## Gaussian deconvolution and the lace expansion

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## Convolution equations on $\mathbb{Z}^d$

• 
$$f * g(x) = \sum_{y \in \mathbb{Z}^d} f(y)g(x - y)$$

• Random walk two-point function (d > 2): Let  $D(x) = \frac{1}{2d} \mathbb{1}\{|x| = 1\}$  and  $\delta(x) = \delta_{0,x} = \mathbb{1}\{x = 0\}$ , then

$$C(x) = \sum_{n=0}^{\infty} D^{*n}(x) = \delta_{0,x} + D(x) + D * D(x) + \cdots$$

### satisfies the convolution equation $C = \delta + D * C$ .

• (Bond) Percolation connection probability (d large): Let  $\tau_p(x) = \mathbb{P}_p(0 \leftrightarrow x)$ , then for  $p \leq p_c$ ,

$$\tau_p = \delta + \Pi_p + pD * (\delta + \Pi_p) * \tau_p.$$

• Self-avoiding walk two-point function (d > 4): For  $z \leq z_c$ ,

$$G_z = \delta + zD * G_z + \Pi_z * G_z.$$

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## Deconvolution

• Random walk: Since  $C(x) = \sum_{n=0}^{\infty} D^{*n}(x)$  satisfies  $C = \delta + D * C$ , we can rearrange it into

$$(\delta - D) * C = \delta.$$

So C is the *deconvolution* of the operator  $\delta - D$ , which is minus the discrete Laplacian. We also call C(x) the *lattice Green function*. It is well-known that as  $|x| \to \infty$ ,

$$C(x) = \frac{a_d}{|x|^{d-2}} + O\left(\frac{1}{|x|^d}\right), \qquad a_d = \frac{d\Gamma(\frac{d-2}{2})}{2\pi^{d/2}},$$

• We consider the convolution equation

$$F * G = \delta$$

with a given  $F : \mathbb{Z}^d \to \mathbb{R}$  and prove  $G(x) \sim \text{const} \cdot |x|^{-(d-2)}$  under some assumptions on F.

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## Fourier transform

• Let  $\mathbb{T}^d = (\mathbb{R}/2\pi\mathbb{Z})^d$  be the continuum torus, which we identify with  $(-\pi,\pi]^d \subset \mathbb{R}^d$ . We will use the  $L^1$  Fourier transform

$$\hat{f}(k) = \sum_{x \in \mathbb{Z}^d} f(x) e^{ik \cdot x} \qquad (k \in \mathbb{T}^d)$$

and the inverse Fourier transform

$$f(x) = \int_{\mathbb{T}^d} \hat{f}(k) e^{-ik \cdot x} \frac{dk}{(2\pi)^d} \qquad (x \in \mathbb{Z}^d).$$

- We will also use the  $L^2$  Fourier transform.
- Random walk example:

$$C(x) = \int_{\mathbb{T}^d} \frac{e^{-ik \cdot x}}{1 - \hat{D}(k)} \frac{dk}{(2\pi)^d}, \qquad \hat{D}(k) = d^{-1} \sum_{j=1}^d \cos k_j$$

(cf. 
$$(\delta - D) * C = \delta$$
)

## Main result: Gaussian deconvolution

We solve  $F * G = \delta$  using Fourier integral  $G(x) = \int_{\mathbb{T}^d} \frac{e^{-ik \cdot x}}{\hat{F}(k)} \frac{dk}{(2\pi)^d}$ .

### Theorem (Hara'08, L.–Slade'23)

Let d > 2. Suppose  $F : \mathbb{Z}^d \to \mathbb{R}$  is a  $\mathbb{Z}^d$ -symmetric function, and suppose there are  $K_1, K_2 > 0$ ,  $\rho > \max(0, \frac{d-8}{2})$  such that, for all  $x \in \mathbb{Z}^d$ and  $k \in \mathbb{T}^d$ ,

$$|F(x)| \le \frac{K_1}{|x|^{d+2+\rho}}, \qquad \hat{F}(0) = 0, \qquad \hat{F}(k) - \hat{F}(0) \ge K_2 |k|^2.$$

Then

$$G(x) \sim \frac{a_d}{\kappa |x|^{d-2}} \qquad as \ |x| \to \infty,$$

where  $\kappa = -\sum_{x \in \mathbb{Z}^d} |x|^2 F(x) \in (0, \infty).$ 

We do not assume  $F(x) \leq 0$  for  $x \neq 0$ .

## Gaussian deconvolution

• The theorem was first proved by Hara in 2008 using intricate Fourier analysis, without the assumption that  $\rho > \frac{d-8}{2}$  (only requiring  $\rho > 0$ ) in

$$|F(x)| \le \frac{K_1}{|x|^{d+2+\rho}}.$$

This extra assumption is satisfied for all known applications.

• The theorem directly applies to self-avoiding walk. For percolation, we combine the theorem with an elementary convolution estimate. We obtain

$$G_{z_c}(x), \tau_{p_c}(x) = \frac{\text{const}}{|x|^{d-2}} + O\left(\frac{1}{|x|^{d-\varepsilon}}\right)$$

with arbitrary  $\varepsilon > 0$ . Hara obtained  $\varepsilon = 2 - 2/d$ .

- Our proof is completely different and is short and simple. It is inspired by the work of Slade on weakly self-avoiding walks in 2022. But to cover percolation, we need new ideas.
- The decay assumption on F(x) can be replaced by regularity assumptions on  $|x|^{2+\varepsilon}F(x)$  and  $|x|^{d-2}F(x)$ .
- (Ongoing) Extension to models on  $\mathbb{R}^d$ , e.g., random connection model.
- (Ongoing) Anisotropic  $|x|^{-(d-2)}$  decay (using only  $\mathbb{Z}_2$ -symmetry).

## Strategy of proof

Recall 
$$G(x) = \int_{\mathbb{T}^d} \frac{e^{-ik \cdot x}}{\hat{F}(k)} \frac{dk}{(2\pi)^d}$$
 and  $\kappa = -\sum_{x \in \mathbb{Z}^d} |x|^2 F(x) \in (0, \infty).$ 

We decompose

$$\hat{G} := \frac{1}{\hat{F}} = \kappa^{-1} \frac{1}{1 - \hat{D}} + \frac{(1 - \hat{D}) - \kappa^{-1} \hat{F}}{(1 - \hat{D}) \hat{F}} = \kappa^{-1} \frac{1}{1 - \hat{D}} + \frac{\hat{E}}{(1 - \hat{D}) \hat{F}},$$

where  $E = A - \kappa^{-1}F$  with  $A = \delta - D$ . The constant  $\kappa$  is chosen to make

$$\sum_{x \in \mathbb{Z}^d} |x|^2 E(x) = 0,$$

so that the remainder would be more regular then the leading term. By inverse Fourier transform, since  $C(x) = \int_{\mathbb{T}^d} \frac{e^{-ik \cdot x}}{1 - \hat{D}(k)} \frac{dk}{(2\pi)^d}$ , we get

$$G(x) = \kappa^{-1}C(x) + f(x),$$

where f is the inverse Fourier transform of  $\hat{f} := \hat{E}/(\hat{A}\hat{F})$ .

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#### From

$$G(x) = \kappa^{-1}C(x) + f(x),$$

and

$$C(x) = \frac{a_d}{|x|^{d-2}} + O\left(\frac{1}{|x|^d}\right),$$

it suffices to prove  $f(x) = o(|x|^{-(d-2)})$  as  $|x| \to \infty$ .

The choice of  $\kappa$  allows us to take d-2 (weak) derivatives of  $\hat{f} = \hat{E}/(\hat{A}\hat{F})$ . We will show all these derivatives are integrable, then by the Riemann–Lebesgue lemma, we get  $|x|^{d-2}f(x) \to 0$  as  $|x| \to \infty$ .

## Intuition

### Since $E = A - \kappa^{-1}F$ is symmetric and satisfies

$$\sum_{x \in \mathbb{Z}^d} E(x) = \sum_{x \in \mathbb{Z}^d} |x|^2 E(x) = 0,$$

we roughly have  $\nabla^{\gamma} \hat{E}(k) \leq |k|^{2+\sigma-|\gamma|}$  for some  $\sigma \in (0, \min\{\rho, 2\})$ . By the assumed infrared bound, we have

$$\left|\frac{1}{\hat{A}(k)}\right|, \left|\frac{1}{\hat{F}(k)}\right| \lesssim \frac{1}{|k|^2}.$$

Taking derivatives roughly gives

$$\left| \nabla^{\gamma} \left( \frac{1}{\hat{A}(k)} \right) \right|, \left| \nabla^{\gamma} \left( \frac{1}{\hat{F}(k)} \right) \right| \lesssim \frac{1}{|k|^{2+|\gamma|}}.$$

Then by the product rule, we get

$$|\nabla^{d-2}\hat{f}| = \left|\nabla^{d-2}\left(\frac{\hat{E}}{\hat{A}\hat{F}}\right)\right| \lesssim \frac{|k|^{2+\sigma}}{|k|^{2+2+d-2}} = \frac{|k|^{\sigma}}{|k|^{d}} \in L^{1}(\mathbb{T}^{d}).$$

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- The intuition works for self-avoiding walk (Slade 2022) but does not work for percolation, where we cannot take enough classical derivatives ( $\Pi(x)$  does not decay fast enough).
- Solution: We use weak derivatives, and replace power-counting by Hölder's inequality.

## Weak derivative

Let  $C_c^{\infty}(\mathbb{T}^d)$  denote the space of infinitely differentiable, compactly supported *test functions*  $\phi : \mathbb{T}^d \to \mathbb{R}$ . (For the torus  $\mathbb{T}^d$ , every function has compact support.)

### Definition (Weak derivative)

Suppose  $u, v \in L^1(\mathbb{T}^d)$  and  $\alpha$  is a multi-index. We say that v is the  $\alpha^{\text{th}}$ weak partial derivative of u, written  $\nabla^{\alpha} u = v$ , if, for all test functions  $\phi \in C_c^{\infty}(\mathbb{T}^d)$ ,

$$\int_{\mathbb{T}^d} u \nabla^{\alpha} \phi = (-1)^{|\alpha|} \int_{\mathbb{T}^d} v \phi.$$

The requirement is the usual integration by parts formula, so u is weakly differentiable if it is classically differentiable.

#### Lemma

The weak derivative satisfies the usual product and quotient rules, provided the result is integrable.

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## Weak derivative and Fourier transform

For us, we just need the fact that the  $L^2$  Fourier transform gives the weak derivative. We write  $\mathcal{F}[f] = \hat{f}$  for the  $L^2$  Fourier transform of  $f \in \ell^2(\mathbb{Z}^d)$ .

#### Lemma

Let  $f : \mathbb{Z}^d \to \mathbb{R}$  and  $\alpha$  be a multi-index. Suppose  $x^{\alpha}f(x) \in \ell^2(\mathbb{Z}^d)$ . Then the  $\alpha^{\text{th}}$  weak partial derivative of  $\hat{f}$  is given by

$$\nabla^{\alpha} \hat{f} = \mathcal{F}[(ix)^{\alpha} f(x)]$$

We use the lemma to make sense of  $\nabla^{d-2}\hat{F}(k)$ . This is the origin of our restriction  $\rho > \frac{d-8}{2}$  on

$$|F(x)| \le \frac{K_1}{|x|^{d+2+\rho}};$$

we need  $|x|^{d-2}F(x) \in \ell^2(\mathbb{Z}^d)$ .

## Proof of main result

We want to show  $\hat{f} = \frac{\hat{E}}{\hat{A}\hat{F}}$  is d-2 times weakly differentiable. By the product and quotient rules,  $\nabla^{\alpha}\hat{f}$  is given by a linear combination of terms of the form

$$\left(\prod_{n=1}^{i} \frac{\nabla^{\delta_n} \hat{A}}{\hat{A}}\right) \left(\frac{\nabla^{\alpha_2} \hat{E}}{\hat{A} \hat{F}}\right) \left(\prod_{m=1}^{j} \frac{\nabla^{\gamma_m} \hat{F}}{\hat{F}}\right),$$

where  $\alpha = \alpha_1 + \alpha_2 + \alpha_3$ ,  $0 \le i \le |\alpha_1|$ ,  $0 \le j \le |\alpha_3|$ ,  $\sum_{n=1}^i \delta_n = \alpha_1$ , and  $\sum_{m=1}^j \gamma_m = \alpha_3$ , provided these terms are integrable.

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where  $\alpha = \alpha_1 + \alpha_2 + \alpha_3$ ,  $0 \le i \le |\alpha_1|, 0 \le j \le |\alpha_3|, \sum_{n=1}^i \delta_n = \alpha_1$ , and  $\sum_{m=1}^j \gamma_m = \alpha_3$ , provided these terms are integrable.

#### Lemma

Let 
$$|\gamma| < \frac{1}{2}d + 2 + \rho$$
, and choose  $\sigma \in (0, \rho)$  such that  $\sigma \leq 2$ . Then

$$\frac{\nabla^{\gamma}\hat{A}}{\hat{A}}, \frac{\nabla^{\gamma}\hat{F}}{\hat{F}} \in L^{q} \quad (q^{-1} > \frac{|\gamma|}{d}), \qquad \frac{\nabla^{\gamma}\hat{E}}{\hat{A}\hat{F}} \in L^{q} \quad (q^{-1} > \frac{2-\sigma+|\gamma|}{d}).$$

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By Hölder's inequality,

$$\left(\prod_{n=1}^{i} \frac{\nabla^{\delta_n} \hat{A}}{\hat{A}}\right) \left(\frac{\nabla^{\alpha_2} \hat{E}}{\hat{A} \hat{F}}\right) \left(\prod_{m=1}^{j} \frac{\nabla^{\gamma_m} \hat{F}}{\hat{F}}\right) \in L^r(\mathbb{T}^d)$$

as long as

$$\frac{1}{r} > \frac{\sum_{n=1}^{i} |\delta_n|}{d} + \frac{2 - \sigma + |\alpha_2|}{d} + \frac{\sum_{m=1}^{j} |\gamma_m|}{d} = \frac{|\alpha| + 2 - \sigma}{d}.$$

Since  $|\alpha| \leq d-2$  and  $\sigma > 0$ , we can take r = 1. This proves that  $\hat{f}$  is d-2 times weakly differentiable and concludes the proof.

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## Bound on $\nabla^{\gamma} \hat{A} / \hat{A}$ .

Recall  $A = \delta - D$  has finite support. If  $|\gamma| = 1$ , by Taylor's theorem and symmetry, we have  $|\nabla^{\gamma} \hat{A}(k)| \leq |k|$ . If  $|\gamma| \geq 2$ , Taylor's theorem gives  $|\nabla^{\gamma} \hat{A}(k)| \leq 1$  instead. Together with the infrared bound, we get

$$\left|\frac{\nabla^{\gamma}\hat{A}}{\hat{A}}(k)\right| \lesssim \frac{1}{|k|^{\min(|\gamma|,2)}} \in L^q(\mathbb{T}^d) \qquad (q^{-1} > \frac{\min(|\gamma|,2)}{d}),$$

which is stronger than the desired result.

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### Bound on $\nabla^{\gamma} \hat{F} / \hat{F}$ .

The  $|\gamma| = 1$  case is the same as for  $\hat{A}$ , because  $\sum_{x} |x|^2 |F(x)|$  is finite. For  $|\gamma| \ge 2$ , the decay assumption  $|F(x)| \le |x|^{-(d+2+\rho)}$  and boundedness of the Fourier transform imply

$$\nabla^{\gamma} \hat{F} \in L^{\frac{d}{|\gamma|-2}}(\mathbb{T}^d) \qquad (2 \le |\gamma| < \frac{1}{2}d + 2 + \rho).$$

Since  $|\hat{F}^{-1}(k)| \leq |k|^{-2} \in L^p$  for all  $p^{-1} > 2/d$  by the infrared bound, it follows from Hölder's inequality that  $\nabla^{\gamma} \hat{F} / \hat{F} \in L^q$  for all  $q^{-1} > (|\gamma| - 2 + 2)/d$ , as desired.

## Bound on $\nabla^{\gamma} \hat{E} / (\hat{A}\hat{F})$ .

Let  $\sigma \in (0, \rho)$  be such that  $\sigma \leq 2$ . We use the fact that  $E = A - \kappa^{-1}F$ has the same  $|x|^{-(d+2+\rho)}$  decay as F. If  $|\gamma| < 2 + \sigma$ , it follows from

$$\sum_{x \in \mathbb{Z}^d} E(x) = \sum_{x \in \mathbb{Z}^d} |x|^2 E(x) = 0,$$

symmetry, and infrared bounds that

$$\left|\frac{\nabla^{\gamma}\hat{E}}{\hat{A}\hat{F}}(k)\right| \lesssim \frac{|k|^{2+\sigma-|\gamma|}}{|k|^2|k|^2} = \frac{1}{|k|^{2-\sigma+|\gamma|}},$$

which is in  $L^q$  for  $q^{-1} > (2 - \sigma + |\gamma|)/d$ , as desired. If  $|\gamma| \ge 2 + \sigma$ , we use the Fourier transform to bound  $\nabla^{\gamma} \hat{E}$ , then use Hölder's inequality (as in the  $|\gamma| \ge 2$  case for  $\hat{F}$ ).

This concludes the proof of the lemma.

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### Theorem (Hara'08, L.–Slade'23)

Let d > 2. Suppose  $F : \mathbb{Z}^d \to \mathbb{R}$  is a  $\mathbb{Z}^d$ -symmetric function, and suppose there are  $K_1, K_2 > 0$ ,  $\rho > \max(0, \frac{d-8}{2})$  such that, for all  $x \in \mathbb{Z}^d$ and  $k \in \mathbb{T}^d$ ,

$$|F(x)| \le \frac{K_1}{|x|^{d+2+\rho}}, \qquad \hat{F}(0) = 0, \qquad \hat{F}(k) - \hat{F}(0) \ge K_2 |k|^2.$$

Then

$$G(x) \sim \frac{a_d}{\kappa |x|^{d-2}} \quad as \ |x| \to \infty,$$
  
where  $\kappa = -\sum_{x \in \mathbb{Z}^d} |x|^2 F(x) \in (0, \infty).$ 

We have proved  $G(x) = \kappa^{-1}C(x) + f(x)$  and  $\nabla^{d-2}\hat{f} \in L^1(\mathbb{T}^d)$ . Error estimate?

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## Better error estimate

We can improve the error to  $f(x) = O(|x|^{-(d-2+\delta)}), \delta > 0$ , by taking more derivatives of  $\hat{f}$ .

For fractional powers of |x|, we use the following integral representation: For  $\delta \in (0, 1)$ ,

$$(\operatorname{sgn} x_1)|x_1|^{\delta} = \frac{1}{c_{\delta}} \int_0^\infty \frac{\sin(x_1 u)}{u^{1+\delta}} du, \quad c_{\delta} = \int_0^\infty \frac{\sin u}{u^{1+\delta}} du \in (0,\infty).$$

Multiplying by  $\sin(x_1 u)$  produces phase shifts in the Fourier space.

### Lemma (Fractional derivative)

Let  $\tilde{u} = (u, 0, ..., 0)$ . Suppose that  $\hat{g} \in L^1(\mathbb{T}^d)$  and that

$$\frac{1}{2ic_{\delta}}\int_0^\infty \frac{1}{u^{1+\delta}} \left\| \hat{g}(\cdot + \tilde{u}) - \hat{g}(\cdot - \tilde{u}) \right\|_{L^1(\mathbb{T}^d)} du < \infty.$$

Then  $\sup_{x \in \mathbb{Z}^d} |x_1|^{\delta} |g(x)| < \infty$ .

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Then  $\sup_{x \in \mathbb{Z}^d} |x_1|^{\delta} |g(x)| < \infty$ .

We use the lemma with  $\hat{g} = \nabla^{\alpha} \hat{f}$  where  $|\alpha| = d - 2$ . Write

$$(U_u\hat{g})(k) = \hat{g}(k+\tilde{u}) - \hat{g}(k-\tilde{u}).$$

Estimates on  $U_u(\nabla^{\alpha} \hat{f})$  then lead to more decay of f(x).

# Estimates on $U_u(\nabla^{\alpha} \hat{f})$

Since  $\nabla^{\alpha} \hat{f}$  is given by a linear combination of terms of the form

$$\left(\prod_{n=1}^{i} \frac{\nabla^{\delta_n} \hat{A}}{\hat{A}}\right) \left(\frac{\nabla^{\alpha_2} \hat{E}}{\hat{A} \hat{F}}\right) \left(\prod_{m=1}^{j} \frac{\nabla^{\gamma_m} \hat{F}}{\hat{F}}\right),$$

and  $U_u$  is taking a finite difference, we apply  $U_u$  to one factor at a time.

#### Lemma

Let  $\gamma$  be a multi-index,  $0 \leq \eta \leq 1$ , with  $|\gamma| + \eta < \frac{1}{2}d + 2 + \rho$ . Choose  $\sigma \in (0, \rho)$  such that  $\sigma \leq 2$ , and choose  $q_1, q_2$  satisfying

$$q_1^{-1} > \frac{|\gamma| + \eta}{d}, \qquad q_2^{-1} > \frac{2 - \sigma + |\gamma| + \eta}{d}.$$

Then for  $0 \le u \le 1$ ,

$$\left\| U_u(\frac{\nabla^{\gamma}\hat{A}}{\hat{A}}) \right\|_{q_1}, \ \left\| U_u(\frac{\nabla^{\gamma}\hat{F}}{\hat{F}}) \right\|_{q_1}, \ \left\| U_u(\frac{\nabla^{\gamma}\hat{E}}{\hat{A}\hat{F}}) \right\|_{q_2} \lesssim u^{\eta}.$$

Estimates on  $U_u(\nabla^{\alpha} \hat{f})$ : an ingredient

$$(U_u\hat{g})(k) = \hat{g}(k+\tilde{u}) - \hat{g}(k-\tilde{u}).$$

#### Lemma ("Sobolev inequality")

Let  $g: \mathbb{T}^d \to \mathbb{C}$  be weakly differentiable. Fix  $1 \leq p < d$ . Assume  $\nabla^{e_j}g \in L^p(\mathbb{T}^d)$  for all j. Let  $0 \leq \eta \leq 1$  and define  $p_\eta$  by  $\frac{1}{p_\eta} = \frac{1}{p} - \frac{1-\eta}{d}$ . Then

 $||U_u g||_{p_\eta} \lesssim u^\eta ||g||_{W^{1,p}},$ 

where  $||g||_{W^{1,p}} = (||g||_p^p + \sum_{j=1}^d ||\nabla_j g||_p^p)^{1/p}.$ 

- Y. Liu and G. Slade. Gaussian deconvolution and the lace expansion. Preprint, arXiv:2310.07635.
- T. Hara. Decay of correlations in nearest-neighbor self-avoiding walk, percolation, lattice trees and animals. Ann. Probab., 36:530–593, (2008).
- G. Slade. A simple convergence proof for the lace expansion. Ann. I. Henri Poincaré Probab. Statist., **58**:26–33, (2022).

If you want to learn lace expansion for spread-out models:

• Y. Liu and G. Slade. Gaussian deconvolution and the lace expansion for spread-out models. Preprint, arXiv:2310.07640.

# Thank You!

The method can be extended to study a family of convolution equations,

$$F_z * G_z = \delta.$$

With  $F_z$  satisfying similar "massive" assumptions, we prove the uniform upper bound

$$G_z(x) \le \frac{c_0}{\max(1, |x|^{d-2})} e^{-c_1 m(z)|x|},$$

where m(z) is the exponential decay rate of  $G_z(x)$ , for  $z \in [z_c - \delta, z_c)$ ,  $\delta > 0$ .

The result applies to strictly self-avoiding walk in dimensions d > 4.

Reference: Y. Liu. A general approach to massive upper bound for two-point function with application to self-avoiding walk torus plateau. Preprint, arXiv:2310.17321.

