dots)) v = -1$$

$$\text{WE EXPAND } v = v_0 + \varepsilon^{1/2} v_1 + \dots$$

$$\left. \begin{array}{l} v_0'' - v_0 = -1 \\ v_0(0) = 0 \end{array} \right\} \rightarrow v_0 = +1 + A e^{-y} + B e^y.$$

NOW THE MATCHING CONDITION IS THAT

$$\lim_{x \rightarrow 0} u_{\text{out}} \sim \lim_{y \rightarrow \infty} v$$

$$u_{\text{out}} = \frac{1}{2-x^2} = \frac{1}{2-(1+\varepsilon^{1/2}y)} \sim \frac{1}{1-\varepsilon^{1/2}y} \sim 1 + \varepsilon^{1/2}y + \dots$$

$$v_0 = 1 + A e^{-y} + B e^y \rightarrow 1 \quad \text{AS } y \rightarrow \infty \quad \text{WHEN } B = 0$$

TO ELIMINATE EXPONENTIAL GROWTH.

THUS  $v_0$  AND  $u_0$  MATCH.

$$\text{FINALLY MAKE } v_0(0) = 0 \rightarrow A = -1.$$

$$u_{\text{INNER}} = 1 - e^{-y} + O(\varepsilon^{1/2}), \quad y = \varepsilon^{-1/2}(1+x)$$

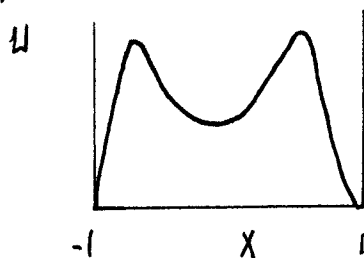
$$u_{\text{out}} = \frac{1}{2-x^2} + O(\varepsilon), \quad -1 + O(\varepsilon^{1/2}) < x < 1 - O(\varepsilon^{1/2})$$

$$u_{\text{INNER}} = 1 - e^{-\hat{y}} + O(\varepsilon^{1/2}), \quad \hat{y} = \varepsilon^{1/2}(1-x)$$

NOTICE THE SYMMETRY UNDER  $x$  TO  $-x$ .

THE COMPOSITE EXPANSION IS

$$U_{\text{COMP}} \sim \frac{1}{2-x^2} - e^{-(1+x)\varepsilon^{-1/2}} - e^{-(1-x)\varepsilon^{-1/2}} + \dots$$



EXAMPLE FIND THE LEADING ORDER INNER AND OUTER SOLUTION FOR

$$\varepsilon U'' - (1-x^2)U = -1 \quad U(\pm 1) = 0, \quad -1 \leq x \leq 1.$$

OUTER REGION: LET  $U = U_0 + \varepsilon U_1 + \dots$

$$\text{THEN } (1-x^2)U_0 = 1 \quad \rightarrow \quad U_0 = 1/(1-x^2).$$

$$+(1-x^2)U_1 = U_0''$$

$$\text{HENCE } U_{\text{out}} \sim U_0 = \frac{1}{1-x^2} + O(\varepsilon).$$

IN BOUNDARY LAYER REGION NEAR  $x=1$  WE LET

$y = \varepsilon^{-p}(x+1)$ ,  $p > 0$ . IN TERMS OF  $y$  NOTICE THAT

$$U_{\text{out}} \sim \frac{1}{1-x^2} = \frac{1}{1-(-1+\varepsilon^p y)^2} = \frac{1}{2\varepsilon^p y + \dots}$$

THIS SUGGESTS THAT THE INNER SOLUTION SHOULD HAVE

THE FORM

$$y = \varepsilon^{-p}(x+1), \quad U(x) = \varepsilon^{-p} V(\varepsilon^{-p}(x+1)).$$

THEN

$$\varepsilon U'' - (1-x^2)U = -1 \quad \text{TRANSFORM TO}$$

$$\varepsilon^{1-3P} V'' - (1 - (-1 + \varepsilon^P \gamma)^2) \varepsilon^{-P} V = -1.$$

$$\varepsilon^{1-3P} V'' - (2 \varepsilon^P \gamma + \dots) \varepsilon^{-P} V = -1$$

TO BALANCE TERMS TAKE  $P = 1/3$  AND

EXPAND  $V = V_0 + \dots$  THIS YIELDS

$$(*) \quad V_0'' - 2\gamma V_0 = -1 \quad 0 \leq \gamma < \infty$$

$$V_0(0) = 0, \quad V_0(\gamma) \sim 1/2\gamma \quad \text{AS } \gamma \rightarrow \infty \quad \text{MATCHING CONDITION}$$

THE SOLUTION CAN BE WRITTEN AS

$$V_0(\gamma) = c_1 A_i(\gamma^{2/3}) + c_2 B_i(\gamma^{2/3}) + V_p(\gamma)$$

WHERE  $A_i(z), B_i(z)$  ARE THE TWO FUNDAMENTAL SOLUTIONS OF

$$\text{AIRY'S EQUATION} \quad W''(z) - z W(z) = 0$$

RECALL THAT  $A_i(z)$  DECAYS EXPONENTIALLY AS  $z \rightarrow +\infty$

$B_i(z)$  GROWS EXPONENTIALLY AS  $z \rightarrow \infty$

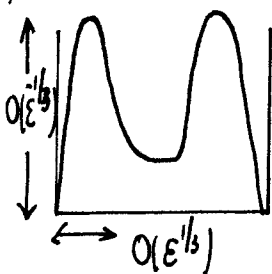
ALSO  $V_p(\gamma)$  IS THE PARTICULAR SOLUTION OF  $(*)$  WHICH

CLEARLY SATISFIES  $V_p(\gamma) \sim 1/2\gamma + d/\gamma^4 + \dots$  AS  $\gamma \rightarrow \infty$ .

HENCE WE SET  $c_2 = 0$  TO KILL EXPONENTIAL GROWTH

AND  $c_1 A_i(0) = -V_p(0) \rightarrow c_1 = -V_p(0)/A_i(0)$ .

$$V_0(\gamma) = V_p(\gamma) - V_p(0) A_i(\gamma^{2/3}) / A_i(0)$$



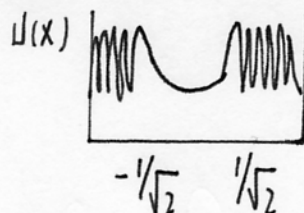
REMARK

(i) IF  $\epsilon u'' - \left(\frac{1}{2} - x^2\right) u = -1$   $-1 \leq x \leq 1$ ,  $u(\pm 1) = 0$  (9)

THEN HIGHLY OSCILLATORY ON  $-1 < x < -\sqrt{\frac{1}{2}}$  AND  $\frac{1}{\sqrt{2}} < x < 1$ .

B.L. THEORY IS NOT APPROPRIATE. NEED WKB TURNING

POINT THEORY



### GENERAL CASE: LINEAR TWO-POINT BVP

$$\epsilon u'' + a(x)u' + b(x)u = 0 \quad 0 \leq x \leq 1 \quad \epsilon \rightarrow 0^+$$

$$u(0) = A, \quad u(1) = B$$

WE WILL ASSUME THAT  $a(x) > 0$  ON  $0 \leq x \leq 1$  OR  $a(x) < 0$  ON  $0 \leq x \leq 1$ . WE WILL EXCLUDE THE "TURNING POINT" CASE WHERE  $a(x_0) = 0$   $x_0 \in (0, 1)$  FOR THE MOMENT.

WE WILL SHOW

(i) IF  $a(x) > 0$  ON  $[0, 1] \rightarrow$  B.L. NEAR  $x = 0$

(ii) IF  $a(x) < 0$  ON  $[0, 1] \rightarrow$  B.L. NEAR  $x = 1$ .

IN THE OUTER REGION WE EXPAND  $u = u_0 + \epsilon u_1 + \dots$  TO GET

$$a(x)u_0' + b(x)u_0 = 0, \quad a(x)u_i' + b(x)u_i = -u_{i-1}'', \quad i \geq 1$$

IN CASE (i) WE SET  $u_0(1) = B$  AND IN CASE (ii) WE TAKE  $u_0(0) = A$ . THEN UPON INTEGRATING

$$(i) \quad u_0(x) = B \exp \left( \int_x^1 \frac{b(s)}{a(s)} ds \right)$$

$$(ii) \quad u_0(x) = A \exp \left( - \int_0^x \frac{b(s)}{a(s)} ds \right)$$

CASE (i)

$Q(X) > 0$ . WE PUT A BL NEAR  $X=0$ .

(10)

WE LET  $\gamma = X/\epsilon^p$  AND  $V(\gamma) = U(\epsilon^p X)$ , WITH  $p > 0$

THEN  $\epsilon^{1-2p} V'' + \epsilon^{-p} a(\epsilon^p \gamma) V' + b(\epsilon^p \gamma) V = 0$ .

$a(\epsilon^p \gamma) \sim a(0) + \dots$   $b(\epsilon^p \gamma) \sim b(0) + \dots$  WE THEN BALANCE  $\epsilon^{1-2p} = \epsilon^{-p}$

TO OBTAIN  $p=1$ . WE THEN EXPAND  $V = V_0 + \epsilon V_1 + \dots$  TO OBTAIN

$$V_0'' + a(0) V_0' = 0, \quad V_0(0) = A$$

$$V_0(\infty) = U_0(0) \quad \text{MATCHING CONDITION}$$

SINCE  $a(0) > 0 \rightarrow V_0'$  DECAYS EXPONENTIALLY.

THE SOLUTION IS  $V_0(\gamma) = (A - C_0) + C_0 e^{-a(0)\gamma}$

THE MATCHING CONDITION GIVES  $A - C_0 = U_0(0)$ , SO THAT

$$C_0 = A - U_0(0) = A - B \exp\left(\int_0^1 \frac{b(s)}{a(s)} ds\right)$$

IN SUMMARY: LEADING ORDER INNER, OUTER, AND COMPOSITE

$$U_{\text{INN}} = B \exp\left(\int_0^1 \frac{b(s)}{a(s)} ds\right) + \left(A - B \exp\left(\int_0^1 \frac{b(s)}{a(s)} ds\right)\right) e^{-a(0)\gamma}$$

$$U_{\text{OUT}} = B \exp\left(\int_x^1 \frac{b(s)}{a(s)} ds\right)$$

$$U_{\text{COMP}} = B \exp\left(\int_x^1 \frac{b(s)}{a(s)} ds\right) + \left(A - B \exp\left(\int_0^1 \frac{b(s)}{a(s)} ds\right)\right) e^{-a(0)X/\epsilon}$$

REMARK (i) IF WE INCORRECTLY ASSUMED A BOUNDARY LAYER AT  $X=1$  INSTEAD OF AT  $X=0$  WE WOULD HAVE

$$V_0'' - a(1) V_0' = 0, \quad 0 < \gamma < \infty \quad \gamma = (1-X)/\epsilon$$

THIS GROWS EXPONENTIALLY AS  $\gamma \rightarrow +\infty$ !!!

CASE 2

NOW ASSUME THAT  $q(x) < 0$  ON  $[0, 1]$ .

(11)

WE INSERT A BL NEAR  $x = 1$ .LET  $y = \varepsilon^{-1}(1-x)$ ,  $W(y) = U[1 - \varepsilon y]$ . THIS GIVES

$$\frac{1}{\varepsilon} W'' - \frac{1}{\varepsilon} a(1 - \varepsilon y) W' + b(1 - \varepsilon y) W = 0 \quad 0 < y < \infty.$$

WE EXPAND  $W = W_0 + \varepsilon W_1 + \dots$ 

$$W_0'' - a(1) W_0' = 0, \quad 0 < y < \infty$$

$$W_0(0) = B, \quad W_0(\infty) = U_0(1) \text{ MATCHING CONDITION.}$$

SINCE  $a(1) < 0$  WE HAVE EXPONENTIAL DECAY AS  $y \rightarrow +\infty$ .

$$\text{THE SOLUTION IS } W_0 = (B - C_1) + C_1 e^{a(1)y}$$

$$\text{MATCHING CONDITION IS } W_0(\infty) = U_0(1) = A \exp\left(-\int_0^1 \frac{b(s)}{a(s)} ds\right) = B - C_1$$

$$\text{THUS } U_{\text{INN}} = W_0(y) = A \exp\left(-\int_0^1 \frac{b(s)}{a(s)} ds\right) + \left[B - A \exp\left(-\int_0^1 \frac{b(s)}{a(s)} ds\right)\right] e^{a(1)(1-x)/\varepsilon}$$

$$U_{\text{OUT}} = A \exp\left(-\int_0^x \frac{b(s)}{a(s)} ds\right)$$

$$U_{\text{COMP}} = A \exp\left(-\int_0^x \frac{b(s)}{a(s)} ds\right) + \left(B - A \exp\left(-\int_0^1 \frac{b(s)}{a(s)} ds\right)\right) e^{a(1)(1-x)/\varepsilon}$$

EXAMPLE

$$\varepsilon U'' + (1 + 2x) U' + U = x$$

$$U(0) = 0, \quad U(1) = 1$$

SINCE  $q(x) = 1 + 2x > 0$  ON  $[0, 1]$  WE WILL HAVE A BLAT  $x = 0$ . IN THE OUTER REGION  $0 < x \leq 1$  WE EXPAND

$$U = U_0 + \varepsilon U_1 + \dots$$

$$(1+2x)U_0' + U_0 = X, \quad U_0(1) = 1$$

$$(1+2x)U_1' + U_1 = -U_0'', \quad U_1(1) = 0$$

SOLVING FOR  $U_0$  WE GET

$$U_0 = \frac{C_0}{(1+2x)^{1/3}} + \frac{1}{3}(x-1) \quad U_0(1) = 1 \rightarrow C_0 = \sqrt{3}$$

THEN,  $(1+2x)U_1' + U_1 = -3C_0(1+2x)^{-5/2}, \quad U_1(1) = 0$  GIVES

$$U_1 = \frac{1}{3} \frac{(2-x-x^2)C_0}{(1+2x)^{5/2}} \quad C_0 = \sqrt{3}.$$

THU TO TWO TERMS WE HAVE

$$U_{out}(x) = \frac{\sqrt{3}}{(1+2x)^{1/2}} + \frac{1}{3}(x-1) + \frac{\varepsilon}{\sqrt{3}} \frac{(2-x-x^2)}{(1+2x)^{5/2}} + \dots$$

NOW PUT A BL NEAR  $x = 0$ . LET  $y = x/\varepsilon, \quad V(y) = U(\varepsilon y)$ .

THEN  $V'' + (1+2\varepsilon y)V' + \varepsilon V = \varepsilon^2 y$

$$V(0) = 0.$$

WE EXPAND  $V = V_0 + \varepsilon V_1 + \dots$

$$(i) \quad V_0'' + V_0' = 0, \quad V_0(0) = 0$$

$$(ii) \quad V_1'' + V_1' = -2yV_0' - V_0, \quad V_1(0) = 0.$$

NOW WE WRITE  $U_{out}(x)$  IN TERMS OF  $y = x/\varepsilon$  OR

$x = \varepsilon y$  AND LET  $\varepsilon \rightarrow 0$ .

$$U_{out} \sim \frac{\sqrt{3}}{(1+2\varepsilon y)^{1/2}} + \frac{1}{3}(-1 + \varepsilon y) + \frac{\varepsilon}{\sqrt{3}} \frac{(2 - \varepsilon y - \varepsilon^2 y^2)}{(1+2\varepsilon y)^{5/2}} + \dots$$

THIS YIELD THAT

(13)

$$U_{out} \sim \sqrt{3} \left( 1 - \varepsilon \gamma + \dots \right) - \frac{1}{3} + \varepsilon \left( \frac{\gamma}{3} + \frac{2}{\sqrt{3}} \right) + O(\varepsilon^2)$$

$$U_{out} \sim \sqrt{3} - \frac{1}{3} + \varepsilon \left( \frac{2}{\sqrt{3}} + \gamma \left( \frac{1}{3} - \sqrt{3} \right) \right) + O(\varepsilon^2)$$

COMPARING WITH  $U_{INN} = V_0 + \varepsilon V_1 + \dots$

WE REQUIRE THAT

$$(ia) \quad V_0 \rightarrow \sqrt{3} - \frac{1}{3}, \quad \gamma \rightarrow +\infty$$

$$(iia) \quad V_1 \rightarrow \frac{2}{\sqrt{3}} + \gamma \left( \frac{1}{3} - \sqrt{3} \right), \quad \gamma \rightarrow +\infty$$

WE NOW SOLVE (i) AND (ii) WITH (ia) AND (iia).

$$\text{WE OBTAIN } V_0(\gamma) = \left( \sqrt{3} - \frac{1}{3} \right) (1 - e^{-\gamma})$$

$$(e^{\gamma} V_1')' = -2\gamma e^{\gamma} V_0' - V_0 e^{\gamma} = -2\gamma \left( \sqrt{3} - \frac{1}{3} \right) - \left( \sqrt{3} - \frac{1}{3} \right) (e^{\gamma} - 1)$$

$$\text{SOLVING THIS WITH } V_1(0) = 0 \text{ AND } V_1 \rightarrow \frac{2}{\sqrt{3}} + \gamma \left( \frac{1}{3} - \sqrt{3} \right)$$

WE OBTAIN

$$V_1 = \left( \sqrt{3} - \frac{1}{3} \right) e^{-\gamma} (\gamma^2 + \gamma + 1) - \left( \sqrt{3} - \frac{1}{3} \right) (\gamma + 1) + \left( \sqrt{3} - \frac{1}{3} + \frac{2}{\sqrt{3}} \right) (1 - e^{-\gamma})$$

EXAMPLE

FIND A TWO-TERM EXPANSION FOR THE SOLUTION TO

(14)

$$\varepsilon u'' + u' - xu = 0, \quad 0 \leq x \leq 1; \quad u(0) = 0, \quad u(1) = e^{1/2}.$$

SINCE  $q(x) = 1 > 0$  ON  $[0, 1]$  WE GET A BL AT  $x = 0$ .IN THE OUTER REGION  $0 < x \leq 1$  WE EXPAND  $u = u_0 + \varepsilon u_1 + \dots$ 

THIS YIELDS:

$$u_0' - xu_0 = 0; \quad u_0(1) = e^{1/2}$$

$$u_1' - xu_1 = -u_0'', \quad u_1(0) = 0$$

WE OBTAIN  $u_0 = e^{x^2/2}$  AND  $u_1 = e^{x^2/2} \left( \frac{4}{3} - x - \frac{x^3}{3} \right)$ .

THUS  $u_{out} = e^{x^2/2} \left( 1 + \varepsilon \left( \frac{4}{3} - x - \frac{x^3}{3} \right) + \dots \right)$

IN THE INNER REGION WE LET  $y = x/\varepsilon$ ,  $v(y) = u(\varepsilon y)$ .

THEN  $\frac{1}{\varepsilon} v'' + \frac{1}{\varepsilon} v' - \varepsilon y v = 0$ .

HENCE  $v'' + v' - \varepsilon^2 y v = 0, \quad 0 \leq y < \infty$

$$v(0) = 0.$$

NOW TAKE THE OUTER SOLUTION AND LET  $x = \varepsilon y$  AND

EXPAND AS  $\varepsilon \rightarrow 0$ :  $u_{out} \sim e^{\varepsilon^2 y^2/2} \left( 1 + \frac{4\varepsilon}{3} + O(\varepsilon^2) \right) \sim 1 + \frac{4\varepsilon}{3} + O(\varepsilon^2)$

MATCHING WITH  $v$  WE CONCLUDE

$$v = v_0 + \varepsilon v_1 + \dots \quad v_0(\infty) = 1, \quad v_1(\infty) = \frac{4}{3}.$$

THEN  $v_0'' + v_0' = 0, \quad v_1'' + v_1' = 0 \rightarrow v_0 = 1 - e^{-y}, \quad v_1 = \frac{4}{3}(1 - e^{-y})$

OUTER  $u_{out} = e^{x^2/2} \left( 1 + \varepsilon \left( \frac{4}{3} - x - \frac{x^3}{3} \right) \right), \quad u_{inn} = (1 - e^{-y}) \left( 1 + \frac{4\varepsilon}{3} + \dots \right)$

THE COMPOSITE IS

(15)

$$U_c = e^{x^2/2} \left( 1 + \varepsilon \left( \frac{4}{3} - x - \frac{x^3}{3} \right) \right) - \left( 1 + \frac{4\varepsilon}{3} \right) e^{-x/\varepsilon}.$$

EXAMPLE (A NONLINEAR PROBLEM)

$$\varepsilon U'' + U' + U^2 = 0, \quad 0 \leq x \leq 1, \quad 0 < \varepsilon \ll 1$$

$$U'(0) = A, \quad U(1) = \frac{1}{2} + 2\varepsilon$$

SINCE  $Q(x) = 1 > 0$  WE EXPECT BL AT  $x = 0$ .

WE THEN EXPAND  $U = U_0 + \varepsilon U_1 + \dots$   $0 < x \leq 1$  IN OUTER REGION.

THIS GIVES

$$U_0' + U_0^2 = 0, \quad U_0(1) = 1/2$$

$$U_1' + 2U_0 U_1 = -U_0'', \quad U_1(1) = 2$$

$$\text{WE SOLVE TO OBTAIN } U_0 = \frac{1}{x+1}, \quad U_1 = \frac{2 \log(x+1)}{(x+1)^2} + \frac{C_1}{(x+1)^2}.$$

WE WANT  $U_1(1) = 2$  SO THAT  $C = 8 - 2 \log 2$ .

THU TO TWO TERMS WE HAVE

$$U_{out}(x) = \frac{1}{x+1} + \varepsilon \left[ \frac{2 \log(x+1)}{(x+1)^2} + \frac{(8 - 2 \log 2)}{(x+1)^2} \right]$$

IN THE INNER REGION WE LET  $y = x/\varepsilon, \quad v(y) = U(\varepsilon y).$

$$\text{THEN WE OBTAIN } \frac{1}{\varepsilon} v'' + \frac{1}{\varepsilon} v' + v^2 = 0, \quad v'(0) = \varepsilon A$$

WE THEN EXPAND  $V = V_0 + \varepsilon V_1 + \dots$

$$V_0'' + V_0' = 0, \quad 0 \leq \gamma < \infty, \quad V_0'(0) = 0$$

$$V_1'' + V_1' = -V_0^2, \quad 0 \leq \gamma < \infty, \quad V_1'(0) = A$$

NOW LET  $X = \varepsilon \gamma$  IN OUTER EXPANSION AND LET  $\varepsilon \rightarrow 0$ .

$$U_{\text{out}} \sim \frac{1}{1 + \varepsilon \gamma} + \varepsilon \left[ \frac{\log(1 + \varepsilon \gamma)}{(1 + \varepsilon \gamma)^2} + \frac{(8 - 2 \log 2)}{(1 + \varepsilon \gamma)^2} \right] \sim 1 - \varepsilon \gamma + \varepsilon (8 - 2 \log 2) + O(\varepsilon^2)$$

WE OBTAIN  $V_0(\infty) = 1$

$$V_1(\gamma) \sim -\gamma + 8 - 2 \log 2 \quad \text{AS } \gamma \rightarrow \infty.$$

THUS  $V_0 = c_0 + d_0 e^{-\gamma} \quad V_0'(0) = 0, \quad V_0(\infty) = 1 \rightarrow V_0(\gamma) \equiv 1$

THEN  $V_1'' + V_1' = -1$  HAS SOLUTION  $V_1 = -\gamma + C_1 + C_2 e^{-\gamma}$ .

$$V_1'(0) = A \rightarrow -1 - C_2 = A$$

$$V_1 \sim -\gamma + 8 - 2 \log 2 \quad \text{AS } \gamma \rightarrow \infty \quad \text{IMPLIES} \quad C_1 = 8 - 2 \log 2.$$

THUS  $V_1(\gamma) = -\gamma + 8 - 2 \log 2 - (1 + A) e^{-\gamma}$

$$U_{\text{INN}} = 1 + \varepsilon \left( -\gamma + 8 - 2 \log 2 - (1 + A) e^{-\gamma} \right)$$

$$U_{\text{out}} = \frac{1}{X+1} + \varepsilon \left( \frac{2 \log(X+1)}{(X+1)^2} + \frac{8 - 2 \log 2}{(X+1)^2} \right)$$

$$U_{\text{COMP}} = U_{\text{out}}(X) - \varepsilon (1 + A) e^{-X/\varepsilon} + O(\varepsilon^2).$$

FIND A 1-TERM INNER, OUTER, AND COMPOSITE SOLUTION FOR

$$\varepsilon U'' - U' + e^U = 0 \quad 0 \leq x \leq 1$$

$$U(0) = A, \quad U(1) = 0 \quad A < 0.$$

SINCE  $a(x) = -1 < 0$  WE EXPECT A BL NEAR  $x=1$ .

IN THE OUTER REGION FOR  $0 \leq x < 1$  WE EXPAND

$$U = U_0 + \varepsilon U_1 + \dots$$

THIS YIELDS  $U_0' = e^{U_0}, \quad U_0(0) = A$

$$U_0(x) = -\ln[e^{-A} - x] \quad \text{SINCE } A < 0 \text{ THE SINGULARITY AT } x = e^{-A} \text{ OCCURS FOR } x > 1.$$

IN THE INNER REGION NEAR  $x=1$  WE PUT  $y = \varepsilon^{-1}(1-x)$ ,  
AND  $V(y) = U(1 - \varepsilon y)$ . THEN

$$\frac{1}{\varepsilon} V'' + \frac{1}{\varepsilon} V' + e^V = 0 \quad y \geq 0$$

$$V(0) = 0, \quad \text{AND WE WRITE } V = V_0 + \varepsilon V_1 + \dots$$

THE MATCHING CONDITION IS  $V_0(\infty) = U_0(1) = -\ln(e^{-A} - 1)$ .

THEN  $V_0'' + V_0' = 0$  SO THAT

$$V_0 = -\ln(e^{-A} - 1)(1 - e^{-y})$$

leading order outer:  $U_{out} = -\ln(e^{-A} - x)$

" " INNER:  $U_{inn} = -\ln(e^{-A} - 1)(1 - e^{-(1-x)/\varepsilon})$

" " COMPOSITE:  $U_{comp} = -\ln(e^{-A} - x) + \ln(e^{-A} - x)e^{-(1-x)/\varepsilon}$

REMARK (i) IF  $A > 0$  THEN  $U_0$  HAS A SINGULARITY INSIDE THE DOMAIN  $0 \leq x \leq 1$ . IN THIS CASE THIS SUGGESTS NO SOLUTION TO THE BVP.

NOW IN THE BL REGION NEAR  $X=0$  WE LET  $\gamma = X/\varepsilon^p$ ,  $V(\gamma) = U(\varepsilon^p X)$ . (18)

THIS YIELDS

$$\varepsilon^{1-2p} V'' + \frac{\varepsilon^p \gamma V'}{\varepsilon^p} - \varepsilon^p \gamma V = 0, \quad 0 \leq \gamma < \infty$$

$$V(0) = 0.$$

TO BALANCE TERMS WE TAKE  $\varepsilon^{1-2p} = \varepsilon^0 \rightarrow p = 1/2$ .

THIS GIVES 
$$\left\{ \begin{array}{l} V'' + \gamma V' - \varepsilon^{1/2} \gamma V = 0 \\ V(0) = 0 \end{array} \right. \quad 0 \leq \gamma < \infty$$

NOW WE TAKE THE OUTER SOLUTION AND LET  $X = \varepsilon^{1/2} \gamma$ .

THIS LEADS TO 
$$U_{out} = e^{\varepsilon^{1/2} \gamma} - \varepsilon e^{\varepsilon^{1/2} \gamma} \left[ \frac{1}{2} \log \varepsilon + \log \gamma \right]$$

$$U_{out} = 1 + \varepsilon^{1/2} \left( \gamma \right) - \frac{1}{2} \varepsilon \log \varepsilon - \varepsilon \left( \log \gamma + \gamma^2/2 \right) + \dots$$

THIS SUGGESTS THE INNER EXPANSION

$$V = V_0 + \varepsilon^{1/2} V_1 + (-\varepsilon \log \varepsilon) V_2 + \dots$$

WE OBTAIN

$$\left\{ \begin{array}{l} V_0'' + \gamma V_0' = 0 \\ V_0(0) = 0, \quad V_0(\infty) = 1 \end{array} \right.$$

$$\left\{ \begin{array}{l} V_1'' + \gamma V_1' = \gamma V_0 \\ V_1(0) = 0, \quad V_1 \sim \gamma \text{ AS } \gamma \rightarrow \infty \end{array} \right.$$

$$\left\{ \begin{array}{l} V_2'' + \gamma V_2' = 0 \\ V_2(0) = 0, \quad V_2(\infty) = + \frac{1}{2} \end{array} \right.$$

REMARK (i) SO FAR WE HAVE ONLY CONSIDERED SECOND ORDER BVP (19) OF THE FORM  $u'' + a(x)u'$  WHERE EITHER  $a(x) > 0$  OR  $a(x) < 0$  ON  $0 \leq x \leq 1$ .

(ii) WHAT HAPPENS IF  $a(x)$  VANISHES AT AN ENDPOINT?  
(BL WIDTH WILL NOT BE  $O(\varepsilon)$ ).

(iii) WHAT HAPPENS IF  $a(x_0) = 0$  FOR SOME UNIQUE  $x_0$  IN  $(0, 1)$ ?  
THESE POINTS ARE CALLED TURNING POINTS.

EXAMPLE (OF (iii) ABOVE).

CONSIDER  $\varepsilon u'' + xu' - xu = 0$ ,  $0 \leq x \leq 1$ ,  $\varepsilon \rightarrow 0$ .

$$u(0) = 0, \quad u(1) = e$$

IT IS STILL CLEAR THAT SINCE  $a(x) = x > 0$  NEAR  $x = 1$  WE CANNOT HAVE A BL NEAR  $x = 1$ .

IN THE OUTER REGION WE EXPAND  $u = u_0 + \varepsilon u_1 + \dots$   
FOR  $0 < x \leq 1$ . THIS YIELDS THAT

$$\left. \begin{aligned} u_0' - u_0 &= 0, \quad 0 < x \leq 1 \\ u_0(1) &= e \end{aligned} \right\} \rightarrow u_0 = e^x$$

$$\left. \begin{aligned} u_1' - u_1 &= -\frac{u_0''}{x} = -\frac{e^x}{x} \\ u_1(1) &= 0 \end{aligned} \right\} \rightarrow u_1 = -e^x \log x$$

THE TWO-TERM OUTER SOLUTION IS

$$u_{\text{out}} = e^x - \varepsilon e^x \log x + \dots$$

$$V_0' e^{\gamma^2/2} = A_0$$

$$V_0(\gamma) = A_0 \int_0^\gamma e^{-\lambda^2/2} d\lambda.$$

BUT  $V_0(\infty) = 1 \rightarrow A_0 \int_0^\infty e^{-\lambda^2/2} d\lambda = 1 \rightarrow A_0 = \left(\frac{2}{\pi}\right)^{1/2}.$

THIS YIELDS  $V_0(\gamma) = \left(\frac{2}{\pi}\right)^{1/2} \int_0^\gamma e^{-\lambda^2/2} d\lambda.$

THEN  $(V_1' e^{\gamma^2/2})' = \gamma e^{\gamma^2/2} V_0$

$$V_1' e^{\gamma^2/2} = \int_0^\gamma \lambda e^{\lambda^2/2} V_0(\lambda) d\lambda$$

$$V_1 = \int_0^\gamma e^{-t^2/2} \left( \int_0^t \lambda e^{\lambda^2/2} V_0(\lambda) d\lambda \right) dt.$$

INDEED  $V_1 \sim \gamma$  AS  $\gamma \rightarrow \infty.$

FINALLY  $V_2 = \frac{1}{2} \left(\frac{2}{\pi}\right)^{1/2} \int_0^\gamma e^{-\lambda^2/2} d\lambda.$

EXAMPLE (INTERNAL LAYER)

WHEN  $q(x)$  CHANGES SIGN INSIDE THE DOMAIN WE COULD HAVE AN INTERNAL LAYER WHERE  $U$  CHANGES RAPIDLY IN A NARROW REGION NEAR WHERE  $q(x)$  CROSSES THROUGH ZERO.

CONSIDER  $\epsilon U'' + (x - 1/2) U' - (x - 1/2) U = 0 \quad 0 \leq x \leq 1$

$$U(0) = \alpha, \quad U(1) = \beta \quad q(x) = x - 1/2.$$

SINCE  $q(0) < 0$  AND  $q(1) > 0$  THERE ARE NO BL NEAR  $x=0$  OR  $x=1$ .  
IN THE OUTER REGION WE EXPAND

$$U = U_0 + \epsilon U_1 + \dots$$

$$(x - 1/2) U_0' - (x - 1/2) U_0 = 0$$

WHERE  $U_0(0) = \alpha$  AND  $U_0(1) = \beta$  BUT  $U_0$  DISCONTINUOUS ACROSS  $x = 1/2$ .

WE WRITE

$$U_0(x) = \begin{cases} c_L e^x, & 0 \leq x < 1/2 \\ c_R e^x, & 1/2 < x \leq 1 \end{cases}$$

SATISFYING THE B.C. WE GET  $U_0 = \begin{cases} \alpha e^x, & 0 \leq x < 1/2 \\ \beta e^{x-1}, & 1/2 < x \leq 1. \end{cases}$

NOW WE INSERT AN INTERNAL LAYER NEAR  $x = 1/2$ . WE LET

$$y = \epsilon^{-p}(x - 1/2), \quad W(y) = U\left[\frac{1}{2} + \epsilon^p y\right].$$

THIS YIELDS,  $\epsilon^{1-2p} W'' + y W' - \epsilon^p y W = 0$ .

WE BALANCE WITH  $p = 1/2$  AND EXPAND  $W = W_0 + \epsilon^{1/2} W_1 + \dots$

$$\left. \begin{aligned} W_0'' + y W_0' &= 0, & -\infty < y < \infty \\ W_0(\infty) &= \beta e^{-1/2} & W_0(-\infty) &= \alpha e^{1/2} \end{aligned} \right\}$$

(FROM MATCHING CONDITION)

WE SOLVE TO OBTAIN  $w_0(y) = C_2 + C_1 \int_0^y e^{-s^2/2} ds$ .

(22)

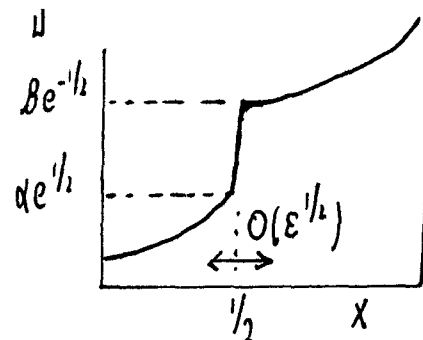
THEN  $C_2 + C_1 (\pi/2)^{1/2} = B e^{-1/2}$   $\int_0^\infty e^{-s^2/2} ds = (\pi/2)^{1/2}$

$C_2 - C_1 (\pi/2)^{1/2} = \alpha e^{1/2}$

THIS YIELDS  $C_2 = \frac{1}{2} (B e^{-1/2} + \alpha e^{1/2})$ ,  $C_1 = \frac{1}{(2\pi)^{1/2}} (B e^{-1/2} - \alpha e^{1/2})$

WHICH COMPLETELY DETERMINES  $w_0(y)$ .

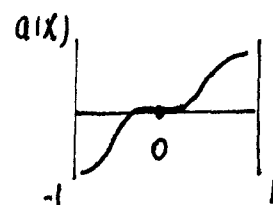
WHEN  $\alpha e^{1/2} < B e^{-1/2}$  WE GET



### EXAMPLE (INTERNAL LAYER)

CONSIDER  $\epsilon u'' + a(x)u' = 0$   $-1 \leq x \leq 1$

$u(-1) = \alpha$ ,  $u(1) = B$ .



SUPPOSE THAT  $a(0) = 0$ ,  $a(x) > 0$  FOR  $x > 0$  AND  $a(x) < 0$  FOR  $x < 0$  AND ASSUME THAT  $a(x) \sim \chi x^{2m+1}$  AS  $x \rightarrow 0$  WHERE  $m > 0$  IS AN INTEGER AND  $\chi > 0$ .

IN THE OUTER REGION  $u = u_0 + \epsilon u_1 + \dots$  SO THAT  $a(x)u_0' = 0$ . THERE WILL BE NO BL NEAR  $x = \pm 1$  SO THAT  $u_0$  MUST SATISFY THE BOUNDARY CONDITIONS. THE OUTER SOLUTION IS

$u_0 = \begin{cases} \alpha & -1 \leq x < 0 \\ B & 0 < x \leq 1 \end{cases}$

NOW CONSTRUCT AN INTERNAL LAYER NEAR  $x = 0$  WITH

$y = x/\epsilon^p$  AND  $w(y) = u(\epsilon^p y)$ .

THIS YIELDS THAT  $\varepsilon^{1-2p} W'' + \varepsilon^{-p} a(\varepsilon^p y) W' = 0$ .

NOW  $a(\varepsilon^p y) = \gamma \varepsilon^{(2m+1)p} y^{2m+1}$ . WE BALANCE  $\varepsilon^{1-2p} = \varepsilon^{-p+(2m+1)p}$ .

THIS YIELDS  $1-2p = 2mp$  SO  $p = 1/(2m+2)$

WE THEN WRITE  $W = W_0 + \dots$

$$W_0'' + \gamma y^{2m+1} W_0' = 0 \quad -\infty < y < \infty$$

$$W_0(\infty) = B, \quad W_0(-\infty) = \alpha.$$

THE INTERNAL LAYER WIDTH IS  $O(\varepsilon^{1/(2m+2)})$  WHICH WIDENS AS  $m$  INCREASES. WE SOLVE TO OBTAIN

$$W_0(y) = C_1 + C_2 \int_0^y e^{-\gamma s^{2m+2}/(2m+2)} ds$$

$$\text{NOW} \quad \int_0^\infty e^{-\gamma s^{2m+2}/(2m+2)} ds = \left(\frac{2m+2}{\gamma}\right)^{1/2m+2} \frac{1}{2m+2} \int_0^\infty r^{1/(2m+2)-1} e^{-r} dr$$

RECALLING  $\int_0^\infty t^{z-1} e^{-t} dt = \Gamma(z)$ ,  $\Gamma$  = gamma FUNCTION.

$$\text{THEN} \quad \int_0^\infty e^{-\gamma s^{2m+2}/(2m+2)} ds = \left(\frac{2m+2}{\gamma}\right)^{1/2m+2} \frac{1}{2m+2} \Gamma(1/2m+2) = I.$$

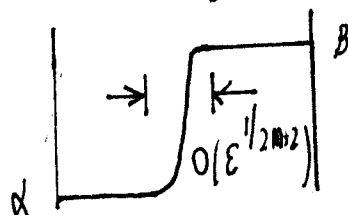
$$\text{THEN} \quad C_1 + C_2 I = B \quad (W_0(\infty) = B)$$

$$C_1 - C_2 I = \alpha \quad (W_0(-\infty) = \alpha)$$

$$\text{THIS YIELDS} \quad C_1 = (B + \alpha)/2, \quad C_2 = (B - \alpha)/2I$$

$$W_0 = \frac{B + \alpha}{2} + \frac{(B - \alpha)}{2I} \int_0^y e^{-\gamma s^{2m+2}/(2m+2)} ds.$$

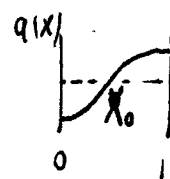
$$\text{WITH} \quad y = X/\varepsilon^{1/(2m+2)}$$



$$\varepsilon u'' + a(x)u' + b(x)u = 0 \quad 0 \leq x \leq 1$$

SUPPOSE THAT FOR SOME UNIQUE  $x_0 \in (0, 1)$  THAT

(i)  $a(x_0) = 0$ ,  $a(x) < 0$  FOR  $x < x_0$ ,  $a(x) > 0$  FOR  $x > x_0$



WITH  $a'(x_0) > 0$ . THEN WE HAVE OUTER SOLUTION ON

EITHER SIDE OF  $x_0$  AND AN INTERNAL LAYER OF WIDTH

$O(\varepsilon^{1/2})$  NEAR  $x_0$ . THERE ARE NO BOUNDARY LAYERS NEAR  $x=0, 1$ .

(ii) IF  $a(x_0) = 0$ ,  $a(x) > 0$  FOR  $x < x_0$ ,  $a(x) < 0$  FOR  $x > x_0$

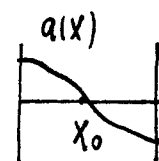
WITH  $a'(x_0) < 0$ . THEN  $u=0$  IS THE OUTER SOLUTION

AND THERE IS NO INTERNAL LAYER NEAR  $x=x_0$ .

INSTEAD THERE ARE BOUNDARY LAYERS NEAR

$x=0, 1$ . THERE ARE SOME EXCEPTIONAL CASES

TO THE STATEMENT THAT  $u=0$  IS OUTER SOLUTION



EXAMPLE (A CORNER LAYER: BOUNDARY LAYER IN THE DERIVATIVE)

CONSIDER  $\varepsilon u'' + xu' - u = 0 \quad -1 \leq x \leq 1 \quad \varepsilon \rightarrow 0$

$$u(-1) = 1, \quad u(1) = 2.$$

SINCE  $q(x) = x$  WITH  $q(-1) < 0$  AND  $q(1) > 0$  THERE ARE NO BOUNDARY LAYERS NEAR  $x=0$  AND  $x=1$ .

THE OUTER SOLUTION IS EXPANDED AS

$$u = u_0 + \varepsilon u_1 + \dots$$

SO THAT  $xu_0' - u_0 = 0$  IN  $[-1, 0)$  AND  $(0, 1]$

WE OBTAIN  $u_0 = cx$  IS THE SOLUTION.

THEREFORE

$$u_0(x) = \begin{cases} 2x, & 0 < x \leq 1 \\ -x, & -1 \leq x < 0 \end{cases}$$

(25)

NOTICE THAT  $u_0$  IS CONTINUOUS BUT  $u_0'$  IS NOT.

WE MUST CONSTRUCT AN INTERNAL LAYER TO SMOOTHEN  $u_0'$ .  
SUCH A LAYER IS CALLED A "CORNER" LAYER

WE LET  $y = x/\varepsilon^p$ ,  $w(y) = u(\varepsilon^p y)$ .

THUS  $\varepsilon^{1-2p} w'' + y w' - w = 0$  CHOOSE  $p = 1/2$ .

THIS GIVES  $w'' + y w' - w = 0$ ,  $-\infty < y < \infty$ .

NOW WRITING THE OUTER EXPANSION  $u_0$  IN TERMS OF  $y$  TO  
OBTAIN

$$u_0 \sim \begin{cases} 2\varepsilon^{1/2} y, & y > 0 \\ -\varepsilon^{1/2} y, & y < 0 \end{cases}$$

THUS WE SET  $w = \varepsilon^{1/2} w_0 + \dots$  WHICH YIELDS

$$w_0'' + y w_0' - w_0 = 0 \quad -\infty < y < \infty$$

$$w_0 \sim 2y \quad y \rightarrow +\infty; \quad w_0 \sim -y \quad \text{as } y \rightarrow -\infty.$$

ONE SOLUTION IS  $w_0 = y$ . TO FIND THE SECOND SOLUTION WE  
CAN USE REDUCTION OF ORDER BY SETTING  $w_0 = yV$ .

THEN  $w_0' = V + yV'$ ,  $w_0'' = yV'' + 2V'$ . THIS GIVES,

$$yV'' + 2V' + yV + y^2V' - yV = yV'' + V'(2 + y^2) = 0$$

$$V'' + \left(\frac{2}{y} + y\right)V' = 0$$

WE SOLVE TO OBTAIN  $w_0 = e^{-y^2/2} + y \int_{-\infty}^y e^{-s^2/2} ds$ .

THIS GIVES THE SOLUTION

$$W_0 = C_1 y + C_2 \left( e^{-y^2/2} + y \int_{-\infty}^y e^{-s^2/2} ds \right) \quad \int_{-\infty}^{\infty} e^{-s^2/2} ds = (2\pi)^{1/2}$$

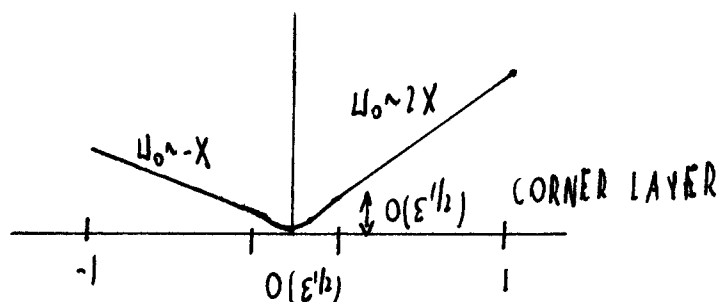
NOW  $W_0 \rightarrow 2y$  AS  $y \rightarrow +\infty$  GIVES  $C_1 + C_2 (2\pi)^{1/2} = 2$

$W_0 \rightarrow -y$  AS  $y \rightarrow -\infty$  GIVES  $C_1 = -1$ .

HENCE  $C_2 = 3/(2\pi)^{1/2}$

$$W_0 = -y + \frac{3}{(2\pi)^{1/2}} \left( e^{-y^2/2} + y \int_{-\infty}^y e^{-s^2/2} ds \right)$$

AND  $W = \varepsilon^{1/2} W_0 + \dots$



### EXAMPLE (INDETERMINACY: CASE OF (ii)).

THIS EXAMPLE SHOWS AN APPARENT FAILURE OF BOUNDARY LAYER THEORY.

$$\varepsilon u'' - Xu' + u = 0 \quad -1 \leq X \leq 1$$

$$u(-1) = 1, \quad u(1) = 2$$

SINCE  $Q(X) = -X$  THERE IS NO INTERNAL LAYER NEAR  $X = 0$ .

INSTEAD WE HAVE BOUNDARY LAYERS NEAR  $X = \pm 1$ .

IN THE OUTER REGION WE EXPAND  $u = u_0 + \varepsilon u_1 + \dots$

THIS YIELDS THAT  $-Xu_0' + u_0 = 0 \rightarrow u_0 = \alpha X$   $\alpha$  ARBITRARY.

NOTICE THAT  $u_0'' = 0$  SO THAT  $u_0 = \alpha X$  SATISFIES

THE FULL EQUATION.

WE NOW CONSTRUCT A BL NEAR  $x = -1$ .

(27)

WE LET  $y = (x+1)/\epsilon$ ,  $w(y) = u(-1 + \epsilon y)$  AND EXPAND  $w = w_0 + \epsilon w_1 + \dots$

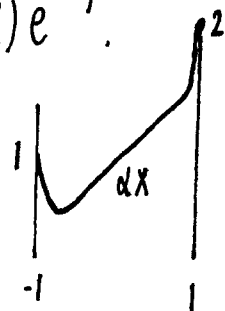
SO THAT

$$\left. \begin{aligned} w_0'' + w_0' &= 0 \\ w_0(0) &= 1, \quad w_0(\infty) = -\alpha \end{aligned} \right\} \Rightarrow w_0 = -\alpha + (\alpha+1)e^{-y}.$$

SIMILARLY NEAR  $x=1$  WE LET  $y = (1-x)/\epsilon$ ,  $v(y) = u(1 - \epsilon y)$

AND WE EXPAND  $v = v_0 + \epsilon v_1 + \dots$  SO THAT

$$\left. \begin{aligned} v_0'' + v_0' &= 0 \\ v_0(0) &= 2, \quad v_0(\infty) = \alpha \end{aligned} \right\} \rightarrow v_0(y) = \alpha + (2-\alpha)e^{-y}.$$



THIS GIVES A COMPOSITE OR UNIFORM SOLUTION

$$u_{\text{comp}} = \alpha x + (\alpha+1)e^{-(x+1)/\epsilon} + (2-\alpha)e^{-(1-x)/\epsilon}$$

A FUNDAMENTAL DIFFICULTY IS THAT WE CANNOT DETERMINE  $\alpha$ .

### REMARKS

(i) A HIGHER ORDER ANALYSIS TO INCORPORATE TERMS OF ORDER  $O(\epsilon^2)$  FOR SOME  $\epsilon$  DOES NOT DETERMINE  $\alpha$ .

(ii) THIS PHENOMENON WAS OBSERVED BY ACKBERG-O'MALLEY (1970), AND GENERATED  $1/\epsilon$  PAPER OVER THE YEARS.

(iii) SIMILAR DIFFICULTIES OCCUR FOR

$$\epsilon u'' + x u' = 0$$

$$u(-1) = 1, \quad u(1) = 2.$$

THE "SOLUTION" IS BY BL THEORY:

$$u_{\text{comp}} = \alpha + (1-\alpha)e^{-(x+1)/\epsilon^{1/2}} + (2-\alpha)e^{-(1-x)/\epsilon^{1/2}}$$